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ADVANCED H_2 – HCI GAS DYNAMIC LASER PHASE II

FINAL TECHNICAL REPORT

January 1976

Submitted to:

Dr. David Howgate
U.S. Army Missile Command
Redstone Arsenal, Alabama 35809

Under Contract No. DAAH01-75-C-0789

Submitted by:

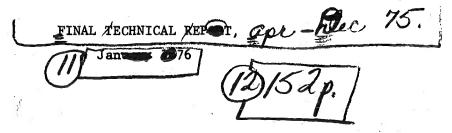
Atlantic Research Corporation 5390 Cherokee Avenue Alexandria, Virginia 22314

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ATLANTIC RESEARCH CORPORATION

ALEXANDRIA, VIRGINIA · 22314

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PHASE II.



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FOREWARD

This final technical report on an "Advanced H₂ - HCl Gas Dynamic Laser" was prepared by Atlantic Research Corporation, 5390 Cherokee Avenue, Alexandria, Virginia, 22314. The Project Engineer was Dr. Robert Cavalleri, the program was directed by Robert Naismith. The major subcontractor was Physical Sciences Inc., 18 Lakeside Office Park, Wakefield, Massachusetts, 01880. The principal investigator of that effort was Dr. Raymond Taylor. Technical support was provided by Stephen Scheffee of Atlantic Research Corporation and Dr. David Rosen, John Coughlin and Dr. Peter Wu, of Physical Sciences, Inc.

The effort was conducted from April 1975 through December 1975.

The effort was administered under the Army Missile Command, Redstone Arsenal, Alabama. Dr. David Howgate was the technical monitor.

This report was submitted February 3, 1976 and has been reviewed and is approved for publication.

SUMMARY

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The initial phase of this work was a feasibility study of an $H_2^{7/2}$ - HCl mixing gas dynamic laser. The results of that study indicated that a high energy laser system based on the $H_2^{7/2}$ - HCl system was indeed feasible. This laser device concept envisions the use of a solid propellant to generate high temperature $H_2^{7/2}$ - HCl mixtures which are expanded supersonically to create the proper conditions for efficient $H_2^{7/2}$ + HCl V \rightarrow V transfer and lasing on HCl. This type of laser device has several significant potential advantages as a high energy device for military application, i.e., the use of compact, field storable solid fuel technology, high specific power (80 - 100 Kj/lb), and lasing of HCl in the 4.00 atmospheric window.

The objective of the current phase was to experimentally verify the V - V and V - T kinetics for the important $\rm H_2$ - HCl processes and to develop solid propellant gas generators that could be used to generate the required gases at attractive weight yields. The results of the current phase have verified that the kinetics are capable with operation of a $\rm H_2$ - HCl laser. In addition, the gas generator development program has succeeded in producing the required gases at weight yields that make the system attractive for Army field application. The potential of even higher gas yields that would make the system even more attractive is also possible but requires additional propellant development. The fundamental remaining problem area is the effect of H-atoms on performance.

TABLE OF CONTENTS

			* · · · · · · · · · · · · · · · · · · ·		Page
1.0	INTR	ODUCTION		•	1
2.0	KINE	rics			2
	2.1	Backgrou	nd		2
	2.2	Experimen	ntal Technique		4
	2.3	Experimen	ntal Results	•	9
	2.4	H-Atom K	inetics	•	20
	2.5	H-HC1 Dea	activation Experiment Design		25
	2.6	HC1 Spec	tral Characteristics	• ;	31
3.0	PROP	ELLANT DE	VELOPMENT		35
	3.1	Objective	e		35
	3.2	Selection	n of Propellant Ingredients		37
		3.2.1	Oxidizers		45
		3.2.2	Fuels		47
		3.2.3	Coolants	•	47
		3.2.4	Chlorides		49
V	3.3	Calculat	ion of Theoretical Yield		50
		3.3.1	Procedure		50
		3.3.2	Discussion of Results		52
		3.3.2.1	Compositions Containing NH ₄ NO ₃	•	53
		3.3.2.2	Compositions Containing N2H5NO3	•	56
		3.3.2.3	Compositions Containing Either H ₂ WO ₄ or WO ₃		56
	3.4	Propertion	es of Candidate Propellants	•	63
		3.4.1	General	•	63
		3.4.2	Properties of Composition 3.3.4	•	63
		3.4.3	Properties of Composition 3.8.1	•	69
		3.4.4	Properties of Composition 3.8.8	•	69
	3.5	Sub-Scale	Motor Firings	•	71
		3.5.1	General	•	71
		3.5.2	Sub-Scale Motor Configurations	•	77
		3.5.3	Sub-Scale Motor Firings Results		80

(Table of Contents - Continued)

									Page
3.0	PI	ROPELLANT DE	VELOPMENT (Co	ntinued)		6 · 1			
		3.5.3.1	Composition (H ₂ WO ₄ /ZrH ₂ /		NH ₄ C10	₄)			80
		3.5.3.2	Composition (NH ₄ NO ₃ /ZrH ₂		/NH ₄ C1	04)			82
		3.5.3.3	Composition (H ₂ WO ₄ /Zr/HD		₄)				84
	3.	6 Summary			• • •			a sin	88
4.0	PI	ROOF EXPERIM	ENT AND LASER	SYSTEM	CHARAC	TERIST	ics	• •	90
	4.	1 Method o	f Characteris	ticsAnal	ysis .	B 2		•	103
	4.	2 Boundary	Layer Analys	is	• • •	· • •			109
	4.	3 Propella	nt Impurity E	ffects .				• •	117
	, 4.	4 Power Ex	traction						119
	4.	5 Advanced System C	H ₂ -HC1 Gas D hafacteristic	ynamic L s	aser	• • • •			126
5.0	CC	NCLUSIONS				8			134

(LIST OF FIGURES - CONTINUED)

FIGURE	PAG
3.1	Effect of Temperature, Pressure and Concentration on H ₂ Dissociation
3.2	Burning Rate of Composition 3.3.4 67
3.3	Burning Rate of Composition 3.8.1
3.4	Burning Rate of HDB
3 .5	Burning Rate of Composition 22 - 19.258% H ₂ WO ₄ + 54.612% Zr + 26.103% NH ₄ ClO ₄
3.6	Sub-Scale Motor Configuration
3.7	Motor in Cyclone Chamber
3.8	Motor Firing Traces
4.1	Performance Curves - Original Contour h _{th} = 0.01, T _o = 3000°K
4.2	Performance Curves - Original Contour h _{th} = 0.015, T _o = 3000°K
4.3	Performance Curves - Original Contour Contour-Mod II, Contour-Mod III
4.4	Characteristic Times for Various Energy Transfer Processes as a Function of Expansion Ratio
4.5	Optimized Performance Map for Baseline Condition of $T_s = 2800 ^{\circ}\text{K}$ and $P_s = 40 \text{atm} .$
4.6	HCL Vibrational Temperature
4.7	Nozzle Subsonic Section Density Distribution 100
4.8	H-Atom Concentration in Subsonic Nozzle Region 101
4.9	H-Atom Concentration
4.10	Characteristic Mesh Details
4.11	Mod VIII Nozzle Contour,θ Versus
4.12	Nozzle Contour Test Case
4.13	Nozzle Contour Test Case
4.14	Variation of N with Re_{θ} for Axisymmetric Air and Nitrogen Nozzle Data
4.15	Comparison Between Theoretical and Experimental F Cf Versus F _{Rd} Re ₀ . ©, Experiments, Adiabatic; O, Experiments with Heat Transfer: - Theory

LIST OF FIGURES

FIGURE	NO.	PAGE
2.0	Shock Tube Facility	5
2.1	Sample Fluorescence and Shock Speed Oscillograms	6
2.2	Optical System Used on Shock Tube	7
2.3	HCl-Ar Shock Tube Run. Semilogarithmic Plot of HCl Vibrational Energy Relaxation Behind Shock	10
2.4	Reduced Particle Relaxation Time versus Temperature for 5% HCl - 95% Ar Mixture	11
2.5	Shock Tube Calculation of HCl Vibrational Energy Relaxation. Shows Regime of Temperature and H ₂ /HCl Concentration Ratio where Relaxation is not Sensitive to V-V rate	13
2.6	Comparison of Present Experimental Data for $HC1/H_2$ V-T Rate with Other Rate Expressions	15
2.7	Interpretation of Bott's Observed HC1/H ₂ Data in Terms of V-V and V-T Processes	16
2.8	Semilogarithmic Plot of Observed HCl Vibrational Relaxation for Mixture of High $\rm H_2/HCl$ Concentration Ratio and Comparison with Calculation for V-V Rate Determination.	18
2.9	Semilogarithmic Plot of Observed HCl Vibrational Relaxation for Mixture of High $\rm H_2/HCl$ Concentration Ratio and Comparison with Calculation for V-V rate Determination.	19
2.10	Interpretation of Bott's Observed $HC1/H_2$ Data in Terms of V-V and V-T Processes	21
2.11	Rate of Deactivation of HCl (v = 1) by H-atoms. Comparison of Experimental Data, Temperature, Extrapolations and Theory	24
2.12		26
2.13	Results of Flow Tube Kinetics Calculation for Design of H-Atom Deactivation Experiment	29
2.14	Shock Tube Relaxation Calculation with H-Atoms Showing Decoupling of HCl and ${\rm H_2}$ Relaxation	30
2.15	Einstein Coefficient for HC1 V-V-1 Transitions	32
2.16	Pressure Broadened Spectral Line Half Widths for HCl P-Branch Fundamental Bands	34

(LIST OF FIGURES - CONTINUED)

4.16	Test Problem Flow Configuration
4.17	Nozzle Boundary Layer Properties
4.18	Effect of N ₂ on Performance
4.19	Maximum Small Signal Gain for the V = 1 to 7 vibrational Levels of HCl. Critical Gain is Calculated for ℓ = 100 cm, R = 0.98, and L = 0.2
4.20	Saturated Gain for V = 1 to 4 Vibrational Levels of HC1. Cavity Conditions: £ = 100 cm, R = 0.98, L = 0.2 124
4.21	Calculated Intracavity Flux for V = 1 to 3 Vibrational Transition of HCl. Also Shown is the Total Flux, i.e., the Sum of All Levels, Including Those not Shown 125
4.22	Calculated Output Power for H.O. Calculation. Cavity Conditions ℓ = 100 cm, R = 0.98 and L = 0.2 127
4.23	Laser Device

LIST OF TABLES

TABLE 1	NO.			PAGE
2.1		H ₂ /HC1 Vibrational Energy Transfer Processes Harmonic Oscillator Model	•	3
2.2		Summary of Updated Vibrational Rate Constants for Harmonic Oscillator Model		22
2.3		Vibration Rate Constants for Atomic Species for H ₂ /HCl System	•	23
2.4		List of Relevant Rate Constants for H-Atom Flow Experiment	•	28
3.1		Equilibrium Yield Values and Heating Values @ P = 750 PSIA for Selected Propellant Ingredients	•	38-42
3.2		Glossary of Candidate Propellant Ingredients		43-44
3.3		Theoretical Values of Flame Temperature and Yield for Selected OMOX Compositions - Compositions Containing NH, NO, and Either Zr or ZrH,		54
3.4		Compositions Containing $\mathrm{NH_4NO_3}$ and Either Al or $\mathrm{AlH_3}$		55
3.5		Compositions Containing $N_2H_5NO_3$ and Either Zr or ZrH_2 .		57
3.6		Compositions Containing $N_2H_5NO_3$ and Either Al or AlH ₃ .		58
3.7		Miscellaneous Compositions Containing NH _L ClO _L		59
3.8		Compositions Containing H ₂ WO ₄ and Either Zr or ZrH ₂		60
3.9		Compositions Containing WO and Either Zr or ZrH2		61
3.10		Compositions Containing ZrO(NO ₃) ₂ ·2H ₂ O and Either Zr or ZrH ₂		62
3.11		Properties of Candidate Propellants		64
3.12		Combustion Products of Candidate Propellants		65
3.13		Properties of HDB and Mix 22		72
3.14		Combustion Products of Two-Compartment HDB Propellant		73
3.15		Subscale Motor Firing Results: Composition 3.8.1 (H ₂ WO ₄ /ZrH ₂ /Zr(OH) ₄ /NH ₄ ClO ₄)	•	81
3.16		Subscale Motor Firing Results: Composition 3.3.4 (NH ₄ NO ₃ /ZrH ₂ /Zr(OH) ₄ /NH ₄ ClO ₄)		83
3.17		Subscale Motor Firing Results: HDB + Composition 22 (H ₂ WO _/ /Zr/NH _/ ClO _/)	•	85

(LIST OF TABLES - CONTINUED)

TABLE	NO.		PAGE
4.1		Contour Characteristics	. 95
4.2		Boundary Layer Characteristics	• 113
4.3		Potential Gas Contaminants	118
4.4		Vibrational Kinetics for N ₂ /H ₂ /HCl Systems	. 120
4.5		Weight and Volume of Components	130
4.6		Weight and Volume of Gas Generator	. 131
4.7		Propellant System Comparison of Candidates	. 132
4.8		Weight and Volume of System	. 133

1.0 INTRODUCTION

The initial phase of this work was a feasibility study of an ${\rm H_2-HC1}$ mixing gas dynamic laser. The results of that study indicated that a high energy laser based on the ${\rm H_2-HC1}$ molecular system was indeed feasible. This laser device concept envisions the use of a solid propellant to generate high temperature ${\rm H_2-HC1}$ mixtures which are expanded supersonically to create the proper conditions for efficient ${\rm H_2\rightarrow HC1}$ V \rightarrow V transfer and lasing on HCl. This type of laser device has several significant potential advantages as a high energy device for military application, i.e., the use of compact, field storable solid fuel technology, high specific power (80-100 Kj/lb), and lasing of HCl in the 4.0 μ atmospheric window.

In the initial study phase the critical kinetic and propellant issues that required further research were identified. In the kinetics area, it was determined that the three most critical kinetic rates that control the laser performance were (i) the H₂/HCl V-VV process for which no data were then available, (ii) the HCl/H₂ V-T process for which only a room temperature data point had been determined, and (iii) the HCl/H V-T process for which only room temperature measurements existed. In the propellant area, it was determined that gas generators with the desired gaseous weight yields and gas composition must be developed.

The purpose of the present investigation (Phase II) was to provide experimental data on the rates of the $\rm H_2/HC1~V \rightarrow V$ and $\rm HC1/H_2~V \rightarrow T$ processes at temperatures in the range of $1000-2500\,^{\circ}\rm K$, to develop the gas generators with the required weight yield and gas composition characteristics as well as to perform additional analytical studies for the design of an appropriate proof experiment. The results of this work are reported in this document.

2.0 KINETICS

2.1 Background

As was emphasized in the early stages of the H₂-HCl GDL feasibility study¹, the vibrational energy transfer kinetics play a critical role in determining potential laser performance. The vibrational temperature, gain, and extractable power are all highly dependent on the rates for V-T and V-V exchange. The principal vibrational energy transfer processes used in the harmonic oscillator model are listed in Table 2.1.

There is presently a large body of experimental data on process (1) obtained from several different experiments $^{2-5}$ and covering a large temperature range from $150 \le T \le 2100$. The rate constant for process (2) has been measured at room temperature 6 , and, quite recently, some shock tubelaser fluorescence data by Bott 7 have become available which provide some additional information about process (2) up to $T = 770^{\circ}K$. These most recent data by Bott will be discussed later in this section.

Experimental data for V-T process (3) exist from three separate experiments - a low temperature experiment covering the temperature range from approximately 50° K to 450° K and two high temperature shock tube experiments covering the temperature range from 1100° K to 2700° K. These data allow us to make a reasonably good curve fit to the rate constant for process (3) over the temperature range of interest for the proposed H_{2} -HCl GDL.

While there is at present no experimental data in the literature for process (4), it is not expected that this reaction will make a significant contribution to the vibrational energy relaxation. According to theories of the SSH type, HCl is expected to be less efficient at deactivating vibrationally excited H₂ than H₂ itself. This fact coupled with the small concentration of HCl relative to H₂ in the laser medium make the contribution of process (4) negligible.

Perhaps the most critical, yet most poorly determined rate is that for the V-V exchange-process (5). Previous to the present experiment, no experimental data existed for the rate of process (5) and, therefore, the laser modeling calculations have had to be done using an estimated rate expression for this process derived from a combination of empirical correlation and theory 1,10.

TABLE 2.1

Harmonic Oscillator Model

V → T Processes

$$HC1 (v = 1) + HC1 \rightarrow HC1 (v = 0) + HC1 + 2886 cm^{1}$$
 (1)

HC1
$$(v = 1) + H_2 \rightarrow HC1 (v = 0) + H_2 + 2886 cm^{-1}$$
 (2)

$$H_2$$
 (v = 1) + $H_2 \rightarrow H_2$ (v = 0) + H_2 + 4159 cm⁻¹ (3)

$$H_2$$
 (v = 1) + HC1 \rightarrow H_2 (v = 0) + HC1 + 4159 cm⁻¹ (4)

V → V Process

$$H_2$$
 (v = 1) + HC1 (v = 0) \rightarrow H_2 (v = 0) + HC1 (v = 1) + 1273 cm⁻¹ (5)

On the basis of the above analysis, it is clear that the important molecular vibrational rate constants that required experimental investigation are those for the V-V process (5) and the V-T process (2). The present shock tube experiment was designed to make measurements of these rate processes over the temperature range of roughly 1000°K to 2500°K.

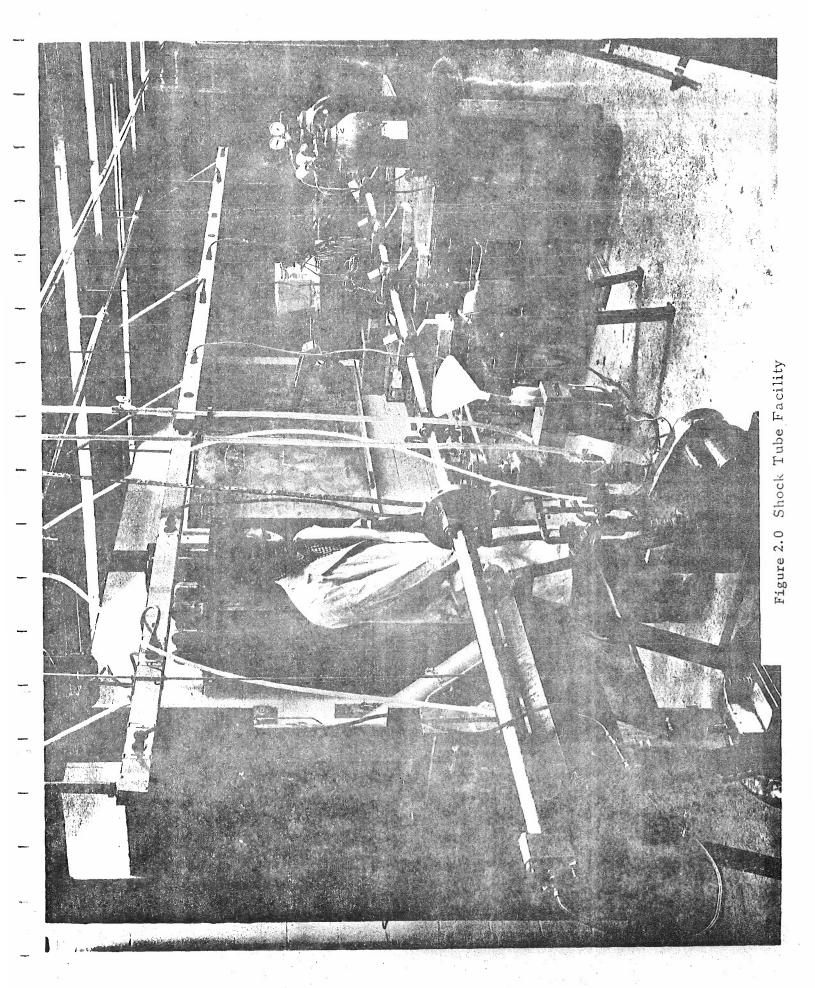
2.2 Experimental Technique

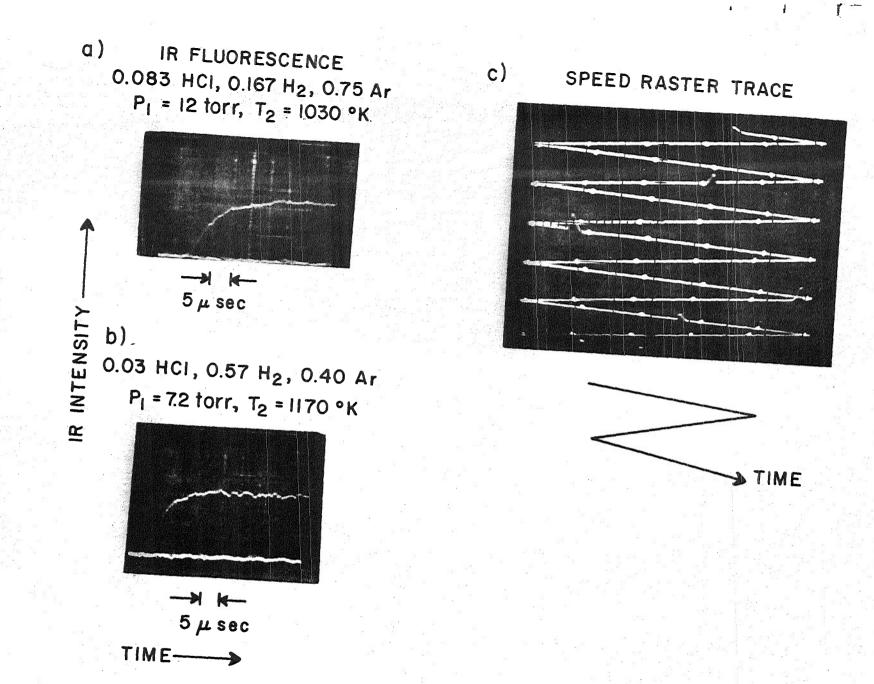
The experiments were carried out on a conventional 1.5-inch diameter shock tube. A photograph of the general shock tube facility is shown in Figure 2.0.

The basic technique employed monitoring the infrared fluorescence emitted by the HCl behind incident shock waves. It has been shown that the infrared radiation intensity from the fundamental vibrational rotation band is proportional to the vibrational energy in that mode. Therefore, following the time history of emission from the fundamental band is equivalent to measuring the rate of change of the vibrational energy.

The shock tube was constructed of 1.5-inch i.d. stainless steel with a 4-foot long driver section and a 15-foot driven section. Shock waves were produced by bursting scored aluminum and cold rolled steel diaphragms with helium. The shock speeds were measured with a series of six thin film, platinum heat transfer gauges mounted at various positions down the tube. The output from the heat transfer gauges were displayed on a speed raster sweep of a suitably modified Tektronix oscilloscope. A typical photograph of such a speed raster sweep with the differentiated outputs of the various heat transfer gauges is shown in Figure 2.1c.

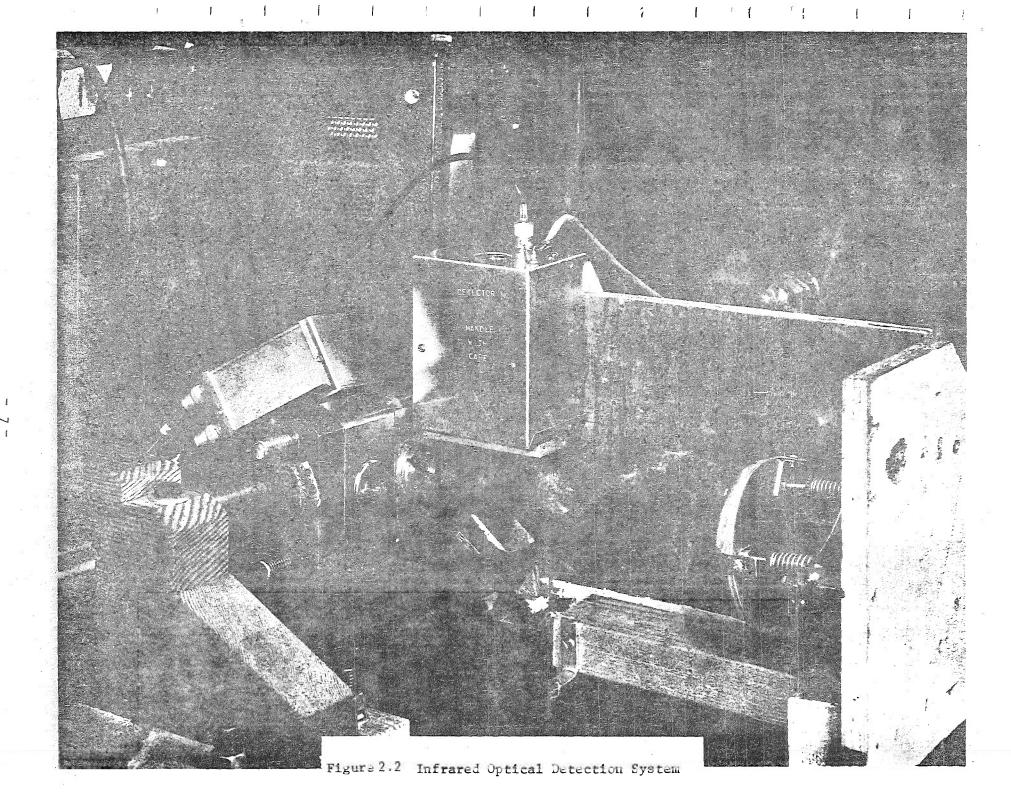
The infrared emission intensity was detected behind the incident shock by means of the optical system shown in Figure 2.2. The radiation first passed through a 0.50-inch diameter CaF_2 window contoured to match the inner radius of the shock tube. The light was then collected by a 10-cm diameter aluminum mirror with a 30-cm radius of curvature, directed to a flat aluminum mirror, and finally imaged with 1:1 magnification on a 1-mm diameter detector element. The detector was a liquid nitrogen cooled InSb detector operated in the photovoltaic mode. An infrared filter with a bandpass of 3.2 to 5.8 μ was placed before the detector to isolate the fundamental HCl emission.





6

Figure 2.1 Sample Fluorescence and Shock Speed Oscillograms



The test section of the shock tube was constructed of Lucite to minimize the stray light that might reach the detector by internal reflections off the metal shock tube walls. Before this Lucite section was utilized, light emission was detected before passage of the shock front by the optical station, presumably from radiation emitted down the tube ahead of the shock and reflected into the optical system by the back wall of the original metal test section. Changing the test section to Lucite, which is a poor reflector in the infrared, eliminated this extraneous light.

The output of the IR detector was amplified by a wide band low noise preamplifier of 20 Db gain before recording by a 545B Tektronix oscilloscope. A 3 MHz high frequency cut-off filter was employed immediately before the input to the final amplifying stage of the oscilloscope. The time response of the optical detection system and associated electronics was approximately 1μ sec. A sample of typical IR fluorescence traces is shown in Figures 2.1a and 2.1b.

Initial fill pressures, p, in the driven section were measured with a Wallace & Tiernan absolute pressure gauge (0 - 200 torr) and an oil manometer. Typical fill pressures ranged from about 3 torr to 20 torr. The gases used in this tudy were Argon: Matheson Purity, minimum purity 99.995%; HC1: Linde Specialty Gases L.E.D. grade, minimum purity 99.995%; and H2: Linde Specialty Gases, ultra high purity, minimum purity of 99.999% and a maximum moisture content of 3 ppm. The argon was further purified to remove any possible water content by passage through a copper coil immersed in a dry iceacetone bath. The HCl was subjected to a freeze-pump-thaw cycle between 77°K and 300°K to remove any residual H, that might have formed due to decomposition. This HCl was then passed through a copper coil at dry ice temperature to remove residual water vapor. Finally, the H2 was passed through a copper coil immersed in a liquid nitrogen bath to trap out any water vapor. The gas mixtures were prepared in a 4 liter stainless steel mixing tank at total pressures up to 100 psig and allowed to mix by diffusion for at least twenty-four hours before utilizing them in the experiments. The shock tube was capable of being pumped to 2×10^{-5} torr before each run with a combined leak/outgas rate of less than 7×10^{-4} torr/minute.

2.3 Experimental Results

Since the rates for vibrational relaxation of HCl in argon are well established 12,13 over a substantial temperature range, the first data were taken on HCl-Ar mixtures for the purpose of calibrating our system. As expected, the time resolved HCl fluorescence behind the incident shock was found to follow a single exponential behavior. Figure 2.3 is a semi-logarithmic plot of the normalized HCl emission intensity versus lab time for a 10% HCl-90% Ar mixture, an initial pressure 7.5 torr, and a shock speed of 1.17 mm/µsec. Fig. 2.4 summarizes the results obtained for a 5% HCl-95% Ar mixture. The data have been reduced to particle relaxation times referenced to a post shock pressure of 1 atmosphere for various post shock temperatures. The large Ar dilution assures an essentially constant post shock temperature during the relaxation. The solid line represents the results obtained by Seery for a similar mixture, and, as can be seen there is good agreement between our data and that of Seery. Data obtained in other HCl-Argon mixtures demonstrated similar good agreement.

With the HCl-Ar runs completed and confidence in our diagnostic system and gas handling procedures established, we proceeded to make shock tube runs for various $HCl-H_2$ -Ar mixtures.

The vibrational energy relaxation equations for a binary gas mixture of harmonic oscillators has been derived previously 14 and will be utilized here without further demonstration. The specific equations used are as follows:

$$-\frac{dX_{A}}{dt} = \left[\psi/\tau_{AA} + \frac{1-\psi}{\tau_{AB}} \right] X_{A}$$

$$-\frac{1-\psi}{\tau_{AB}^{e}} \left[\frac{X_{B}(1-X_{A})}{1-\exp(-\theta_{A})} - \frac{X_{A}(1-X_{B})}{1-\exp(-\theta_{B})} \right] [1-\exp(-\theta_{A})]$$

and

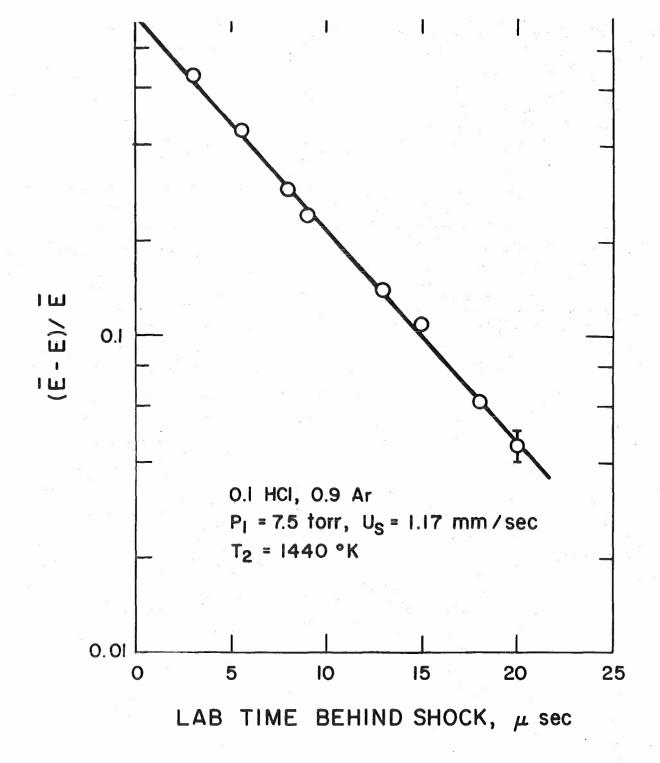


Figure 2.3 HC1-Ar Shock Tube Run. Semilogarithmic Plot of HC1 Vibrational Energy Relaxation Behind Shock.

TEMPERATURE °K

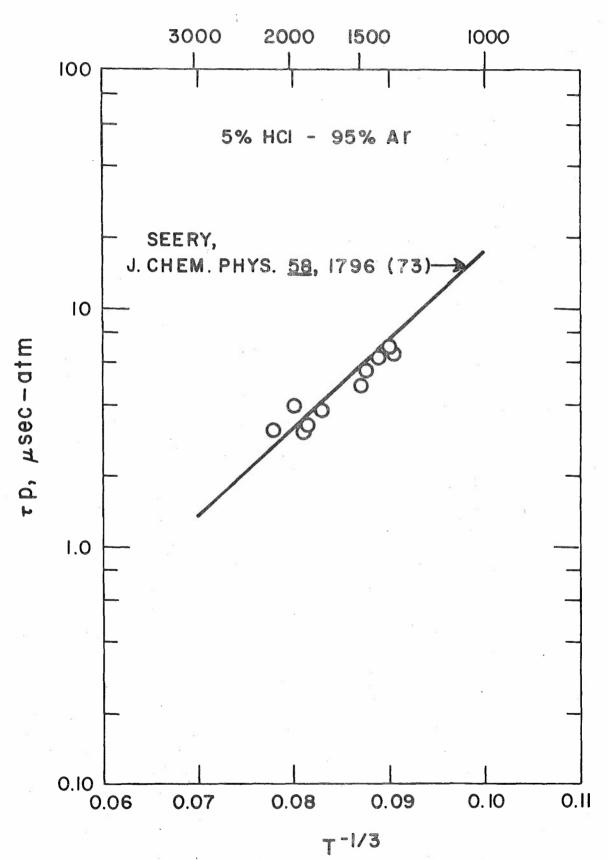


Figure 2.4. Reduced Particle Relaxation Times versus Temperature for 5% HCl - 95% Ar Mixture

$$-\frac{dX_{B}}{dt} = \left[\frac{1-\psi}{\tau_{BB}} + \frac{\psi}{\tau_{BA}}\right] X_{B} +$$

$$= \exp(\theta_{B}) \left[\frac{X_{B}(1-X_{A})}{X_{A}(1-X_{B})} + \frac{X_{A}(1-X_{B})}{X_{A}(1-X_{B})}\right] \times \left[1 - \exp(-\theta_{A})\right]$$

$$\frac{\psi}{\tau \xrightarrow{\text{e}} \frac{\exp(\theta_{\text{B}})}{\exp(\theta_{\text{A}})}} \times \left[\frac{X_{\text{B}}(1-X_{\text{A}})}{1-\exp(-\theta_{\text{A}})} - \frac{X_{\text{A}}(1-X_{\text{B}})}{1-\exp(-\theta_{\text{B}})} \right] \times \left[1-\exp(-\theta_{\text{A}}) \right]$$

where $X = (\overline{E} - E)/\overline{E}$, \overline{E} is the equilibrium vibrational energy at translational temperature T, and E is the vibrational energy at the local vibrational temperature T_v ; ψ and $1 - \psi$ are the mole fractions of species A and B respectively; and $\theta = hc\omega/kT$. Denoting H_2 as species A and HCl as species B, τ_{AA} , τ_{AB} , τ_{BB} , τ_{BA} , τ_{AB}^e would be the characteristic time constants for process (3), (4), (1), (2), and (5), respectively of Table 2.1.

While the above equations may be linearized to provide approximate solutions ^{15,16}, for the present analysis solutions were generated numerically by computer. With input of initial fill pressure, temperature (295°K), gas composition, and incident shock velocity, the computer program calculates the post-shock conditions and, depending on the values chosen for the unknown rate constants, predicts normalized temporal profiles of both the HCl and H₂ vibrational energies following passage of the incident shock.

From a number of initial computer simulation runs of various experimental conditions, it became apparent that determination of the V-V rate for process (5) from the experimentally observed relaxation profiles would require a reasonably accurate prior knowledge of the rate for V-T process (2). Further, it was determined that there was a regime of experimental conditions where the rate for process (2) could be extracted essentially independent of an evaluation of the V-V rate constant, namely, at relatively low H₂/HCl concentration ratios, i.e., 1:1, 2:1 and 3:1, and at the lower end of the experimental temperature range. Under such conditions, energy sharing between HCl and H₂ is minimal and the relaxation profile is only very weakly dependent on the V-V rate. This fact is illustrated by Figure 2.5 which is a calculation of the temporal profile of the HCl vibrational energy behind the incident shock for a 10% HCl-20% H₂-70% Ar mixture at an initial pressure of 10 torr, an incident shock speed of u₈ = 1.0 mm/usec, and a post shock temperature of 915°K. Comparison of the two

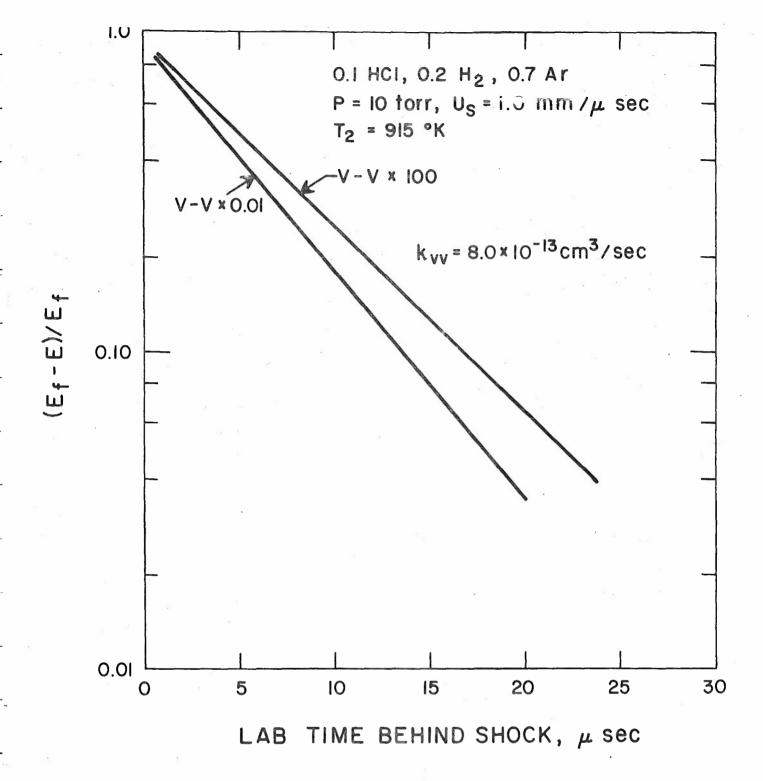


Figure 2.5. Shock Tube Calculation of HCl Vibrational Energy Relaxation. Shows Regime of Temperature and $\rm H_2/HCl$ Concentration Ratio Where Relaxation is not Sensitive to V-V Rate

calculated curves shows only a small difference in the predicted relaxation behavior even though the assumed V-V rate has been varied a factor of 100 around the baseline theoretical rate.

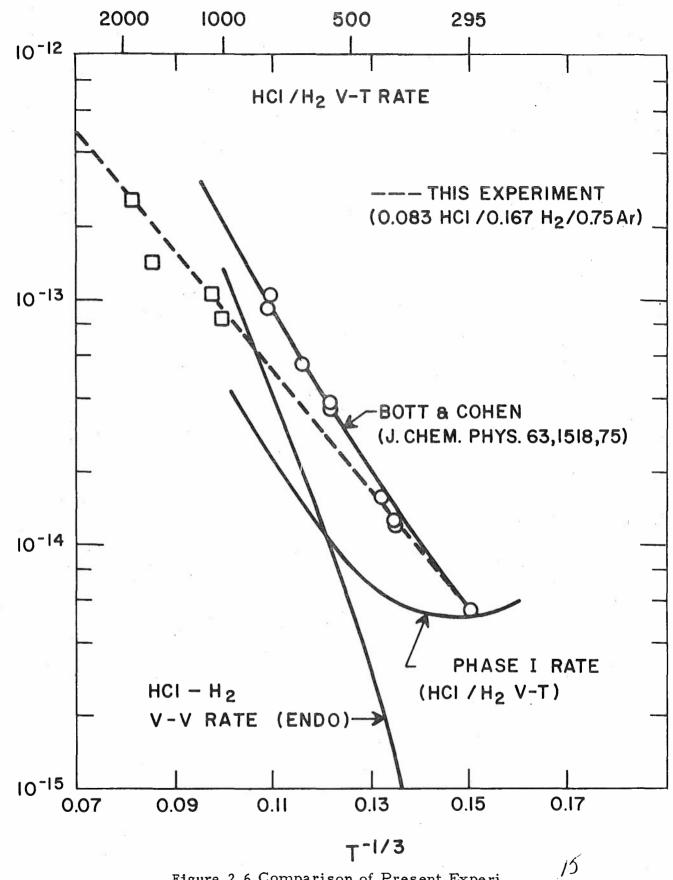
Such computer stimulations made clear the logical sequence of experiments to be performed to determine both the $\mathrm{HCl-H_2}$ V-T rate and the $\mathrm{HCl-H_2}$ V-V rate. First, runs would be made for mixtures with relatively low $\mathrm{H_2}/\mathrm{HCl}$ concentration ratios in order to ascertain the rate for V-T process (2). Following an accurate determination of this rate, runs would be made for mixtures of high $\mathrm{H_2}/\mathrm{HCl}$ concentration, where the relaxation is expected to be strongly dependent on the rate for the V-V coupling process.

A series of data was taken on a mixture of 8.3% HCl-16.7% H₂-75% Ar over a temperature range from 1000°K - 2000°K. Figure 2.1a is a sample oscillogram of one such run. A few of these runs were analyzed for the rate of V-T process (2), and the results are shown as the square data points in Figure 2.6. A curve fit to these results combined with the established room temperature measurement is best described by the expression:

$$k_2 = 2.5 \times 10^{-11} \exp(-55.9/T^{1/3}) \text{ cm}^3/\text{particle}^{-1} - \text{sec}^{-1}$$
 (8)

The circles shown on Figure 2.6 represent the data of Bott and Cohen between 300°K and 800°K. On comparing these two independent results, there appears to be a discrepancy. Better agreement between the two sets of data is possible if one takes into account the possible contribution of the H₂-HCl V-V rate process (5) to Bott and Cohen's observed HCl relaxation. Such an anlysis of Bott and Cohen's high temperature data is shown in Figure 2.7 and described later. As can be seen, assuming a combination of a V-V rate based on our theoretical estimate and an HCl-H₂ V-T rate based on the present results extrapolated to 770°K, one can predict a relaxation profile that follows Bott's observed decay quite closely over nearly two e-folding times. While such an interpretation of Bott's data is premature pending the experimental determination of the V-V rate constant, it does provide a preliminary explanation for the apparent discrepancy between the two experiments.

TEMPERATURE °K



RATE CONSTANT, k, cm3 - sec-

Figure 2.6 Comparison of Present Experimental Data for HC1/H2 V-T rate with Other Rate Expressions

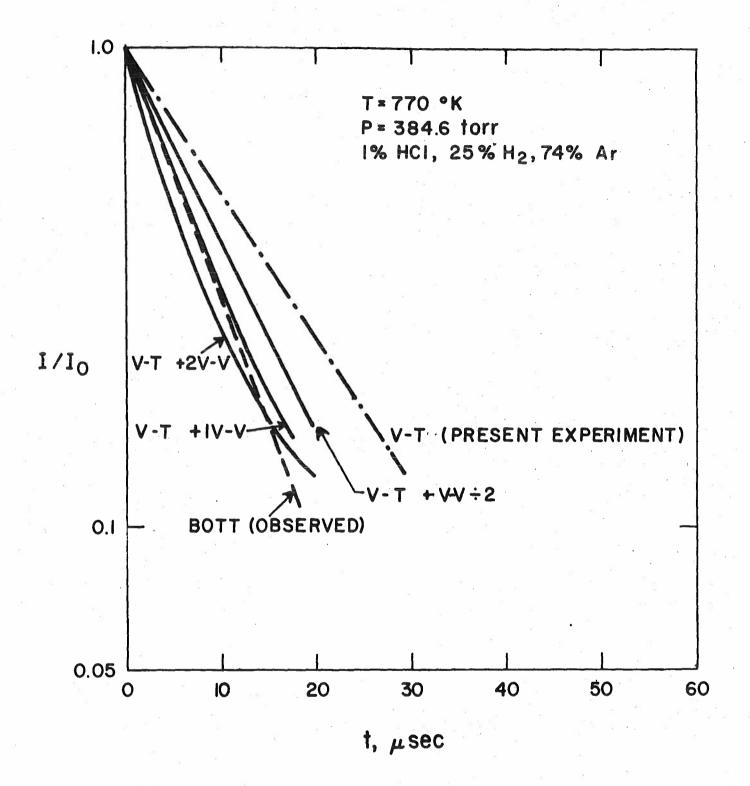


Figure 2.7. Interpretation of Bott's Observed $HC1/H_2$ Data in Terms of V-V and V-T Processes

Using the value for the HCl-H, V-T rate constant established from the present experiment, runs made in mixtures of high H2/HCl concentration were analyzed to obtain information on the value of the V-V rate constant. Plotted in Figure 2.8 is the intensity data from the oscillogram shown in Figure 2.1b taken in a 3% HC1-57% $H_0-40\%$ Ar mixture at a post shock temperature of 1170°K. The solid lines are computed curves for various values of the assumed V-V rate, the baseline value being the theoretically estimated rate expression. The dashed line indicates the relaxation predicted if the contribution from the V-V process were negligible. It is clear that the data plotted in Figure 2.8 indicate a V-V rate constant between 1/5 to 5 times the theoretical baseline value at 1170°K, with the best fit shown being that utilizing the baseline value. Figure 2.9 is another similar plot for a 2.6% HC1-47.4% H₂-50% Ar mixture at a post shock temperature of 1070°K. Again, the solid lines are computed temporal histories for various assumed values of the V-V rate, and the dashed line represents the relaxation profile that would be expected for a negligible contribution from the V-V process. It would appear that the experimental data obtained from this run would best be fit with a value of the V–V rate constant approximately 1/3 of the baseline value. Analysis of other shock tube data taken for mixtures of high H₂/HCl concentrations at various post shock temperatures indicated similar results; that is, the experimental data are best fit with a value of the V-V rate constant between 1/3 to 3 times the theoretically derived baseline value. Determination of the V-V rate constant to better precision will require; (a) a more precise determination of the $HC1-H_2$ V-T rate constant and, (b) acquisition of more highly resolved data taken for higher H2/HCl concentration ratios.

An additional test may be applied as a constraint on the value of the V-V rate constant. Using the value for the HCl-H₂ V-T rate obtained from the present experiments, one can analyze the data of Bott and Cohen near 800°k to see what value of the V-V rate constant is required to best fit their data. Presented in Figure 2.7 are the results of such an analysis. The three solid curves shown are predicted temporal profiles of the HCl laser induced fluorescence decay calculated for assumed values of the V-V rate constant of 1/2, 1 and 2 times the baseline theoretical value at 770°K. The dashed curve is the relaxation profile observed by Bott. It is apparent that the data can best

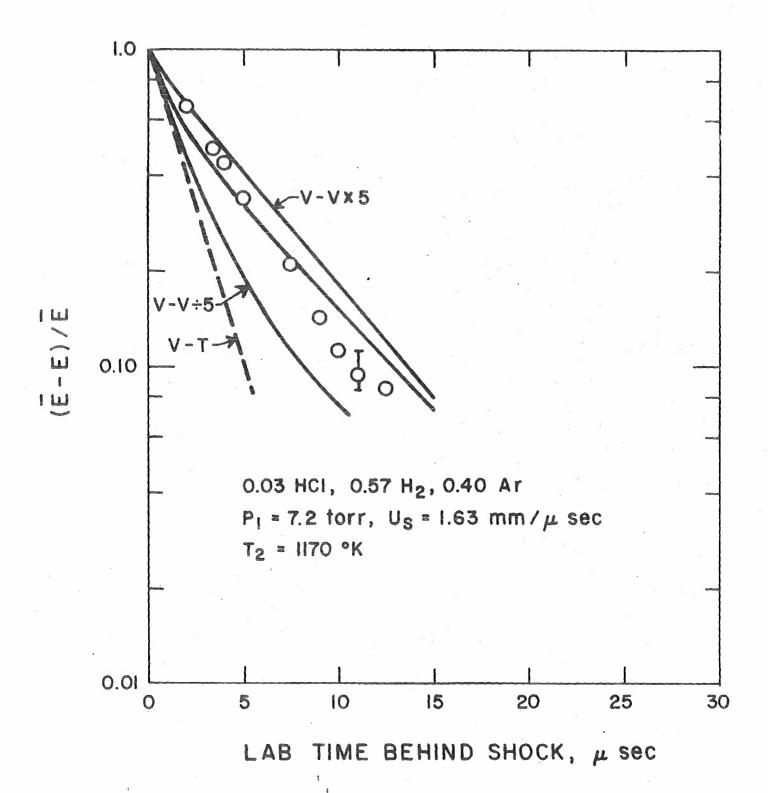


Figure 2.8 Semilogarithmic Plot of Observed HCl Vibrational Relaxation for Mixture of High H2/HCl Concentration Ratio and Comparison with Calculation for V-V Rate Determination

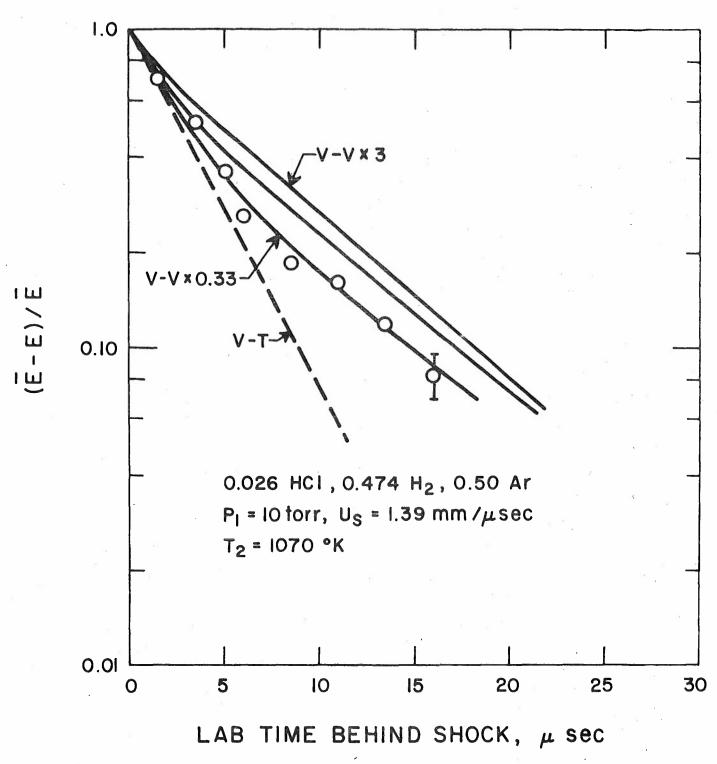


Figure 2.9 Semilogarithmic Plot of Observed HCl Vibrational Relaxation for Mixture of High H2/HCl Concentration Ratio and Comparison with Calculation for V-V Rate Determination

be fit with a value for the V-V rate constant very near the baseline value. A similar comparison made with Bott's data taken at T = 554°K is shown in Figure 2.10 and indicates essentially the same result - the data are best fit with a value for the V-V rate constant very near our previously assumed theoretical value. In addition, a recent as yet unpublished, laser fluorescent experiment in a HF-HC1-H₂ mixture at room temperature was reported by Bott²⁷. Analysis of these data for the rate of the H₂/HC1 V-V process yields a value which is within 70% of the base line value at 300°K. Therefore, it is concluded that the present experimental data support our theoretically derived H₂/HC1 V-V rate constant from 300 to ~1000°K. Table 2.2 lists the current values of rate constants for the vibrational energy exchange processes for the H₂/HC1 system.

2.4 H-Atom Kinetics

In the Phase I investigation the potentially serious performance degradation of atoms, particularly the H-atom, produced in the high temperature stagnation region and subsequently frozen in the rapid expansion, was discussed. At that time only one measurement of HCl (v = 1) deactivation by H-atoms was available 18 , and that measurement was at room temperature. Hence, a large and uncertain temperature extrapolation was necessary. Table 2.3 summarizes the previously derived rate expressions for the atomic rate processes.

Early in the current program an unpublished result was brought to our attention* that the rate in reference 18 had been confirmed²⁰. Working over the limited temperature range of 200 - 400°K, an activation energy of about 1 kcal/mole had also been determined. This information has been converted into a temperature dependent rate constant and is shown on Figure 2.11 and compared to the rate based on Equation (9) of Table 2.3. Clearly, the data of Glass et al, lead to a steeper temperature dependence and larger rate for temperatures above 300°K. Also shown in Figure 2.11 is the theoretical calculation by R. Wilkens²¹, which has been divided by 10. More recently, another measurement at room temperature ¹⁹ has provided a further confirmation of the low temperature rate. This result is also shown in Figure 2.11.

^{*}Private communication with Dr. Thomas Barr (USAMICOM)

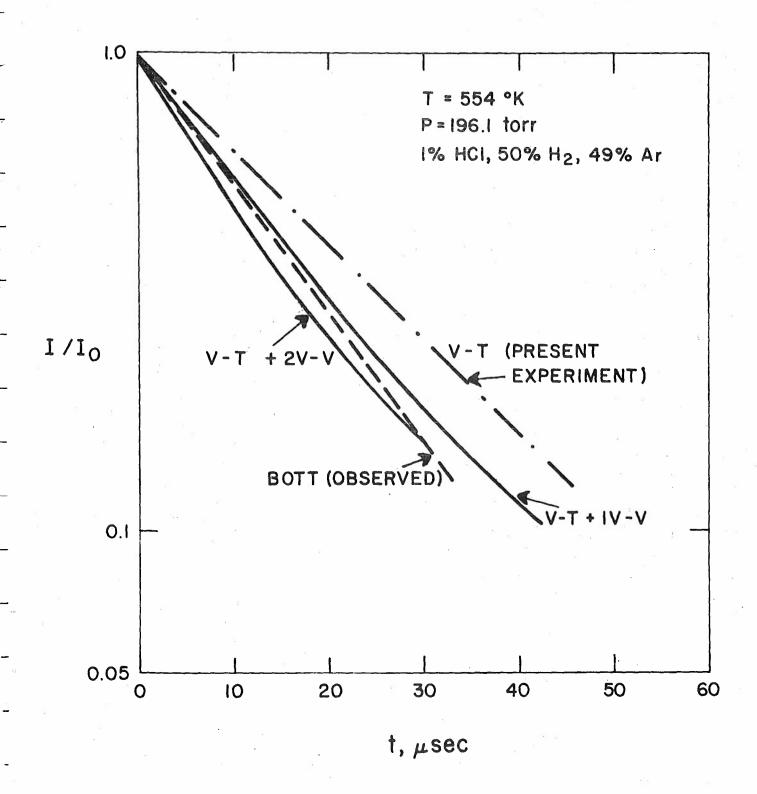


Figure 2.10 Interpretation of Bott's Observed HCl/H₂ Data in Terms of V-V and V-T Processes

TABLE 2.2

SUMMARY OF UPDATED VIBRATIONAL RATE CONSTANTS

FOR HARMONIC OSCILLATOR MODEL

V → T Processes

(1) HC1
$$(v = 1) + HC1 (v = 0) \rightarrow HC1 (v = 0) + HC1 (v = 0)$$

(2) HC1
$$(v = 1) + H_2 (v = 0) \rightarrow HC1 (v = 0) + H_2 (v = 0)$$

(3)
$$H_2$$
 (v = 1) + H_2 (v = 0) \rightarrow H_2 (v = 0) + H_2 (v = 0)

(4)
$$H_2$$
 (v = 1) + HC1 (v = 0) \rightarrow H_2 (v = 0) + HC1 (v = 0)

$k (cm^3 - part^{-1} - sec^{-1})$

$$k_1 = 4.68 \times 10^{-10} \exp (-76.8T^{-1/3})$$

$$+ 3.13 \times 10^{-16} \exp (27.9 \text{ T}^{-1/3})$$

$$k_2 = 2.5 \times 10^{-11} \exp (-55.9 \text{ T}^{-1/3})$$

$$k_3 = 2.18 \times 10^{-9} \exp (-114.9 \text{ T}^{-1/3})$$

$$k_{\Delta} = 4.35 \times 10^{-10} \exp (-114.9T^{-1/3})$$

V → V Process

(5)
$$H_2 9v = 1) + HC1 (v = 0) \rightarrow H_2 (v = 0) + HC1 (v = 1)$$

$$k_5 = 1.71^{+1.71} \times 10^{-14} \text{ T exp } (-30.0 \text{ T}^{-1/3})$$

TABLE 2.3

VIBRATION RATE CONSTANTS FOR ATOMIC SPECIES FOR H2/HC1 SYSTEM

(Rate constants in units of cm³ - particle⁻¹ - sec⁻¹)

(Rate expressions are for the forward process)

V → T Processes

HC1 (v = 1) + H
$$_{+}^{\rightarrow}$$
 HC1 (v = 0) + H (9)
k = 3.75 x 10⁻¹³ T^{1/2}

HC1
$$(v = 1) + C1 \xrightarrow{+} HC1 (v = 0) + C1$$
 (17)
 $k = 7.50 \times 10^{-13} T^{1/2}$

$$H_2 (v = 1) + H \xrightarrow{+} H_2 (v = 0) + H$$
 (18)
 $k = 1.73 \times 10^{-14} \text{ T}^{1/2}$

$$H_2$$
 (v = 1) + C1 $\stackrel{\rightarrow}{\leftarrow}$ H_2 (v = 0) + C1
 $k = 1.73 \times 10^{-15} \text{ T}^{1/2}$ (19)

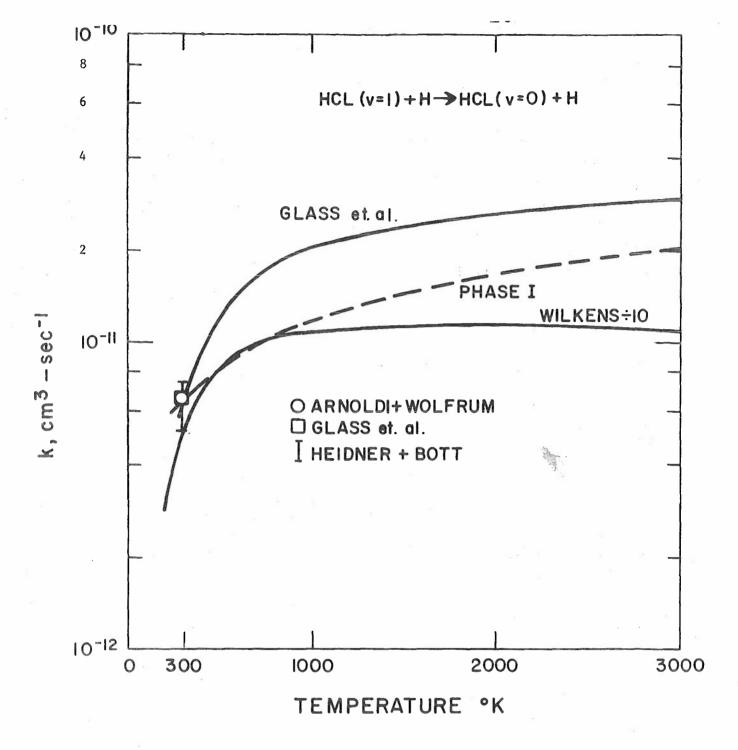


Figure 2.11 Rate of Deactivation of HC1 (v = 1)
by H-atoms. Comparison of
Experimental Data, Temperature
Extrapolations, and Theory

Clearly, the data at and around room temperature are in close agreement and support a large rate for the deactivation of HCl (v = 1) by H-atoms. The discrepancy between the measured values for HCl/H and HF/H still exists and requires adequate theoretical explanation. Nevertheless, the large room temperature value of HCl/H combined with the uncertainty in temperature dependence make this process of key importance in understanding potential laser performance. No additional data on the rates of other processes shown in Table 2.2 or 2.3 have become available during this program.

2.5 H-HCl Deactivation Experiment Design

An experiment to measure the rate of H-atom deactivation of HCl for T > 1000°K is shown schematically in Figure 2.12. The basic concept consists in converting the present 1.5 inch diameter shock tube into a glass flow reactor/shock tube. A primary flow of H₂ and Ar will be passed through a microwave discharge and mixed upstream with a secondary flow of HCl external to the shock tube. This mixture will then pass through the shock tube. Flow rates of H₂, Ar, and HCl will be measured using mechanical flow meters. H-atom concentrations will be determined by the standard NOCl titration technique 22. Once a steady flow of determined concentration is established, fast acting valves will close off the tube, and a shock wave will be initiated. I.R. emission from the HCl will then be monitored behind the incident shock wave using the InSb detector as in the present experiment

In designing the above experiment, various prereactions that might occur in the flow must be considered. These chemical reactions are:

$$H + HC1 \stackrel{\rightarrow}{\leftarrow} H_2 + C1 \tag{10}$$

$$C1 \xrightarrow{\text{Wall}} 1/2 C1_2$$
 (11)

$$H \xrightarrow{\text{wall}} 1/2 \text{ H}_2 \tag{12}$$

$$H + C1_2 \stackrel{?}{\leftarrow} HC1 + C1$$
 (13)

$$H + H + M \rightarrow H_2 + M \tag{14}$$

$$C1 + C1 + M \rightarrow C1_2 + M$$
 (15)

$$H + C1 + M \rightarrow HC1 + M \tag{16}$$

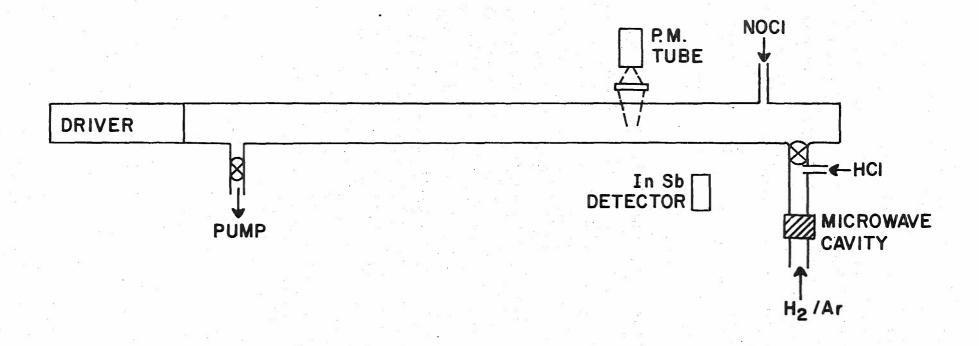


Figure 2.12 Shock Tube Flow Reactor for H-Atom Deactivation Experiment

Rate constants for these various reactions were obtained and a computer program used to calculate the flow history of the relevant species. Table 2.4 lists the values of the rate constants utilized in the program. For the flow pressures of interest and wall accommodation coefficients, γ , for H and Cl atoms of less than 10^{-3} , the rate limiting processes for wall removal of atoms are calculated to be reactions (11) and (12) and not diffusion. See reference (24) for a fuller explanation of the theory of wall effects on active species.

One of the principal results of the above flow kinetics calculations was that even at the lowest HCl concentrations permissible for adequate fluorescence signal and for the fastest flow speeds reasonably attainable through the shock tube, the time constant associated with reaction (10) is expected to be short compared to the flow time between the HCl injector and the test region. Such a situation is undesirable because it produces both a depletion of H-atoms and an undesirable concentration of Cl atoms. cluding that it is not possible to design a practical shock tube flow experiment to compete with the rapid rate of reaction (10), it was decided to prevent reaction (10) from proceeding too far in the forward direction by operating under conditions where the H2 concentration is large compared to HCl, and the backward reaction is favored. Shown in Figure 2.13 are the results of a calculation for a pre-shock pressure of 10 torr, an ambient temperature of 300°K, a flow velocity of 10^3 cm/sec, and a mixture of 1% HCl -33% $\rm H_2$ - 0.3% $\rm H$ - 65.7% Ar. Plotted are the concentration profiles as a function of flow distance down the tube. As can be seen, for a wall accomodation coefficient, $\gamma = 10^{-5}$ or 10^{-4} (values that can easily be achieved on a properly treated glass wall) 22, the H-atom concentration will remain essentially constant, and there will be negligible formation of C1-atoms.

Operating in mixtures of high $\rm H_2/HC1$ ratios does, however, present an additional potential problem. The energy sharing between $\rm H_2$ and HC1 due to the V-V process can complicate the interpretation of the effect of H-atom deactivation. That is, if HC1 and $\rm H_2$ are strongly coupled together, does one have to consider the contribution from H-atom deactivation of $\rm H_2$, as well? Figure 2.14 shows the results of a calculation which indicates an experimental regime where the interpretation of the observed HC1 relaxation is essentially independent of the effect of the H on $\rm H_2$. The calculation

TABLE 2.4

LIST OF RELEVANT RATE CONSTANTS FOR H-ATOM FLOW EXPERIMENT

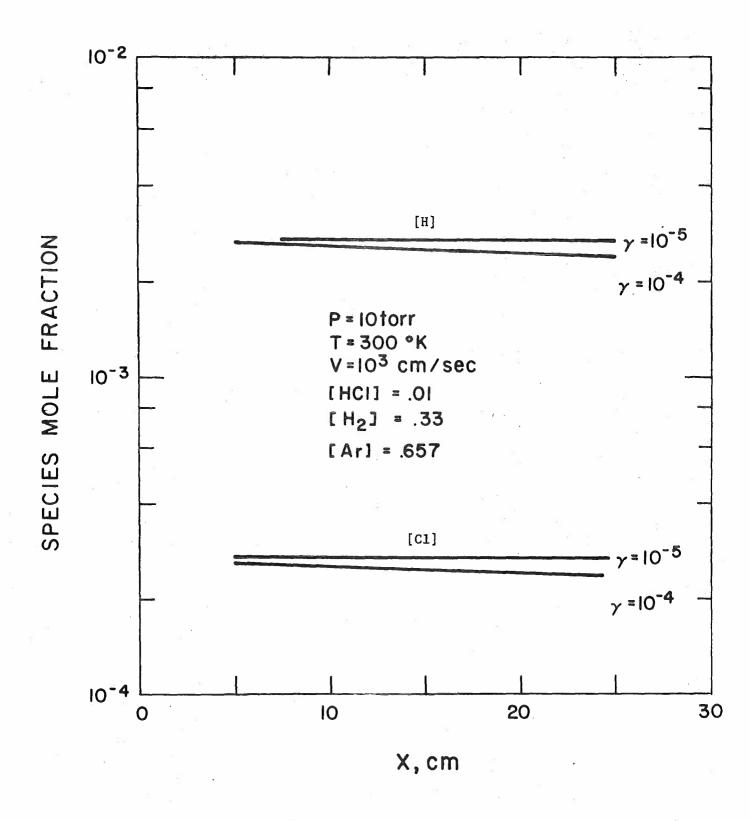


Figure 2.13 Results of Flow Tube Kinetics
Calculation for Design of H-Atom
Deactivation Experiment

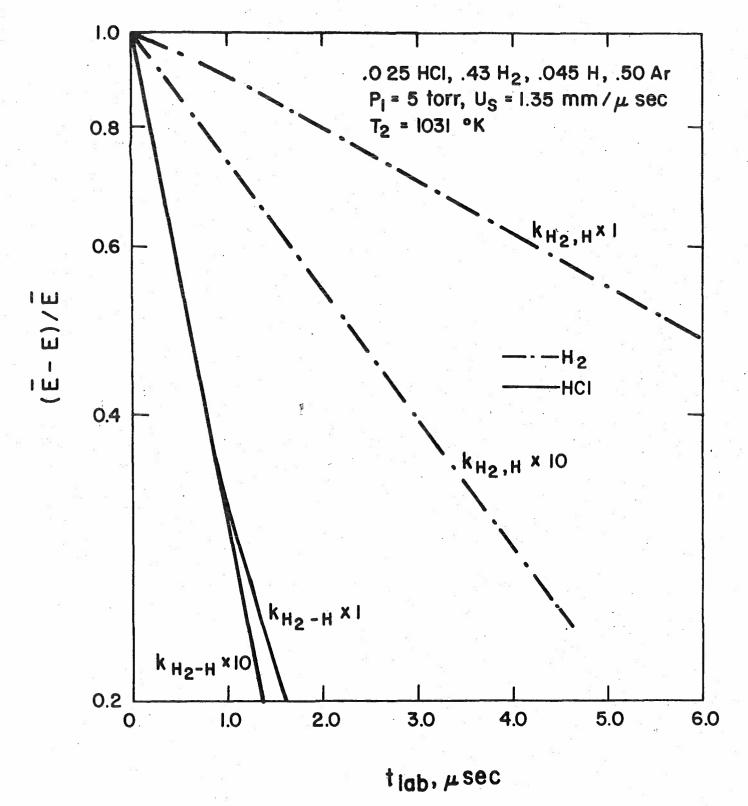


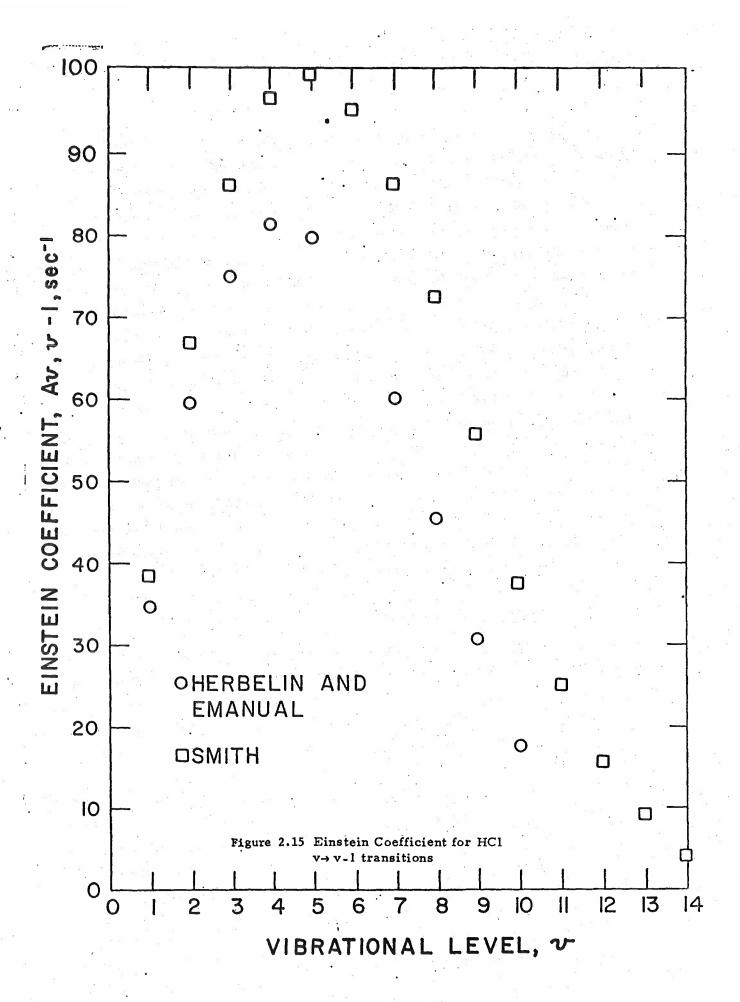
Figure 2.14 Shock Tube Relaxation Calculation with H-Atoms Showing Decoupling of HCl and H₂ Relaxation

was performed for an incident shock speed of 1.35 mm/usec in a 5 torr mixture of 2.5% HCl - 4.5% H - 43% H $_2$ - 50% Ar. Such conditions yield a post shock temperature of 1030°K. Even at this increased H-atom concentration, calculations indicate that 3 body recombination of H is still slow compared to the pertinent flow time. The two dash-dot curves plotted in Figure 2.13 show the predicted relaxation profile of the Ho vibrational energy assuming our present base line value for the H₂/H rate (upper curve), and a rate ten times that value (lower curve). The two solid curves plotted show the predicted relaxation of the HCl vibrational energy for, again, an assumed value of the H2/H deactivation rate equal to our baseline rate, and 10 times that value. Clearly, the value for the rate of H-atom deactivation of H, has little effect on the HCl relaxation profile under these conditions. Similar calculations were performed for other conditions, and the conclusion is reached that a determination of the rate for H-atom deactivation of HCl can be made independent of the rate for H-atom deactivation of H, provided that the energy sharing between H_2 and HC1 is not allowed to progress appreciably. Such a condition can be achieved by operating at the lower shock temperatures and/or creating sufficient H-atom loading to make the V-T processes fast compared to the rate for V-V coupling.

2.6 HCl Spectral Characteristics

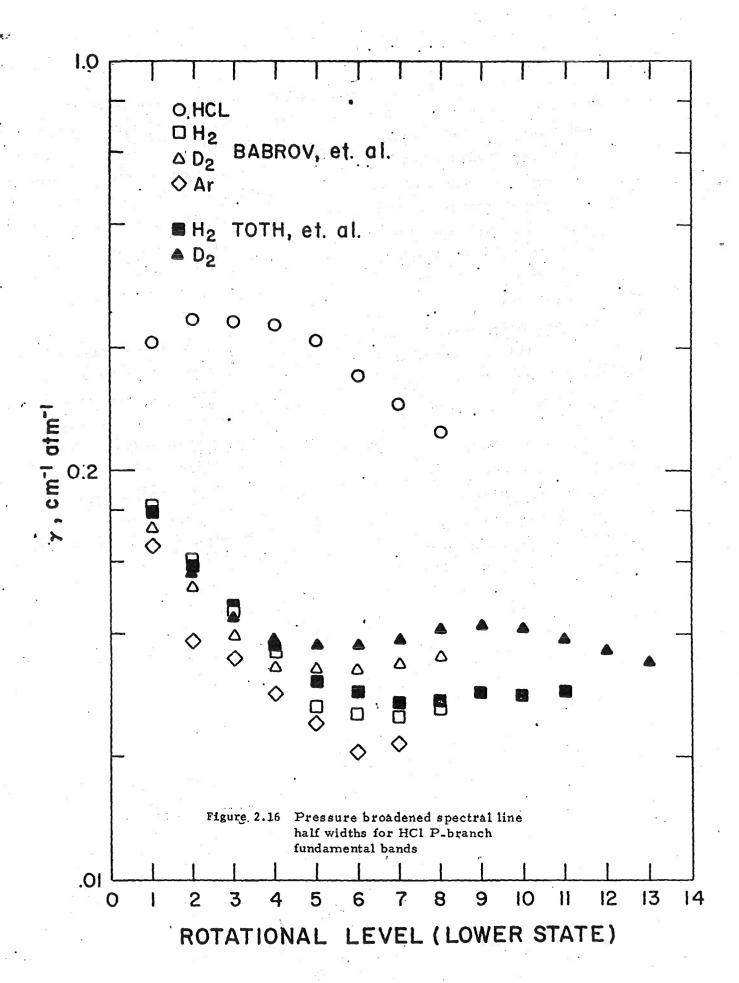
Besides the basic kinetic processes discussed above, it was decided to assess two other important parameters used in the modeling - the Einstein coefficients and line broadening cross-sections for HCl. Both of these latter parameters are needed in the calculation of gain and subsequent power extraction.

The experimental data on Einstein coefficients for the fundamental vibration/rotation band of HCl exist for only the lower vibrational levels. Theory must be used to extrapolate these data to the upper levels. The results of two such theoretical extrapolations ^{28,29} are shown in Figure 2.15. These results are in close agreement at the lower vibrational levels but exhibit an increasing divergence with increasing v. While this descrepancy is of some fundamental importance, its impact on the present laser modeling is small, since the emphasis is on energy extraction from the lower vibrational levels to provide laser wavelengths with good atmospheric transmission. For the present modeling, the Einstein coefficients of Smith²⁸ are used.



Another important physical parameter required in the laser modeling is the optical broadening cross-section. The experimental data base for HCl is somewhat limited and is shown in Figure 2.16 as the half width, γ , for the P-branch $v=1 \to 0$ transition $^{30},^{31}$. The collision broadening shows an interesting variation with rotational level. This type of variation has also been noted with other hydrogen halides such as HF 32 . Also, the self-broadening (HCl by HCl) is considerably larger than other homonuclear species such as H $_2$. Finally, the data of these two investigations for similar species are in close agreement. No data exist for the higher vibrational levels of HCl and for higher temperature.

Meredith and co-workers have extensively modeled the optical broadening cross-sections for HF and DF in response to the needs of chemical laser development. For example, they have shown 32 that the cross-section can vary as $T^{-1/2}$ to $T^{-1/4}$ depending upon the type of molecular interaction involved. Lacking this type of theoretical calculation for HCl, we have determined an optical broadening cross-section from the data base of Figure 2.16 and assumed no temperature or vibrational level dependence. For high J transitions (J > 8), we have assumed that the cross-section is constant and equal to its value at J = 8. Clearly, refinement of those assumptions is necessary as the concept of an H_2/HCl GDL advances.



3.0 PROPELLANT DEVELOPMENT

3.1 Objective

The objective of the propellant development is to develop solid propellant gas generators that yield combustion products suitable for use as the working fluid in an advanced H₂-HCl gas dynamic laser. There are certain basic constraints concerning the properties of the combustion products produced by the gas generators. These constraints are enumerated below:

1. Flame Temperature

Initially, in Phase I it was believed that the required flame temperature value was 3000°K, but nozzle optimization calculations for available power indicated that a decrease in flame temperature was permissable. In addition, at lower temperatures there are less H-atoms present. As shown in Figure 3.1, H-atoms increase with temperature due to H₂ dissociation. Thus, at a lower temperature the H-atom problem is somewhat alleviated. A parametric study of available power for a range of flame temperatures in the range 2400 to 3000°K indicated 2800°K to be the best compromise value.

2. Combustion Pressure

A chamber or combustion pressure of 750 psia was selected. This pressure, and the above flame temperature, are the stagnation values for the working fluid.

3. Gas Composition

The molar ratio selected ($\rm H_2/HC1$) in the combustion products is 97/3. Initially, values as low as 90/10 were considered. More recent results indicate that an optimum ($\rm H_2/HC1$) value is 97/3 and values as high as 98/2 may be desirable.

4. Gas Purity (Molar Content)

The $\rm H_2/HCl$ purity in the combustion products must be >90%. Initially, the requirement was for a best possible value, and hydrocarbons - which are an excellent off-the-shelf source of hydrogen - were considered seriously as propellant ingredients. However, subsequent calculations of the effect of contaminants primarily CO and $\rm N_2$ on gain coefficient and available power indicated that a relatively pure gas composition for the combustion products is required.

5. Condensed Species

The only way to generate the flame temperatures of interest is by the use of metal fuels (principally A1 and Zr). Hydrocarbons are not good fuels because combustion to C0 is not energetic enough while combustion in excess of CO yields H₂O, which is detrimental to laser performance. Combustion of A1 and Zr results in the production of A1₂O₃ and ZrO₂, both of which are condensed at the flame temperatures and pressure of interest, and must be removed from the combustion products, either by clinker formation or by mechanical means (filtration and/or inertial separation).

6. Weight Yield

The weight yield of the H₂/HCl mixture should be as high as possible by weight of propellant. The results of a preliminary laser systems analysis indicates that a 5% yield is desirable. This 5% yield, however, was not considered to be a hard requirement, and two propellants were developed which were satisfactory and feasible in all respects except with respect to weight yield.

The one propellant ingredient that is necessary to meet all of the above requirements is hydrazine diborane ($^{8}2^{N}2^{H}_{10}$ referred to as HDB throughout this report). The compound is available in limited quantities on a special order basis, but synthesis is a straightforward and a well-known process.

Three candidate propellant compositions were developed under this program. The three are pressed compositions comprised of an oxidizer, a fuel, a coolant and a chloride. Candidates are based upon selection of ingredients and computed values of theoretical flame temperature and combustion product composition. For a given set of ingredients, propellant composition is systematically varied to map ($\rm H_2$ + HCl) weight yield and purity versus the molar ratio ($\rm H_2/HCl)$) within the constraint $\rm T_f$ = 2800°K. Candidate compositions were then characterized for safety and ballistic properties and fired in subscale motors. Additionally, various filter arrangements were evaluated to retain condensed combustion products in the motor. The methodology and results of this work are described below.

TABLE 3.1 EQUILIBRIUM YIELD VALUES AND HEATING VALUES @ $P_c = 750$ PSIA FOR SELECTED PROPELLANT INGREDIENTS

*		librium Y: 00 gms @ :				Heating Vagms, @ T,	
PROPELLANT COMPOSITION, WT %	H ₂	HC1	Σng	3000	2800	2600	2400
40.467% N ₂ H ₅ NO ₃ + 59.533% ZrH ₂	1.685	· -	2.351	-19.98	-40.91	-47.98	-54.66
36.402% NH ₄ NO ₃ + 63.598% ZrH ₂	1.576		2.055	-18.77	-40.13	-46.78	-53.06
38.654% NH ₄ C10 ₄ + 61.346% ZrH ₂	1.144	0.3048	1.649	-47.50	-67.44	-73.25	-78.74
63.078% Zr(OH) ₄ + 36.922% ZrH ₂	1.181	e - -	1.195	52.70	30.48	25.17	20.11
45.025% Zro(NO ₃) ₂ ·2H ₂ 0 + 54.975% ZrH ₂	0.9180		1.100	-19.74	-40.82	-45.78	2
38.041% Zr(NO ₃) ₄ ·5H ₂ O + 61.959% ZrH ₂	1.098	, – .	1.291	-33.73	-55.14	-60.48	
62.375% WO ₃ + 37.625% ZrH ₂	0.4012	-	0.4059	6.762	- 4.599	- 7.347	- 9.986
57.264% H ₂ WO ₄ + 42.726% ZrH ₂	0.6835	-	0.6915	11.74	- 1.565	- 5.030	- 8.334
59.972% NC + 40.028% ZrH ₂	1.241	-	2.855	1.287	-15.40	-22.58	-29.42
57.776% TMETN + 42.224% ZrH ₂	1.450	_	2.948	- 8.813	-26.35	-33.83	-40.94
69.064% NitroGuanidine + 30.936% ZrH ₂	1.635	-	3.657	28.75	12.68	4.269	- 3.691
32.783% Al(NO ₃) ₃ ·9H ₂ 0 + 67.217% ZrH ₂	1.492	_ 3	1.650	- 6.505	-29.01		
52.730% A1(OH) ₃ + 47.270% ZrH ₂	1.506	· _	1.533	67.03	47.67		
* 'See Table 3.2 for Glossary of Terms		103	0*				2
	a et	22°	9			•	
	**		8			V	÷.
			# .		8		#10
		2 ·				8	

3.2 Selection of Propellant Ingredients

The propellants are comprised of oxidizer, fuel, coolant and chloride, each ingredient selected on the basis that the constraints of the problem are achieved. Candidate ingredients were selected on the basis of hydrogen content, chemical energy content, commercial availability, ease of use, stability of propellant ingredients and stability of combustion products.

For any given set of propellant ingredients, there is a unique composition that yields the desired values of flame temperature and $\rm H_2/HC1$ composition. To minimize the appearance of $\rm H_2O$ as a combustion product, the stoichiometry of candidate propellants was constrained to an oxidation ratio $\rm (O_R)$ of unity, i.e., the oxygen in the propellant is equal to that required to burn the metals to the metal oxides $\rm (ZrO_2,Al_2O_3)$ and the carbon (if present) to CO. The oxides must be stable with respect to $\rm H_2$ at the stagnation or chamber conditions to prevent formation of $\rm H_2O$. The resultant stoichiometry is referred to throughout this report as the OMOX point. Compositions fuel-rich of this stoichiometry are not particularly desirable because metal vapors and suboxides are produced as combustion products, which may interfere with lasing or result in condensed species in the optical cavity.

Equilibrium values of $\rm H_2$ yield, HCl yield, total gas yield ($\rm Eng$), and heating value ($\rm \Delta H_T^{\circ}$) were computed for selected OMOX combustions of: (1) oxidizer and fuel; (2) coolant and either oxidizer or fuel, and (3) chloride (and fuel if necessary). Calculations were carried out with the ARC equilibrium program using JANAF thermochemical data on an IBM 370/168 computer at assigned values of 750 psia combustion pressure and temperatures of 3000 to 2400°K with 200°K increments. The results of these calculations are given in Table 3.1. The heating values in the table are derived quantities which are defined as the enthalpy at equilibrium of the combustion products at a temperature T and 750 psia less the enthalpy of the propellant. Negative values of $\rm \Delta H_T^{\circ}$ indicate available heat while positive values indicate that heat is required to attain the temperature T. For example, the OMOX composition

$$N_2H_5NO_3 + 1.5 ZrH_2 \rightarrow 1.5 ZrO_2 + 4H_2 + 1.5 N_2$$

is comprised of 40.467% $N_2H_5NO_3$ + 59.533% ZrH_2 , (by weight) and has heating

TABLE 3.1

(Continued)

EQUILIBRIUM YIELD VALUES AND HEATING VALUES @ P = 750 PSIA FOR SELECTED PROPELLANT INGREDIENTS

		librium Y 00 gms @ 2		Equilibrium Heating Values, K cal/100 gms, @ T, °K					
PROPELLANT COMPOSITION, WT %	H ₂	HC1	Σng	3000	2800	2600	2400		
	49				2 2				
0.995% N ₂ H ₅ NO ₃ + 59.005% Zr	1.066	_	1.732	-59.96	-79.64	-85.44	-90.99		
36.909% NH ₄ NO ₃ + 63.091% Zr	0.9124	<u>-</u> -1 _y ,	1.389	-61.49	-86.81	-86.81	-91.88		
39.174% NH ₄ C10 ₄ + 60.826% Zr	0.4972	0.3051	1.000	-89.06	-107.7	-112.2	-116.5		
3.581% Zr(OH) ₄ + 36.419% Zr	0.7938		0.8031	28.625	7.156		- 1.699		
2.886% WO ₃ + 37.114% Zr	0.0	-	0.0	-18.15	-28.64	-30.56	-32.45		
7.798% H ₂ WO ₄ + 42.202% Zr	0.2300	: <u> </u>	0.2326	-16.54	-28.87	-31.40	-33.87		
0.495% NC + 39.505% Zr	0.8233	· -	2.446	-25.32	-41.10	-47.42	-53.49		
8.308% TMETN + 41.692% Zr	1.011	6 × _	2.517	-36.96	-53.56	-60.15	-66.45		
9.529% Nitroguanidine + 30.471% Zr	1.315	- ·	3.346	8.407	- 6.956	-14.72	-22.09		
	9 8 2					10 B			
3.791% N ₂ H ₅ NO ₃ + 36.209% A1	1.657		2.700	-89.38	-101.6	-112.8	-123.4		
9.734% NH ₄ NO ₃ + 40.266% Z1	1.475	· - 3	2.252	-95.52	-107.2	-118.1	-128.5		
2.021% NH ₄ C10 ₄ + 37.979% A1	0.7689	-	1.593	-135.2	-145.7	-155.5	-164.7		
74.301% A1(OH) ₃ + 25.699% A1	1.415		1.440	51.79	40.20	29.30	18.84		
7.643% H ₂ WO ₄ + 22.357% A1	0.3077		0.3131	-20.47	-25.24	-29.82	-34.26		
81.573% Zr(OH) ₄ + 18.427% A1	1.014	-	1.032	38.76	20.63	.13.72	7.126		
•	21	8:			7	#S #8			
			· ·	*					

TABLE 3.1 (Continued) EQUILIBRIUM YIELD VALUES AND HEATING VALUES @ P $_{\rm c}$ = 750 PSIA FOR SELECTED PROPELLANT INGREDIENTS

		librium Yi 00 gms @ 2		Equilibrium Heating Values, K cal/100 gms, @ T, °K				
PROPELLANT COMPOSITION, WT %	H ₂	HC1	Σng	3000	2800	2600	2400	
69.316% N ₂ H ₅ NO ₃ + 30.684% CH ₂	3.955	-	7.305	163.8	147.8	133.2	119.5	
65.544% NH ₄ NO ₃ + 34.456% CH ₂	4.038	.	7.383	188.7	172.4	157.7	143.8	
57.680% NH ₄ C10 ₄ + 32.320% CH ₂	3.129	0.5697	6.353	125.5	111.4	98.54	86.61	
35.022% Zr(OH) ₄ + 14.978% CH ₂	2.114		3.214	167.9	147.5	139.1	131.1	
34.638% WO ₃ + 15.362% CH ₂	1.082	· -	2.196	108.4	102.9	97.92	93.21	
31.662% H ₂ WO ₄ + 18.338% CH ₂	1.616	. <u>.</u> .	2.950	135.2	128.1	121.6	115.5	
73.134% Zr0(NO ₃) ₂ ·2H ₂ 0 + 26.866% CH ₂	2.428		4.660	141.5	124.6	114.3		
57.112% Zr(NO ₃) ₄ ·5H ₂ O + 32.888% CH ₂	3.074	· _	5.804	154.0	136.4	124.2	1 (4.1)	
58.857% N ₂ H ₅ NO ₃ + 31.143% HTPB	3.414	_	6.796	138.0	123.1	109.6	96.92	
55.056% NH ₄ NO ₃ + 34.944% HTPB	3.430	_	6.812	159.4	144.6	131.0	118.3	
57.208% NH ₄ C10 ₄ + 32.792% HTPB	2.565	0.5654	5.824	98.54	85.75	73.96	63.03	
		л		30 30 at at	6 4 8 5 60 60		K100	
		¥			520			
		um ^{to}	3	5-		20		
		<i>13</i>			, .	2		

TABLE 3.1
(Continued)

EQUILIBRIUM YIELD VALUES AND HEATING VALUES @ P = 750 PSIA FOR SELECTED PROPELLANT INGREDIENTS

8 8		librium Y:		1 -		Heating Vagms, @ T,	
PROPELLANT COMPOSITION, WT %	H ₂	HC1	Σng	3000	2800	2600	2400
29.465% NH ₄ NO ₃ + 70.525% PEG 200	3.949	₆ –	7.299	183.5	167.4	152.8	139.1
31.493% NH ₄ C10 ₄ + 68.507% PEG 200	3.530	0.2652	6.823	154.2	139.2	125.4	112.6
52.223% NC + 47.777% PEG 200	2.894	-	6.306	135.7	122.0	109.5	97.77
49.957% TMETN + 50.043% PEG 200	3.153	-	6.550	133.3	119.1	106.1	93.84
61.959% Nitroguanidine + 38.041% PEG 200	2.907	-	6.319	132.9	119.2	106.7	94.95
	y 11.	*					
55.552% NH ₄ NO ₃ + 44.448% TCE	1.360	· – 20	4.863	63.21	53.10	43.68	34.73
57.910% NH ₄ ClO ₄ + 42.090% TCE	0.7264	0.4838	4.115	45.22	36.49	28.41	20.77
76.589% NC + 23.411% TCE	1.038	-	4.558	49.93	40.58	31.81	23.44
74.923% TMETN + 25.077% TCE	1.294	-	4.800	40.37	30.40	21.11	12.27
	7.0	. 18		20			
100% NH ₄ C1	2.791	1.849	5.641	227.6	214.4	202.7	192.0
100% N ₂ H ₅ C1	2.903	1.444	5.870	169.1	155.6	24	2
100% HDB	8.200	_	8.546	227.7	128.9	87.65	59.99
<i>€</i>		¥	-				
		P32.p					
				179 11			5 (8)

TABLE 3.1
(Continued)

EQUILIBRIUM YIELD VALUES AND HEATING VALUES @ P = 750 PSIA FOR SELECTED PROPELLANT INGREDIENTS

• 2		librium Y .00 gms @				Heating Va	
PROPELLANT COMPOSITION, WT %	H ₂	HC1	Σng	3000	2800	gms, @ T, 2600	2400
79.923% Zr(OH) ₄ + 20.077% A1H ₃	1.988	_ 1	2.023	64.16	42.60	34.52	26.07
57.154% NH ₄ NO ₃ + 42.846% A1H ₃	3.532	_	4.313	-37.62	-54.82	-70.16	-84.34
79.439% WO ₃ + 20.561% A1H ₃	1.018		1.036	4.110	- 2.143	- 7.844	-13.19
75.745% H ₂ WO ₄ + 24.255% AlH ₃	1.501	1.501 -		10.45	2.422	- 4.823	-11.56
61.303% N ₂ H ₅ NO ₃ + 38.697% A1H ₃	3.508	· .	4.544	-37.31	-54.44	-69.68	-83.75
59.489% NH ₄ ClO ₄ + 40.511% AlH ₃	2.727	0.4244	3.578	-77.42	-93.34	107.74	120.99
72.221% A1(OH) ₃ + 27.779% A1H ₃	2.750	·	2.799	85.18	70.05	56.26	43.34
						-	
			** ** **	* 2 * 2 * 2	* , *		

TABLE 3.2
GLOSSARY OF CANDIDATE PROPELLANT INGREDIENTS

		INGREDIENT	Chemical Formula or Acronym
Α.	OXID	DIZER	
	1.	Hydrazine Nitrate	N ₂ H ₅ NO ₃
	2.	Ammonium Nitrate	NH, NO 3
	3.	Ammonium Perchlorate	NH _A C10 _A
	4.	Tungsten Oxide	wo ₃
	5.	Tungstic Acid	H ₂ WO ₄
	6.	Zirconyl Nitrate Dihydrate	Zro(No ₃) ₂ ·2H ₂ o
	7.	Zirconium (IV) Nitrate Pentahydrate	Zr(NO ₃) ₄ ·5H ₂ 0
ů.	8.	Nitrocellulose (A Resin used in Solid Propellant Binders)	NC
•	9.	Trimethylolethane Trinitrate (A co- plasticizer for Nitrocellulose)	TMETN
D	ोल राज्य		
В.	FUEL		Zr
	1.	Zirconium	
	2. 3.	Zirconium Hydride Aluminum	ZrH ₂ Al
٥٠			_
	5.	Aluminum Hydride Tetracyanoethylene (A High-energy Organic Compound, C ₂ (CN) ₄)	AlH ₃ TCE
c.	COOL	ANTS	
	1.	Zirconium (IV) Hydroxide	Zr(OH)
	2.	Aluminum Hydroxide	A1(OH) ₃
	3.	Hydrazine Diborane (B ₂ N ₂ H ₁₀)	HDB
	4.	Polyethylene, Wax, or Long-chain Hydrocarbon	CH ₂

TABLE 3.2 (GLOSSARY OF CANDIDATE PROPELLANT INGREDIENTS - CONTINUED)

2 5	; 	Ingredient		*	Chemical Formula or Acronym
		6 7		E 1 48	# # \$ 8
	5.	Hydroxy Terminated Polybuta Isocyanate-cured Polyuretha System Widely used in the Polyurethan Industry)	ne Binder		нтрв
	6.	Polyethylene Glycol with a l Weight of 200 a.m.u.	Molecular		PEG-200
	7.	Nitroguanidine			
D.	CHL	ORIDES			•
	1.	Ammonium Chloride		ŧ	NH _A C1
	2.	Hydrazine Chloride		*	N ₂ H ₅ C1
	3.	Zirconium (IV) Chloride			zrcl ₄

 \tilde{v} alues, ΔH_{T}° = -19.98, -40.91, -47.98 and -54.66 Kcal/100 gm @ 3000, 2800, 2600 and 2400°K, respectively. These heating values indicate the heat available at the indicated temperature. Similarly, the OMOX composition

$$Zr(OH)_4 + ZrH_2 \rightarrow 2 ZrO_2 + 3 H_2$$

is comprised of 63.078% $\rm Zr(OH)_4 + 36.922\%$ $\rm ZrH_2$ (by weight) and has heating values of 52.70, 30.48, 25.17 and 20.17 Kcal/100 gm at 3000, 2800, 2600 and 2400°K, respectively, indicating that heat is required to attain these temperatures. Note that in the above two examples the discontinuity in $\Delta H_{\rm T}^{\circ}$ vs temperature between 2800°K and 3000°K. This is due to the heat of fusion of $\rm ZrO_2$, which takes place at 2950°K. Note that the tabulated values of $\rm H_2$, HCl and $\rm \Sigma ng$ are equilibrium values at 2800°K, not the stoichiometric values expressed above. The two do not differ much, especially at the lower temperatures and differ primarily due to $\rm H_2$ dissociation.

The results in Table 3.1 are discussed in terms of $\rm H_2$ yield, $\rm H_2$ yield purity ($\rm H_2/\Sigma ng$), heat release and oxygen content (of oxidizers). At the start of the program, $\rm H_2$ yield was held most important and the effort was directed primarily toward maximizing this parameter. For this reason organic propellant ingredients were initially favored. Later on in the program, it became apparent that contaminants such as CO greatly reduce lasing performance, and for this reason the effort was redirected toward maximizing $\rm H_2$ yield at values of $\rm H_2$ yield purity >95% and the use of organics was eliminated. Heat release and oxygen content of oxidizers are most important when high-H₂-content compounds such as HDB or AlH₃ are used in the propellants. Salient features of the results in Table 3.1 are discussed below.

3.2.1 Oxidizers

In terms of H₂ yield the ingredients that generate the most H₂ in descending order are N₂H₅NO₃ > NH₄NO₃ > NH₄ClO₄ > H₂WO₄ > WO₃. In addition to the above there are other oxidizers which are less effective than NH₄NO₃ and are valued in the order Al(NO₃)₃·9H₂O > TMETN > NC > NH₄ClO₄ > Zr(NO₃)₄·5H₂O > ZrO(NO₃)₂·2H₂O > H₂WO₄. In terms of gas purity of H₂ yield (H₂/ Σ ng), H₂WO₄ = WO₃ > NH₄ClO₄ > NH₄NO₃ > N₂H₅NO₃; in terms of heating value, NH₄ClO₄ > NH₄NO₃ > N₂H₅NO₃ > WO₃ > H₂WO₄, and in terms of oxygen content, NH₄NO₃ > NH₄ClO₄ > N₂H₅NO₃ > H₂WO₄, and in terms of oxygen content, NH₄NO₃ > NH₄ClO₄ > N₂H₅NO₃ > H₂WO₄ > WO₃.

All of these oxidizers are commercially available and relatively easy to use, although both non-metallic nitrates require specially-dried processing facilities (relative humidity <50% at 70°F) for use. Of the two nitrates, $N_2H_5NO_3$ is preferred because this compound does not undergo a low-temperature phase transition as does $\mathrm{NH_4NO_3}$. Both of the nitrates, as well as $\mathrm{NH_4ClO_4}$ yield substantial amounts of N₂ as a combustion product, which is a contaminant. This problem, however, is not severe with NH2C10, because typically only small amounts of this oxidizer are used in any given composition due to the constraint H2/HC1 = 97/3. Best (H_2 + HC1) purity values are obtained by use of either WO₃ or H_2 WO₄ with hydrogen-containing fuels, coolants and chlorides. Both of these oxidizers yield the metal W(c) (condensed phase) as a combustion product which can be filtered out or retained in the motor, to prevent contamination of the working fluid. Tungsten is thermodynamically stable at the temperatures of interest (2800°K) with respect to H2, HCl, and metal oxides, so that contamination of the working fluid with tungsten reaction products is not a problem. Other possible oxidizers include aluminum nitrate, A1(NO3)3.9H2O, zirconium nitrate, Zr(NO3)2. $5\text{H}_2\text{O}$ and zirconyl nitrate, $2\text{rO}(\text{NO}_3)_2 \cdot 2\text{N}_2\text{O}$. The best of these is Al-nitrate, yielding 2.4 mols $\rm H_2/100~gm$, which is seen to be inferior to both $\rm N_2H_5NO_3$ and NH, NO 2. The use of hydrates as propellant ingredients is not good practice. The organic oxidizers, NC and TMETN, received only minor attention here because contamination of the H_2 yield with CO and N_2 is significant. Calcium and magnesium nitrates were not considered because of the instability at 2800°K of CaO(c) and MgO(c), which are major combustion products of these oxidizers. Both oxides react significantly with H, and HCl at high temperatures to yield H₂O and other contaminants (Ca, Mg, CaOH, MgOH, CaCl, CaCl₂, MgCl₂). Similarly, the other alkaline earth nitrates, Sr(NO3)2 and Ba(NO3)2, which are well-known oxidizers used in colored flare signals, are not considered usable here. Alkali metal nitrates and perchlorates (e.g., $NaNO_3$, $CaClO_4$, KNO_3 , etc.) yield the chlorides quantitatively in combustion products. These species, although gaseous at stagnation or chamber conditions, would condense in the cavity and for this reason these oxidizers are not considered as propellant ingredient candidates.

3.2.2 Fuels

As mentioned previously, hydrocarbons are not desirable as fuels, because combustion to CO is not energetic enough to yield a flame temperature of $2800\,^{\circ}\text{K}$, and combustion to CO_2 yields H_2O due to the water-gas reaction. For the attainment of high temperature without H_2O contamination, metallic fuels must be employed.

In terms of $\rm H_2$ yield, $\rm AlH_3$ > $\rm ZrH_2$ > $\rm Al$ > $\rm Zr$; in terms of $\rm H_2$ yield purity, $\rm AlH_3$ > $\rm ZrH_2$ > $\rm Al}$ = $\rm Zr$, and in terms of heating value, $\rm Al$ > $\rm Zr$ > $\rm ZlH_3$ > $\rm ZrH_2$. Since $\rm AlH_3$ has stability and sensitivity problems, it is not a viable candidate. Consequently, the value of the fuels in propellant compositions is considered primarily as the source of energy rather than of $\rm H_2$. Either $\rm ZrH_2$ or $\rm Zr$ are favored in preference to $\rm Al$. This is because $\rm ZrO_2$, the major $\rm Zr$ combustion product at the OMOX point, is thermodynamically more stable than $\rm Al_2O_3$. Even though the use of $\rm Al$ compares favorably with $\rm Zr$ and $\rm ZrH_2$ in terms of $\rm H_2$ yield purity, $\rm Al_2O_3$ dissociates somewhat and reacts with $\rm H_2O$ as a combustion product. In $\rm Zr$ -containing propellants, the $\rm H_2/H_2O$ molar ratio is $\rm \sim 250$, but is only $\rm \sim 80$ in the combustion products of propellants containing Al. The additional effect of this phenomenon on the combustion products composition at throat conditions is discussed in Section 3.3.

For reasons stated previously, due to CaO and MgO stability, neither ${\rm CaH_2}$ nor MgH $_2$ were considered as fuels.

3.2.3 Coolants

The combustion of stoichiometric mixtures of candidate oxidizers and fuels yields values of flame temperature greater than 2800° K. In order to reduce the flame temperature to 2800° K, coolants are required and those of choice include Zr(OH)_4 , Al(OH)_3 , $(\text{CH}_2)_2$ and HDB. Stoichiometric mixtures of the hydroxides with candidate fuels serve as coolants to obtain a flame temperature of 2800° K, and yield H_2 as the only gaseous combustion product. The as-received form of Zr(OH)_4 is the hydrate $\text{Zr(OH)}_4 \cdot \text{nH}_2 \text{O}$ (n = 4 to 5) which must be dried carefully prior to use (48 hours at 145° F in a forced-draft oven has been found to be satisfactory). Care must be used because excessive drying at too high a temperature drives the as-received product to ZrO_2 .

Stoichiometric mixtures of (CH2)n and candidate oxidizers are coolants if the oxidation ratio (0_R) of the mixture, defined as the molar ratio of oxygen to carbon, is unity. At values of 0_R <1, soot and hydrocarbon (HC) fragments (CH, CH2, CH3, CH4, etc.) appear as combustion products, while at $0_p > 1$, CO_2 and H_2O appear as combustion products. Note that while the H_2 yield of (CH2) n is second only to HDB in terms of mols/100 gm, the molar purity of the H₂ yield, assuming C \rightarrow CO is only 50%, (CH₂ \rightarrow CO + H₂, H₂/(H₂ + CO) = 1/2). Because of the high degree of purity required, the use of hydrocarbons in candidate propellants was given only minor consideration. The purity effect can be minimized by using the hydrocarbon in propellants containing no oxygen for burning the carbon. However, soot and HC fragments would result. It is possible to filter out the soot, depending on the ability of the composition to form a clinker on combustion (if the composition does not clink, the soot will be generated as angstrom-sized particles and aggomerates thereof, which will probably be very difficult to remove). Due to this and the unknown effect that HC fragments would have on gain and available power, hydrocarbons were ruled out.

Of the three coolants, only HDB is not commercially available in large quantities. However, synthesis of this compound is straight forward and results in a product of good yield and high purity. It is the best $\rm H_2$ gas generator ingredient of all the candidate ingredients. It is capable of self-sustained combustion, with an adiabatic flame temperature of 1877°K ($\rm B_2N_2H_2 \rightarrow \rm 5H_2 + 2BN(c)$). The compound is not a good candidate for systems of flame temperature equal to 3000°K because of significant dissociation of BN(c) at this temperature, and subsequent recombination in the throat, resulting in a dirty dust-laden working fluid. However, BN(c) is essentially undissociated at 2800°K and the $\rm H_2$ yield is essentially cleam. HDB is a fairly reactive solid, and has exhibited chemical incompatibility with many propellant ingredients of interest. Storage and use is limited to very dry facilities (<40% relative humidity at 70°F). Its use with NH4ClO4 may result in explosions on ignition, though this seems to be a phenomenon dependent on HDB purity and the absolute and relative contents of each component in the propellant composition.

The oxidizers and the coolants are the primary H_2 sources in the propellants. The requirement (see Section 3.1)that (H_2 + HCl) yield be >5% by weight @ (H_2 /HCl) = 97/3 (molar) dictates that the H_2 yield values as tabulated in Table 3.1) be >1.59 mol/100 gm of propellant. Disregarding the organics and AlH₃, it is seen that the only oxidizer-fuel combinations that

qualify are $N_2H_5NO_3$ + either ZrH_2 or A1. The only coolant that qualifies is HDB. Of the two H_2 sources, HDB is obviously superior, not only in terms of H_2 yield, but also H_2 gas purity.

3.2.4 Chlorides

The chlorides are included in propellants as C1-donors to yield HC1 in the combustion products. All are commercially available and easily used and handled. The H-containing compounds also yield N_2 as a combustion product, but this is not too important because the chloride content of propellants is not high due to the requirement for a high value of $H_2/HC1$ (97/3). For compositions in which no contaminants are permitted in the combustion products, $2rCl_4$ can be used as the C1-donor. For example, the equimolar mixture of $2rCl_4$ and $2r(OH)_4$ yields HC1 and $2rQ_2$ (c) as the only combustion products.

In terms of H₂ yield, N₂H₅Cl > NH₄Cl > (NH₄ClO₄ + ZrH₂) > (NH₄ClO₄ + Al) > (NH₄ClO₄ + Zr). In terms of HCl yield, NH₄Cl > N₂H₅Cl > (NH₄ClO₄ + Al) > (NH₄ClO₄ + Zr) = (NH₄ClO₄ + ZrH₂). Values of (H₂ + HCl) yield purity range from 74% for N₂H₅Cl up to 88% for NH₄ClO₄ + ZrH₂. The NH₄ClO₄-fuel combinations are of course more energetic than either of the chlorides.

3.3 <u>Calculation of Theoretical Yield</u>

3.3.1 Procedure

The selection of candidate propellant compositions is based primarily on computed values of theoretical flame temperature and desired combustion product composition. Basically the propellant can be constructed from any combination of oxidizer, fuel, coolant and chloride from the binary combinations in Table 3.1 by means of two equations for a heat and molar ratio balance; these equations are:

$$\Gamma(X\Delta H_{2800}^{\circ})_{i} = 0$$
 Heat Balance (1)

$$\frac{\Sigma(XH_2)_i}{\Sigma(XHC1)_i} = \frac{H_2}{HC1}$$
 Molar Ratio Balance (2)

where X_1 are multipliers. In (1) the condition is that combustion is adiabatic, while in (2), H_2/HCl = 97/3 generally. In all cases, there are three multipliers (e.g., one for oxidizer-fuel, one for fuel-coolant, and one for chloride) any one of which is arbitrary and can be set to unity. Simultaneous solution of (1) and (2) for the other two multipliers and normalization to 100 result in propellant compositions that yield the flame temperature and H_2 and HCl very close to the desired values. Equilibrium calculations of resulting compositions yield slightly different values than predicted by using (1) and (2) because of mixing effects on entropy and the resultant equilibria. If required, compositions are then sometimes varied slightly to attain the desired results. For example, consider the system $N_2H_5NO_3 + Zr H_2 + Zr(OH)_4 + NH_4ClO_4$. What is the OMOX composition with a flame temperature of $2800^{\circ}K$ and $H_2/HCl = 97/3$? Consider the three binaries:

A.
$$40.467\% \text{ N}_2\text{H}_5\text{NO}_3 + 59.533\% \text{ Zr H}_2, \Delta\text{H}_2^\circ = -40.91, \\ \text{H}_2 = 1.685, \text{HCl} = 0$$

B.
$$63.078\% \text{ Zr}(OH)_4 + 36.922\% \text{ Zr H}_2, \Delta H_{2800}^{\circ} = 30.48,$$

 $H_2 = 1.181, HC1 = 0$

C.
$$38.654\% \text{ NH}_4\text{ClO}_4 + 61.346\% \text{ ZrH}_2, \Delta \text{H}_2^\circ = -67.44, \\ \text{H}_2 = 1.144, \text{ HCl} = .3048$$

Let x_A , x_B and x_C be the multipliers for compositions A, B and C, respectively, and set $x_A = 1$.

From (1)

$$-40.91 + 30.48X_B - 67.44X_C = 0$$

From (2)

$$\frac{1.685 + 1.181X_B + 1.144X_C}{0.3048X_C} = \frac{97}{3}$$

From which we determine

$$X_{R} = 2.5287$$

$$X_C = 0.53625$$

The resultant composition is comprised of $100\mathrm{X}_{\mathrm{A}}$ = 100 gm of Composition A plus $100\mathrm{X}_{\mathrm{B}}$ = 252.87 gm of Composition B and $100\mathrm{X}_{\mathrm{C}}$ = 53.625 gm of Composition C, and the weight percentages of each ingredient are obtained by normalizing to the sum X_{A} + X_{B} + X_{C} = 4.0650:

$$N_2H_5NO_3 = 40.467X_A/(X_A + X_B + X_C)$$

= 40.467/4.0650
= 9.955

$$Z_{\text{TH}_2}$$
 = $(59.533X_{\text{A}} + 36.922X_{\text{B}} + 61.346X_{\text{C}})/(X_{\text{A}} + X_{\text{B}} + X_{\text{C}})$
= $185.79 / 4.065$
= 45.706

$$Z_r(OH)_4 = 63.078X_B/(X_A + X_B + X_C)$$

= 159.50/4.065

= 39.239

$$NH_4C10_4 = 38.654X_C/(X_A + X_B + X_C)$$
= 20.728/4.065

= 5.100

Slight variations in the composition to compensate for mixing effects, as above, are generally derived by solving for NH_4ClO_4 content ([NH_4ClO_4]) around the above value. Thus, replace (2) with

$$\frac{38.654X_{C}}{X_{A} + X_{B} + X_{C}} = [NH_{4}C10_{4}]$$

and solve for selected values of $[NH_4ClO_4]$, say $[NH_4ClO_4] = 3$, 4, 5, 6 and 7%.

Calculations of theoretical flame temperature and combustion product composition were carried out for adiabatic combustion at 750 psi followed by isentropic expansion of the gaseous combustion products to throat conditions. The condensed species $(Al_2O_3, ZrO_2, W, BN, etc.)$ in these calculations were assumed to be retained in the combustion chamber and not allowed to undergo the expansion process with the gases. Tabulated values for H_2 yield, HC1 yield, and other gaseous species are throat values.

The results of these calculations and the finalized compositions, for $\rm H_2/HC1 \stackrel{\sim}{\sim} 97/3$ are summarized in Tables 3.3 to 3.10. Results are discussed below.

3.3.2 <u>Discussion of Results</u>

The results of the thermochemistry calculations are summarized by oxidizer for compositions containing $\mathrm{NH_4NO_3}$ (Tables 3.3 and 3.4), $\mathrm{N_2H_5NO_3}$ (Tables 3.5 and 3.6), $\mathrm{NH_4C1O_4}$ (Table 3.7), $\mathrm{H_2WO_4}$ (Table 3.8), $\mathrm{WO_3}$ (Table 3.9) and $\mathrm{ZrO(NO_3)_2 \cdot 2H_2O}$ (Table 3.10). Some calculations were carried out for temperatures other than 2800°K and these also are presented in the Tables, but without discussion.

3.3.2.1 Compositions Containing NH, NO

The results for compositions containing either Zr or ${\rm ZrH_2}$ are shown in Table 3.3 whereas results for either Al or ${\rm AlH_3}$ are shown in Table 3.4. All of the compositions containing organic coolants (${\rm CH_2}$ or HTPB) easily meet the (${\rm H_2}$ + HCl) weight yield requirement and are easily reduced to practice, but are very poor from the standpoint of (${\rm H_2}$ + HCl) gas purity. The composition that contains off-the-shelf ingredients and that comes closest to meeting the yield requirement (>90%) is 3.3.4, which is comprised of ${\rm NH_4NO_3}$, ${\rm ZrH_2}$, ${\rm Zr(OH)_4}$, and ${\rm NH_4ClO_4}$. This composition was selected as one of the candidates and is considered to be representative of the best that can be done, in terms of yield and purity, with off-the-shelf ingredients. Notice though that the use of HDB as a coolant (Compositions 3.3.21 and 3.3.22) greatly increases both yield and gas purity.

One of the most serious drawbacks to the use of Al or ${\rm AlH}_3$ as a fuel is the recombination of reaction products in the throat. As noted previously, ${\rm Al}_2{\rm O}_3$ is less stable than ${\rm ZrO}_2$ at the desired stagnation conditions which result in the reactions

$$A1_20_3$$
(c) + $2H_2$ + $2HC1$ + $2A1C1$ + $3H_20$ (1)

$$A1_2O_3(c) + H_2 + 4HC1 + 2A1C1_2 + 3H_2O$$
 (2)

In the combustion chamber. The reverse of these reactions occurs at the throat. Although only gaseous species are expanded to throat conditions, recombination of AlCl or AlCl $_2$ with $\rm H}_2{\rm O}$ results in small amounts of $\rm Al}_2{\rm O}_3(c)$ formation at throat conditions. The yield of $\rm Al}_2{\rm O}_3$ through recombination is 1-2% by weight of the working fluid. Since the nozzle does not allow time for equilibrium to be obtained, these are strictly thermodynamic results and the amount of $\rm Al}_2{\rm O}_3$ that actually forms and the extent of interference this phenomenon may have on nozzle and/or laser performance is not known. A similar phenomenon is found for HDB-containing propellants due to recombination of the decomposition and reaction products of BN(c). However, this compound can be used in a two-compartment motor system, and the BN can be filtered prior to admixture of the $\rm H}_2$ with the other combustion products. In this way, the recombination phenomenon is obviated.

Comp		P	ropellant	Compos	ition, Wt	%	3		T _f , °K,*	H ₂ /HC1*	Yie (H ₂ +1	≘1d H C1) *
ID	Oxid	izer	Fue	1	Coo	lant	Ch1	loride	<u>@750 psia</u>	(molar)	Wt %	Vol %
3.3.1 3.3.2	29.144 17.950	NH ₄ NO ₃	57.325 49.203	ZrH ₂	10.948 30.482	Zr(OH) ₄	2.583 2.365	NH ₄ C1	3000 2799	32.1 32.2	4.85 4.46	79.4 85.1
3.3.3 3.3.4	20.241 10.482	NH ₄ NO ₃	54.914 47.132	ZrH ₂	19.845 38.386	Zr(OH) ₄	5 4	NH ₄ C10 ₄	3000 2801	33.1 38.2	4.35 3.84	83.0 89.5
3.3.5	4.008		42.838		48.440		4.714		2400	30.7	3.93	94.6
3.3.6 3.3.7 3.3.8	36.955 39.538 40.574	NH ₄ NO ₃	58.194 52.007 49.532	ZrH ₂	1.917 5.137 6.426	(CH ₂) _n	2.934 3.318 3.468	NH ₄ C1	3000 2801 2602	32.2 32.3 32.3	5.51 6.27 6.57	73.5 69.4 68.2
3.3.9 3.3.10 3.3.11	34.015 35.990 37.162	NH ₄ NO ₃	56.126 49.702 47.145	ZrH ₂	3.859 7.308 8.693	(CH ₂) _n	6 7 7	NH ₄ C10 ₄	3000 2802 2602	3 6. 3 34.2 35.9	6.45 6.26 6.47	70.9 67.3 66.2
3.3.12 3.3.13 3.3.14 3.3.15	33.918 36.803 38.122 39.461	NH ₄ NO ₃	54.736 47.670 44.740 41.768	ZrH ₂	4.655 8.527 10.137 11.771	НТРВ	6.691 7 7 7	NH ₄ C10 ₄	3000 2802 2602 2402	31.0 33.1 34.5 35.9	5.58 6.11 6.30 6.47	69.0 64.4 63.0 61.5
3.3.16 3.3.17 3.3.18 3.3.19	39.628 41.904 42.118 43.208	NH ₄ NO ₃	43.982 38.876 36.226 33.784	Zr	10.394 13.220 14.656 16.008	HTPB	6 6 7 7	NH ₄ C10 ₄	3000 2801 2601 2401	31.7 35.6 32.1 33.7	5.07 5.51 6.02 6.21	56.1 55.2 54.9 54.5
3.3.20	23.507	NH ₄ NO ₃	40.182	Zr	30.311	HDB	6	NH ₄ C1	2825	34.0	10.1	89.3
3.3.21	2.482	NH ₄ NO ₃	33.744	Zr	44.774	HDB	19	NH ₄ ClO ₄	2812	32.0	12.7	95.3
3.3.22	7.782	NH ₄ NO ₃	34.896	Zr	40.322 + 1	HDB (CH ₂) _n	16	NH ₄ C10 ₄	2818	34.1	11.8	92.6

^{*} Notes: T_f is adiabatic flame temperature at a combustion pressure of 750 psia. $H_2/HC1$ and yield (H_2+HC1) refer to combustion product values at throat conditions. Yield (H_2+HC1), vol. % is the yield purity value.

Table 3.4 Theoretical Values of Flame Temperature and Yield for Selected OMOK Compositions: Compositions Containing $^{\rm NH}_4{^{\rm NO}}_3$ and Either Al or AlH $_3$.

Comp		Pr	onellant	Compos	ition, Wt		3		T _f , °K,	H ₂ /HC1	Yid (H ₂ +1	eld HCl)
ID	Oxidi		Fuel		Cool		Ch1	oride	@750 psia	(molar)	Wt %	Vol %
3.4.1 3.4.2	23.140 18.016	NH ₄ NO ₃	30.571 29.462	Al	43.289 50.069	A1(OH) ₃	3 2.453	NH ₄ C1	2989 2791	33.8 37.8	4.56 4.39	80.6 84.8
3.4.3 3.4.4	13.647 10.212	NH4 ^{NO} 3	30.219 29.179	Al 🗆	50.119 55.609	A1(OH) ₃	6.015 5	NH ₄ C10 ₄	2995 2794	34.4 38.2	4.19 4.10	85.3 89.2
3.4.5 3.4.6 3.4.7	59.033 58.973 59.442	NH ₄ NO ₃	26.953 25.097 23.294	A1	10.014 11.430 13.082	(CH ₂) _n	4 4.5 4.182	NH ₄ C1	2971 2776 2584	39.4 33.0 34.6	6.67 7.48 7.73	59.9 60.2 60.0
3.4.8 3.4.9 3.4.10	52.190 52.463 53.005	NH ₄ NO ₃	25.431 23.541 21.791	A1	12.379 13.996 15.502	(CH ₂) _n	10 10 9.702	NH ₄ C10 ₄	2970 2777 2585	34.4 32.9 33.0	7.02 7.54 7.92	58.6 59.0 59.0
3.4.11 3.4.12 3.4.13 3.4.14	52.234 53.454 53.707 53.952	NH ₄ NO ₃	23.772 21.870 19.947 18.100	A1	13.994 15.676 17.346 18.948	НТРВ	10 9 9	NH ₄ C10 ₄	2970 2778 2585 2393	32.2 34.3 33.2 32.6	6.76 6.97 7.36 7.68	55.5 55.5 55.2 54.8
3.4.15	36.026	NH ₄ NO ₃	35.610	A1H ₃	22.364	A1(OH)3	6	NH _A C1	2784	33.9	10.1	85.7
3.4.16	39.289	NH ₄ NO ₃	35.162	A1H ₃	18.549 + 1	A1 (OH) 3	6	NH ₄ C1	2783	34.5	10.2	83.6

3.3.2.2 <u>Compositions Containing N₂H₅NO</u>₃

These results are shown in Table 3.5 for compositions containing either Zr or ${\rm ZrH}_2$ as a fuel and in Table 3.6 for Al and ${\rm AlH}_3$ as fuels. Compared to ${\rm NH}_4{\rm NO}_3$, the use of this oxidizer in comparable formulations results in slight increases in yield and decreases in gas purity. For example, compare 3.3.2 and 3.5.2, at 2800°K the composition containing NH₄ has a weight yield of 4.46% and a purity of 85.1%, whereas the composition containing N₂H₅NO₃ has a weight yield of 4.64% and a purity of 81.4%.

3.3.2.3 Compositions Containing Either H₂WO₄ or WO₃

These oxidizers are ideal for use in thermite systems, resulting in excellent clinkers. Both result in high values of ($\rm H_2$ + HCl) gas purity since no N₂ is introduced in the combustion products. When used in propellants made with off-the-shelf ingredients, H₂WO₄ is preferred. The composition 3.8.1, comprised of H₂WO₄,ZrH₂, Zr, and NH₄ClO₄ is representative of the best that can be attained from thermite systems made with off-the-shelf materials and was selected as one of the propellant candidates. Note that the yield of this composition, 2.59%, is very low; however, it is also very pure. When used with HDB, there is not much advantage between the two oxidizers and, in fact, the performance is comparable to the use of NH₄NO₃ or N₂H₅NO₃ (>10% weight yield). The composition selected as one of the candidates was composition 3.8.8.

Table 3.5 Theoretical Values of Flame Temperature and Yield for Selected OMOX Compositions: Compositions Containing $^{N}2^{H}5^{NO}3$ and Either Zr or $^{ZrH}2$

				**		25	3 4.14 2.	_	Yield				
Comp		Pro	pellant (Composi	tion, Wt	%			T _f , °K,	H ₂ /HC1	(H ₂ -	+HC1)	
ID	Oxid:	izer	Fue	<u>l</u>	Coo	lant	Ch	loride	@750 psia	(molar)	Wt %	Vo1_%	
3.5.1	31.997	N2H5NO3	53.994	ZrH ₂	11.484	Zr(OH)	2.725	NH _A C1	3000	32.1	5.11	75.1	
3.5.2	19.832		47.097	<u> </u>	30.615	.,	2.456	•	2800	32.2	4.64	81.4	
3.5.3	31.621	N2H5NO3	53.292	$^{\mathrm{ZrH}}_{2}$	11.571	Zr(OH) ₄	3.516	N2H5C1	3000	32.1	5.15	74.6	
3.5.4	22.126	N2H5NO3	52.444	ZrH ₂	20.430	Zr(OH) ₄	5	NH ₄ ClO ₄	3000	34.4	4.46	79.2	
3.5.5	41.560	$^{\mathrm{N}_{2}\mathrm{H}_{5}\mathrm{NO}_{3}}$	52.696	ZrH ₂	2.541	(CH ₂) _n	3.203	NH ₄ C1	2950	32.2	6.02	68.9	
3.5.6	43.917	N ₂ H ₅ No ₃	47.368	-	5.188		3.527		2801	32.3	6.66	66.3	
3.5.7	40.759	N2H5NO3	53.131	ZrH ₂	2.056	(CH ₂) _n	4.054	N2H5C1	3000	32.2	5.94	68.9	
3.5.8	37.070	N2H5NO3	51.781	z_{rH_2}	4.164	$(CH_2)_n$	6.985	NH4C104	3000	32.2	5.97	67.5	
3.5.9	35.114	N ₂ H ₅ NO ₃	51.658	$^{ZrH}_{2}$	9.290	HDB	3.938	NH ₄ C1	2999	32.1	7.38	79.4	
3.5.10	25.758	$^{\mathrm{N}}2^{\mathrm{H}}5^{\mathrm{NO}}3$	37.074	Zr	29.168	HDB	8	N2H5C1	2826	32.7	10.4	85.7	
3.5.11	3.969	N2H5NO3	33.661	Zr	44.370	HDB	18	NH4C104	2817	33.2	12.6	94.8	
3.5.12	23.808	N ₂ H ₅ NO ₃	34.267	Zr	31.825 +4.100	HDB Zr(OH) ₄	6 2 2	ZrC1 ₄	2831	36.8	9.84	88.7	
						4							

Table 3.6 Theoretical Values of Flame Temperature and Yield for Selected OMOX Compositions:

Compositions Containing $\mathrm{N_2H_5NO_3}$ and Either Al or $\mathrm{A1H_3}$

			Cor	npositi	ons Contai	ning N2H5	NO_3 and 1	Either Al o	r AlH ₃		Yi	.eld
Comp		Pro	pellant	Composi	tion, Wt	%			T _f , °K,	H;/HC1	(H ₂ +	HC1)
ID	0xid		Fue			lant	Ch	loride	@750 psia	(molar)	Wt %	Vo1 %
3.6.1	25.786	N ₂ H ₅ NO ₃	29.177	A1	42.037	A1(OH) ₃	3	NH ₄ C1	2989	35.8	4.70	76.6
3.6.2	20.438	. 2 3 3	28.294		48.268	J	3	22	2789 ⁻	32.0	4.86	80.8
3.6.3	25.523	N2H5NO3	28.876	Al	41.601	A1(OH) ₃	4	N_2H_5C1	2988	34.6	4.81	75.8
3.6.4	14.546	N ₂ H ₅ NO ₃	29.383	A1	49.571	A1(OH) ₃	6.5	NH ₄ C10 ₄	2995	32.5	4.39	82.6
3.6.5	62.206	N ₂ H ₅ NO ₃	23.884	A1	8.910	(CH ₂) _n	5	NH _A C1	2969	32.4	7.47	58.4
3.6.6	63.059	3 3	21.867		10.861		4.213		2777	37.4	7.57	58.4
3.6.7	61.228	N ₂ H ₅ NO ₃	23.477	A1 .	8.796	(CH ₂) _n	6.5	N_2H_5C1	2968	32.1	7.54	58.3
3.6.8	54.611	N ₂ H ₅ NO ₃	22.537	A1	11.852	(CH ₂) _n	11	NH ₄ C10 ₄	2967	32.5	7.49	57.3
3.6.9	36.180	N ₂ H ₅ NO ₃		A1	34.784	HDB	8.5	NH _A C1	2950	32.0	12.8	84.5
3.6.10	47.622	$N_2H_5NO_3$	34.316	AlH ₂	11.062	A1(OH) ₃	7	NH _A C1	2982	32.6	10.44	78.7
3.6.11	37.348	N ₂ H ₅ NO ₃	32.764	A1H ₃	23.888	A1(OH)3	6	NH ₄ C1	2746	33.1	10.1	82.8
3.6.12	24.568	N ₂ H ₅ NO ₃	30.736	AlH ₃	39.588	A1(OH) ₃	5.108	NH ₄ C1	2391	34.0	9.52	87.7
3.6.13	57.759	N ₂ H ₅ NO ₃	32.379	AlH ₃	2.862	A1 (OH) 3	7	NH ₄ C1	2978	34.2	10.8	73.5
3.6.14	59.143	N ₂ H ₅ NO ₃	29.044	A1H ₃	5.813	(CH ₂) _n	6	NH ₄ C1	2780	37.2	10.7	71.5
3.6.15	18.264	N ₂ H ₅ NO ₃	33.600	AlH ₃	36.136	A1(OH) ₃	12	NH ₄ C10 ₄	2794	34.3	9.20	88.8
3.6.16	6.248	N ₂ H ₅ NO ₃	31.246	AlH ₃	51.506	A1(OH) ₃	11	NH _Z C10 _Z	2395	32.2	8.99	9,5.0
3.6.17	32.337	N ₂ H ₅ NO ₃	28.895	AlH ₃	33.768	Zr(OH)	5 .	NH _A C1	2784	35.6	8.63	82.9
3.6.18	12.474	N ₂ H ₅ NO ₃		A1H ₃	50.226	Zr(OH) ₄		NH ₄ C10 ₄	2794	33.5	7.56	90.0

Table 3.7 Theoretical Values of Flame Temperature and Yield for Selected OMOX Compositions: Miscellaneous Compositions Containing NH $_4$ ClO $_4$

Comp		Pro	pellant (Composi		4 T _f , °K,	H ₂ /HC1	Yie (H ₂ +1	eld HC1)			
ID	Oxidi	zer	Fue]		Cool	Lant	Ch	loride	@750 psia	(molar)	Wt %	Vo1 %
3.7.1	42.794	NH ₄ C10 ₄	26.206	A1	31	HDB			2798	42.4	9.28	76.5
3.7.2	16.725	NH ₄ C10 ₄	9.717	A1	70.161	HDB	3.397	ZrC1 ₄	2347	34.4	18.6	98.2
3.7.3	13.317	NH ₄ C10 ₄	20.677	Zr	66.006	HDB			2394	50.6	15.6	98.8
3.7.4	13.586	NH ₄ C10 ₄	19.921	Zr	63.493	HDB .	3	ZrC1 ₄	2393	32.9	17.0	98.7
3.7.5	3.095	NH ₄ C10 ₄	37.850	Zr	0.205	(CH ₂) _n			2800	30.2	2.55	96.2
3.7.6	3.000	NH ₄ C10 ₄	35.728 +2.596	Zr ZrH ₂	58.676	Zr(OH) ₄			2800	31.6	2.54	98.0
3.7.7	16.714	NH ₄ C10 ₄	26.526	ZrH ₂	56.760	HDB		7%	2384	37.1	15.7	98.4

Table 3.8 Theoretical Values of Flame Temperature and Yield for Selected OMOX Compositions: Compositions Containing ${\rm H_2WO_4}$ and Either Zr or ZrH₂

F-350 to	0xidizer		Propellant Composition, Wt %				T _f , °K,			н,/нс1	Yield (H ₂ +HCl)	
Comp ID			Fuel Fuel		Coolant		Chloride		@750 psia	(molar)	Wt %	Vol %
3.8.1	40.879	H ₂ WO ₄	42.968	ZrH ₂	13.153	Zr(OH)	3	NH ₄ C10 ₄	2799	32.6	2.59	98.0
3.8.2	20.641	H ₂ WO ₄	41.453	ZrH ₂	34.051	Zr(OH)		NH _Z C10 _Z	2400	30.8	3.21	98.3
3.8.3	54.657	H ₂ WO ₄	42.018	ZrH ₂	0.825	(CH ₂) _n	2.5	NH ₄ C10 ₄	2800	35.9	2.30	91.4
3.8.4	57.018	H ₂ WO ₄	41.649	ZrH ₂	0.127	(CH ₂) _n	1.206	ZrC1	2798	33.5	2.09	98.2
3.8.5	55.282	H ₂ WO ₄	41.856	ZrH ₂	0.615	$(CH_2)_n$	0.465 +1.782	ZrCl ₄ NH ₄ ClO ₄	2799	31.9	2.31	93.1
3.8.6	32.348	H ₂ WO ₄	30.354 +9.827	ZrH ₂ Zr	11.363 +9.292	HDB Zr(OH) ₄	6.816	NH ₄ C10 ₄	2780	34.6	4.92	96.4
3.8.7	47.996	H ₂ WO ₄	35.044	Zr	14.533	HDB	2.427	NH,C1	2802	41.5	4.25	96.5
3.8.8	11.821	H ₂ WO ₄	33.475	Zr	38.704	HDB	16	NH ₄ C10 ₄	2808	33.3	10.8	95.9
3.8.9	21.952	H ₂ WO ₄	33,623	Zr	30 +1	HDB (CH ₂) _n	13.425	NH ₄ C10 ₄	2802	32.2	8.91	93.7
3.8.10	48.988	H ₂ WO ₄	7.096 +33.816	Zr ZrH ₂	6.600 +1	Zr(OH) ₄ (CH ₂) _n	2.5	NH ₄ C10 ₄	2800	35.2	2.26	89.9

Table 3.9 Theoretical Values of Flame Temperature and Yield for Selected OMOX Compositions: Compositions Containing WO $_3$ and Either Zr or ZrH $_2$

Comp			Propellant (Composi	tion. Wt	%			T, °K,	H ₂ /HC1	Yi (H ₂ +	eld HC1)
ID	Oxidi		Fuel	_	Cool		Chl	loride	@75 0 psi a	(molar)	Wt %	Vo1 %
3.9.1	45 .1 81	WO ₃	38.685	ZrH ₂	14.124	Zr(OH)	2	NH ₄ C10 ₄	2789	36.4	1.86	98.2
3.9.2	60.983	WO ₃	36.097	ZrH2	.1.084	(CH ₂) _n	1.836	NH ₄ C10 ₄	2800	31.4	1.55	85.1
3.9.3	62.957	WO ₃	35.697	ZrH ₂	0.596	(CH ₂) _n	0.750	ZrC14	2799	33.0	1.30	90.6
3.9.4	62.143	WO ₃	11.593	Zr	25.637	ZrH ₂	0.627	N2H5C1	3000	32.2	0.916	95.9
3.9.5	60.239	WO ₃	7.178	Zr	31.199	ZrH ₂	1.384	NH ₄ C10 ₄	3000	30.2	1.13	97.2
3.9.6	9.960	WO ₃	28.959	Zr	49.327 +9.254	Zr(OH) 4 ZrH ₂	2.5	NH ₄ C10 ₄	2800	35.4	2.28	98.1
3.9.7	53.819	WO ₃	31. 763	Zr	12 845	HDB	1.846	NH _A C1	2998	32.2	3.45	97.4
3.9.8	20.113	WO ₃	32.056	Zr	34.831	HDB	13	NH ₄ C10 ₄	2811	36.2	9.27	96.0
3.9.9	17.840	wo ₃	10.529 +20.032	Zr ZrH ₂	37.889	HDB	13	NH ₄ C10 ₄	2600	34.2	11.0	98.1

Table 3.10 Theoretical Values of Flame Temperature and Yield for Selected OMOX Compositions: Compositions Containing $Zr0(NO_3)_2\cdot 2H_20$ and Either Zr or ZrH_2

Comp	P	Composit: ropellant			5 2	·2H ₂ O and	d Either Zr	or ZrH ₂ T _f , °K,	H ₂ /HC1		leld -HC1)
ID	0xidizer	Fuel			lant	Ch	loride	<u>@750 psia</u>	(molar)	Wt %	Vol %
3.10.1	34.593 ZrO(NO ₃) ₂ ·2H ₂ O	50.168	ZrH ₂	13.549	Zr(OH) ₄	1.690	NH ₄ C1	3000	32.1	3.17	86.7
3.10.2	21.348 Zro(No ₃) ₂ ·2H ₂ o	44.812	ZrH ₂ .	32, 024	Zr(OH) ₄	1.816	NH ₄ C1	2799	32.2	3.43	91.5
3.10.3	26.311 Zro(NO ₃) ₂ .2H ₂ O	49.998	ZrH ₂	19.691	Zr (OH) 4	4	NH ₄ C10 ₄	3000	30.3	3.29	89.1
3.10.4	12.840 Zro(No ₃) ₂ ·2H ₂ o	44.598	ZrH ₂	38.562	Zr(OH) ₄		NH ₄ C10 ₄	2800	32.6	3.46	94.0
3.10.5	46.846 ZrO(NO ₃) ₂ ·2H ₂ O	48.784	ZrH ₂	2.532	(CH ₂) _n	1.838	NH ₄ C1	3000	32.2	3.45	74.4
3.10.6	50.069 Zro(NO ₃) ₂ ·2H ₂ o	42.045	ZrH ₂	5.743	(CH ₂) _n	2.143	NH ₄ C1	2802	32.3	4.05	68.0
3.10.7	44.307 Zro(No ₃) ₂ ·2H ₂ o	47.924	ZrH ₂	3.769	(CH ₂) _n	4	NH ₄ C10 ₄	3000	34.2	3.56	71.8
3.10.8	46.725 Zr0(N0 ₃) ₂ ·2H ₂ 0	41.084	ZrH ₂	7.191	(CH ₂) _n	5	NH ₄ C10 ₄	2802	32.1	4.28	66.1

3.4 Properties of Candidate Propellants

3.4.1 General

The three candidate compositions chosen for evaluation in motor firings are Compositions 3.3.4, 3.8.1 and 3.8.8. The first two selections are about the best that can be obtained with off-the-shelf ingredients with regard to yield and purity, while the last is a composition containing HDB. As indicated in the tables of the previous section, there are other propellant selections besides 3.8.8 that have a high yield; however, only one high-yield highrisk system was selected. This selection was based on the system that appeared to have the highest probability of success. The Composition 3.3.4 is comprised of $\mathrm{NH_4NO_3}$, $\mathrm{ZrH_2}$, $\mathrm{Zr(OH)_4}$, and $\mathrm{NH_4C1O_4}$; Composition 3.8.1 is comprised of $\mathrm{H_2WO_4}$, ZrH_2 , $Zr(OH)_4$, and NH_4C1O_4 , and Composition 3.8.8 is comprised of H_2WO_4 , Zr, HDB and NH ClO. The safety and ballistic properties of each composition were measured, and are given in Table 3.11. Detailed theoretical results for flame and throat temperatures for the combustion product composition at a chamber pressure of 750 psia are given for each of the compositions in Table 3.12. Burning rate curves of rate vs pressure are shown for compositions 3.3.4 and 3.8.1 in Figures 3.2 and 3.3, respectively.

All of the compositions are prepared by first blending the ingredients and then pressing them to the desired size and shape at 30 KSI. For each composition, two values of the nozzle discharge coefficient are given, one for the gas-only flow condition (which is the design condition) and the other for the flow of both gases and solids (given in parenthesis). Burning rates are for specimens ~0.5" dia x ~1"long pressed at 30 KSI in steel tubes. Burning rates derived from subscale motor firings (see Section 3.5) are also shown. Results are discussed below.

3.4.2 Properties of Composition 3.3.4

This composition is the best (in terms of gas purity) of the nitrate-containing systems comprised of off-the-shelf ingredients. The composition presses up fairly well at 30 KSI, ignites easily and burns uniformly in the pressure range 400-1500 psia. Combustion at 750 psia yields combustion products comprised of 8.04% by weight gases and 91.96% solids. At this pressure, and at higher values, the solids form a good clinker. At low pressure (<100 psia) the condensed combustion products do not adhere and are entrained by the gases.

Table 3.11
PROPERTIES OF CANDIDATE PROPELLANTS

المراجعة المراجعة	Property	Comp	3.3.4	Comp	3.8.1	Comp	3.8.8	
A .	Composition, Wt%							
	Oxidizer	10.482	NH _A NO ₃	40.879	H ₂ WO ₄	11.821	H ₂ WO ₄	
	Fuel	47.132	ZrH ₂	42.968	ZrH ₂	33.475	Zr	
٠.	Coolant	38.186	Zr (OH)	13.153	Zr(OH)	38.704	HDB	
	Chloride	4.000	NH ₄ C10 ₄	3.000	NH ₄ C10 ₄	16.000	NH ₄ C10 ₄	
В.	Theoreticals							
	Flame Tempera- ture, °K	2	799	2	799	2	80 8	
	(H2/HC1) molar	38	3.7	3:	32.6		33.3	
	Yield (H2+HC1)							
	Wt. %	3	. 84	2	.59	10	0.8	
	Vol. %	89	9.5	98	98.0		95.9	
	Discharge Co- efficient,	0.0	00312	0.0	00249	0.	00292	
	C_{D} , sec, $^{-1}$	(0.	01055)	(0.0	01369)	(0.	00686)	
. C.	Compaction							
	Density, gm/cc,@30 KSI	2	. 60	3	. 35	1	.53	
D.	Burning Rate, in/sec, vs P _c (PSI pressed @ 30 KSI	A) 0.48(P/	1000) 0.29	0.21(P/	1000)		_	
Ε.	Sensitivity to Ignition							
	Impact	225	kg-cm	>300	kg-cm	42	kg-cm	
	Friction		so screw) psi (ABL)		so screw) psi (ABL)		so screw) psi (ABL)	
	Electrostatic Discharge	>6J @	5000 V	>6J @	5000 V	0. 0 056J	@ 5000 V	
	Autoignition Temperature	>4	00°C	>4	00°C	22	9°C	

TABLE 3.12

COMBUSTION PRODUCTS OF CANDIDATE PROPELLANTS

						•		
Cor	mposition 3.3.	4	Compo	sition 3.8.	1	Сотр	osition 3.8.	8
Specie	Chamber	Throat	Specie	Chamber	Throat	<u>Specie</u>	Chamber	Throat
Cr	0.00028	0.00011	C L	0.00020	0.00008	В	0.00011	0.00002
н	0.01575	0.00597	HC1	0.02487	0.02532	BCL2H	0.00052	0.00046
Н2	1.28937	1.29426	Н2	0.82210	0.82506	вн2	0.02901	0.01109
ин3	0.00016	0.00012	NH3	0.00004	0.00003	BN	0.00003	0.00001
N2	0.14712	0.14714	N2	0.01249	0.01250	вон	0.00106.	.0.00052
ОН	0.00001	0.00000	ZRCL4	0.00003	0.00001	B20	0.00044	0.00012
ZRN8	0.00155	0.00000	ZRO28	0.54281	0.00010	CL	0.00078	0.00061
ZRCL4	0.00001	0.00000	H2O	0.00129	0.00110	нв02	0.00074	0.00067
NH2	0.00002	0.00001	WCL2	0.00004	0.00002	н2	3.55937	3. 58573
ZRCL3	0.00003	0.00000	ZRCL2	0.00002	0.00000	NH2	0.00003	0.00001
HC1	0.03363	0.03391	W8	0.16356	0.00003	WCL2	0.00015	0.00012
H20	0.00328	0.00321	н	0.00962	,0.00368	ZRCL3	0.00711	0.00410
ZRCL2	0.00001	0.00000	ZRCL3	0.00007	0.00001	ZRO	0.00011	0.00001
ZR028	0.74496	0.00004	ZRN8	0.00050	0.00000	BN\$	1.26001	0.04158
T °K	2799	2479	T °K	2799	2481	ZR02\$ '	0.35327	0.00525
			9.		9 .	BCL2	0.00039	0.00027
	•					ВН	0.00017	0.00004
	•				* *	вн3	0.00249	0.00100
			•			BOCL	0.002040	.0.00179
			a * a		40	B203	0.00015	0.00012
	•					H	0.04350	0.02849
		8			0.0	ZRCL2	0.00305	0.00096
					12	ZRH	0.00001	0.00000
		f		4	H	BC1	0.00933	0.00503
	•		1			BCL3	0.00001	0.00001
						BH2N	0.01234	0.00543
	•					ВО	0.01106	0.00570
						B202	0.00422	0.00243
	* a .			e)	10	в3н303	0.00001	0.00001
					200	HC1	0.09305	0.11072
					(9)	H20	0.00093	0.00084
						инз	0.00023	0.00014
			g 5			ZRCL	0.00032	0.00005
: \$ -	Condensed Spec	ie				ZRCL4	0.00240	0.00263
					**	W\$	0.03831	0.00003
	20					N2	0.10551	0.08825
					3 89	T °K	2807	2619

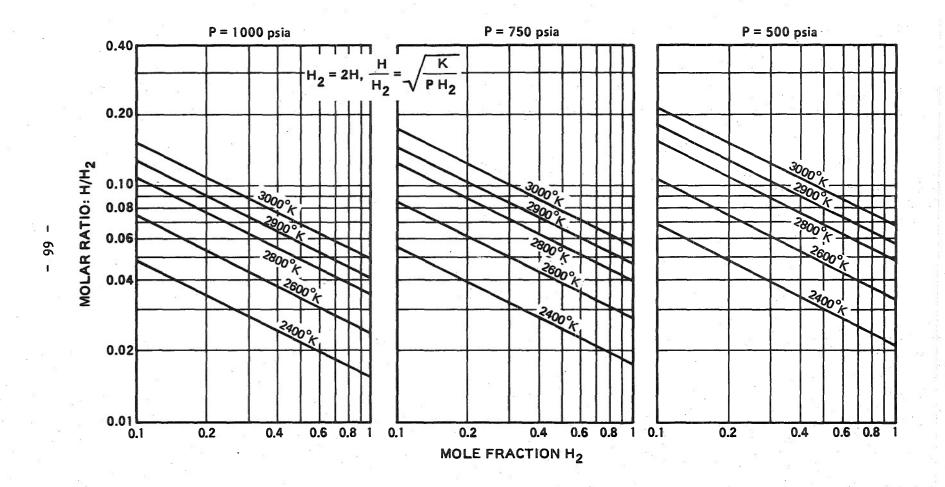


Figure 3.1. Effect of Temperature, Pressure and Concentration on ${\rm H_2}$ Dissociation.

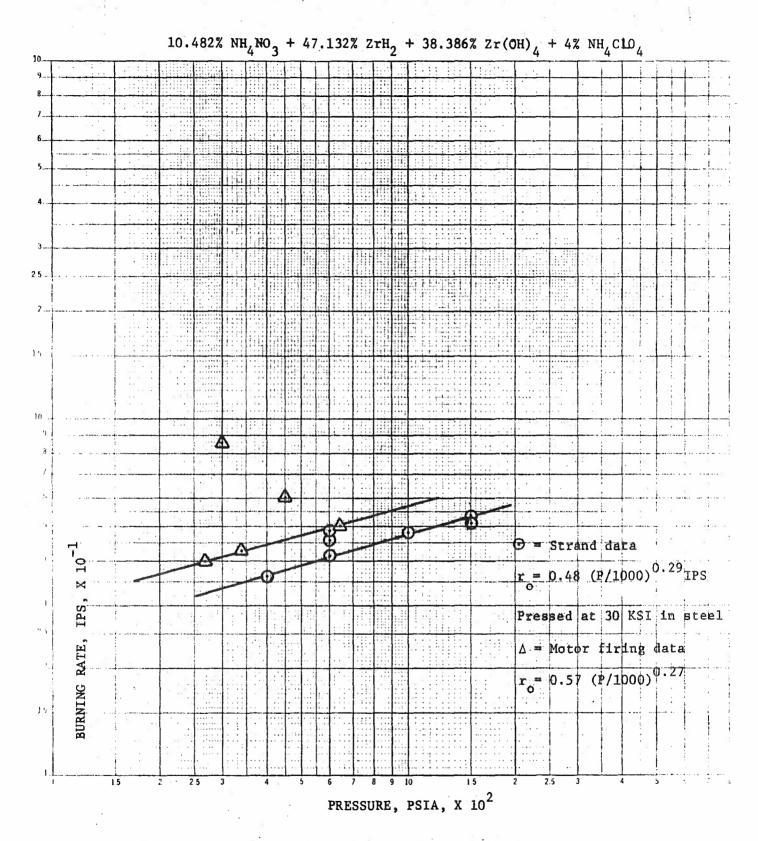


Figure 3.2. Burning Rate of Comp. 3.3.4.

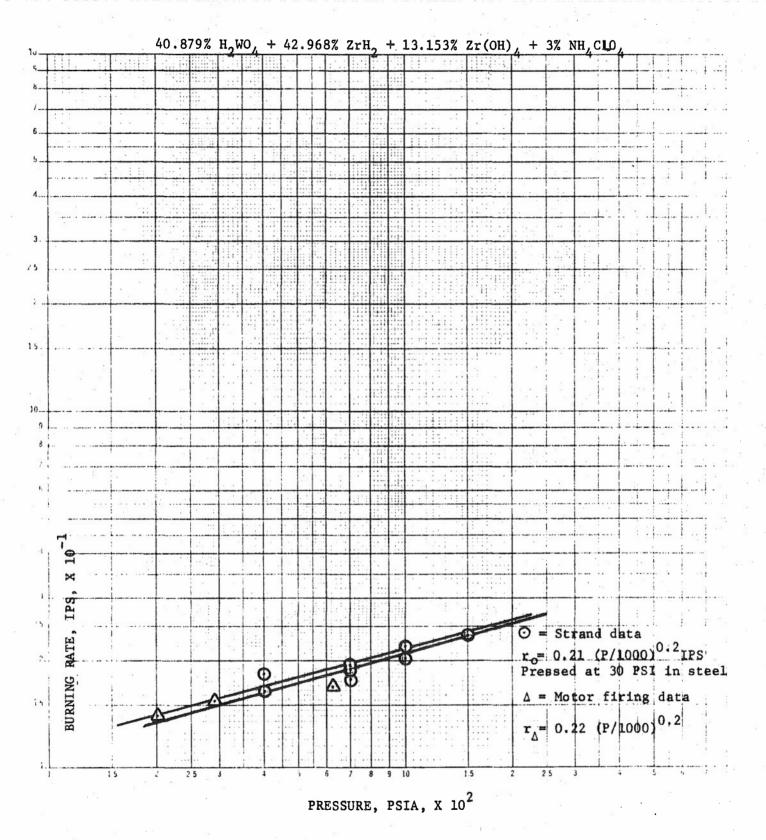


Figure 3.3. Burning Rate of Comp. 3.8.1.

The propellant composition has excellent safety properties. Because of the use of $\mathrm{NH_4NO_3}$ and $\mathrm{Zr(OH)_4}$, the composition must be processed in a low humidity facility (<50% relative humidity at 70°F is satisfactory). Because of the phase change exhibited by $\mathrm{NH_4NO_3}$, the propellant cannot withstand temperature cycles of -40/+140°F without loss of some structural strength.

3.4.3 Properties of Composition 3.8.1

This composition is the best of the thermite systems (in terms of yield and purity) comprised of off-the-shelf ingredients. An excellent clinker is formed on combustion at all pressures. The total gas yield of the composition is only 2.59% by weight, the remainder being ${\rm ZrO}_2(\ell)$ and ${\rm W(c)}$. This propellant also has excellent safety properties. The use of ${\rm Zr(OH)}_4$ requires dry processing and handling facilities but presents no major problem. The composition can be temperature cycled with no adverse affects.

3.4.4. Properties of Composition 3.8.8

This composition is one of the best of any of the systems evaluated, in terms of yield and purity. It contains HDB, and is the result of an extensive laboratory-scale search for ways and means to use the compound in propellants. It proved to be a difficult goal to obtain, because the material on hand was from various production lots and this seemed to affect propellant properties to a great extent. The propellants studied proved to be sensitive to ignition by mechanical means and difficult to burn uniformly over a wide pressure range, and difficult to store and use, even in a dry environment.

Initially, work with the Composition 3.9.8 (WO $_3$ /Zr/HDB/NH $_4$ ClO $_4$) indicated unstable combustion when 1/2 inch diameter pellets pressed at 30 KSI were burned at 1500 psi. At combustion pressures (Pc) <1000 psi, the pellets burn uniformly at a rate r = 2.51 (Pc/1000) $^{0.58}$ in/sec. At Pc = 1500 psi, these pellets explode when ignited. When pressed at 20 KSI, the pellets exploded when burned at 400 psi. Substitution of H $_2$ WO $_4$ for WO $_3$ (i.e., Composition 3.8.8) reveals a similar result, although there was some evidence of HDB batch effects on this phenomenon since one of the batches of HDB gave uniform burns at 1500 psi, while another exploded. Substitution of hydrazine nitrate (N $_2$ H $_5$ NO $_3$) for the oxidizer also resulted in explosive burning at 1500 psi. These results indicate that the substitution of oxidizer is not sufficient to stabilize high pressure combustion of HDB propellants.

The substitution of either ${\rm ZrH_2}$ for ${\rm Zr}$ and/or ${\rm NH_4C1}$ (or ${\rm N_2H_5C1}$) for ${\rm NH_4C10_4}$ in propellants containing HDB and either WO $_3$ or ${\rm H_2WO_4}$ gives yield values in the range 3-5% by weight (Comps 3.8.6 and 3.8.7 and 3.9.7) and as a result are not particularly interesting. Neither compound is energetic and consequently results in compositions containing low values of HDB content and therefore low values of H $_2$ yield. Substitution of Al for Zr yields combustion products which are significantly contaminated with dissociation products of ${\rm Al_2O_3}$ (which is not as thermally stable as ${\rm ZrO_2}$). In these latter systems, species such as AlBO $_2$ and BOC1 are also important combustion products.

The use of another chloride as a substitute for $\mathrm{NH_{\Delta}C10_{\Delta}}$ is possible in propellants containing N2H5NO3 as the oxidizer. For compositions containing other oxidizers, the favored chlorides are mixtures of NH_4C10_4 and a fuel (e.g., Zr or ZrH2) stoichiometric to the major metal oxide (e.g., ZrO2). These are high-energy compositions that release HCl as a major combustion product and require large amounts of coolant (HDB, etc.) to maintain the flame temperature at the desired value of 2800°K. As noted above, the use of NH_LCl in place of NH_L ${\rm C1O_4}$ results in low values of ${\rm H_2}$ yield with either ${\rm WO_3}$ or ${\rm H_2WO_4}$ as oxidizers, because the NH Cl is a coolant as well as an HCl donor and the HDB content is accordingly reduced. The use of $N_2H_5NO_3$ overcomes this difficulty because of the greater energy of ${\rm N_2H_5NO_3/Zr}$ or ${\rm ZrH_2}$ mixtures compared to those of ${\rm WO_3}$ or $\mathrm{H_2WO_4/Zr}$ or $\mathrm{ZrH_2}$. Thus, the use of low-energy chlorides requires the use of a high-energy oxidizer in order to maintain acceptably high values of H2 yield. An alternate candidate oxidizer is NH2NO3. Toward this end, studies were directed toward the system $N_2H_5NO_3/Zr/HDB/N_2H_5C1$ (composition 3.5.10). The commonality of the ingredients is conducive to stability and compatibility. The composition burns uniformly at pressure up to 1500 psi at a rate of r =1.05 $(P/1000)^{0.65}$ ips, when compacted in 1/2 inch diameter steel tubes at 30 KSI (compaction density = 1.63 gm/cc). Unfortunately, the (H_2 + HCl) gas purity is (85.7%) which is too low to be of interest. Also, the composition is extremely moisture sensitive. At values of relative humidity of 40% (at 70°F), the composition physically degrades through chemical reaction. Under desiccated conditions, no evidence of chemical reaction is seen. A slight improvement in gas purity can be attained by use of the stoichiometric mixtures $Zr(OH)_{\Lambda} + ZrCl_{\Lambda}$ as the HCl generator (Composition 3.5.12).

The best solution seemed to be through the use of HDB as a monopropellant coupled with the oxidizer/fuel/chloride as a second component in a bipropellant motor configuration. The propellant composition 3.8.8 can be divided into two propellants, one comprised of HDB, the other comprised of the remainder of composition 3.8.8, i.e., H2WO4, Zr, and NH4ClO4 (called The physical, ballistic, and safety properties of these two propellants are given in Table 3.13 and 3.14 and Figures 3.4 and 3.5. presses well, and ignites and burns uniformly at all pressures. The Mix 22 burns with a zero pressure exponent, which is characteristic of thermites. this technique, the compatibility problems of chemical reaction and unstable combustion of the HDB at high pressures are obviated. It should be noted that neither composition forms clinkers on combustion, the HDB because it is too gassy, and Mix 22 because it is too hot due to ZrO, boiling at 4590°K. However, it should be noted that the combustion products of Mix 22 form clinkers when mixed with the combustion products of HDB in the bipropellant configuration. This effect is described in Section 3.5.3.3 below.

3.5 Sub-Scale Motor Firings

3.5.1 General

The objective of the sub-scale motor firings was to determine a satisfactory separation technique between gases and solids for each of the candidate propellants. For these firings, two sizes of propellant grains were evaluated, one being 1-5/8"in diameter and the other being 2.5" in diameter. Both were cylindrical end-burners, inhibited on the circumferential surface and one end. The nominal grain length was 1". Since the effectiveness of filtration (or other separation technique) governs the mass discharge rate and therefore the combustion pressure, it was the practice to use oversized nozzles in first shots, with both filter design and nozzle throat size changed by steps to approach the desired value of combustion pressure. It is recalled that two values for discharge coefficient were given for each candidate in Table 3.11. They are discussed below.

TABLE 3.13
PROPERTIES OF HDB AND MIX 22

	Property	HDB	M1x 22
A.	Composition, Wt. %	100 HDB	19.285 H ₂ WO ₄
			54.612 Zr
			26.103 NH ₄ ClO ₄
В.	Theoreticals		
	1. Flame Temperature, °K	1878	4590 2.88
	2. (H ₂ /HC1) molar		2.80
	3. Yield (H ₂ +HCl)		
	Wt. %	16.9	4.39
	Vol. %	100.0	41.3
	4. Discharge coefficient, sec-1.	(.00237) (.00554)	(0.00526) (0.0101)
C.	Compaction Density, gm/cc @ 30 KSI	0.837	3.18
D.	Burning Rate, in/sec, vs P (psia)	0.70(P/1000) ^{0.48}	2.54 @ all values of P
E.	Sensitivity to Ignition		
	1. Impact	41 kg-cm	125 kg-cm
	2. Friction	Neg (Esso Screw) Neg @ 90°, 400 psi	Neg (Esso Screw) Neg @ 75° 100 psi
	3. Electrostatic Discharge	0.0013 J @ 5000 V	.00135 @ 5200

COMBUSTION PRODUCTS OF TWO-COMPARTMENT HDB PROPELLANT

	HDB		91 SE	MIX 22	
SPECIE	CHAMBER	THROAT	SPECIE	CHAMBER	THROAT
BH2N	0.00009	0.00001	CL.	0.06385	0.06463
NH3	0.00005	0.00001	HC1	0.09957	0.10410
BN8	3.34872	0.00048	Н2	0.29424	0.29971
вн2	0.00022	0.00001	N	0.00053	0.0 0038
н2	8.37201	8.37296	NH3	0.00001	0.00001
внз	0.00019	0.00002	N2	0.10961	0.11005
н	0.00078	0.00010	он .	0.01629	0.01252
N2	0.00018	0.00001	WCL	0.00008	0.00006
T °K	1877	1609	ZRCL	0.00220	0.00169
•			ZRCL4	0.00002	0.00002
	152 ·		ZRO2	0.00671	0.00547
			ZRO2°	0.54271	0.01171
•		a .	80 18		
			CL2	0.00006	0.00005
		9	Н20	0.02176	0.01863
• .			NH	0.00010	0.00007
•			· NO	0.00224	0.00159
		ā ·	02	0.00026	0.00016
	1000		WCL2	0.02599	0.02386
	* 192	K.O.	WO2	0.00114	0.00075
		S^{2}	ZRCL2	0.00140	0.00121
			ZRH	0.00008	0.00005
*			н	0.29482	0.28954
*		• •	H2W04	0.00001	0.00001
			NH2	0.00006	0.00004
		•	O .	0.00900	0.00673
		· a	W	0.00016	0.00010
			WO	0.00160	0.00110
			WO3	0.00003	0.00002
		·	ZR .	0.00034	0.00021
	·		ZRCL3	0.00049	0.00045
		i la	ZRO	0.04472	0.03515
		· · · · · · · · · · · · · · · · · · ·	.M.	0.04816	0.00313
	· ·	_ 72 _	T °K	4608	4398

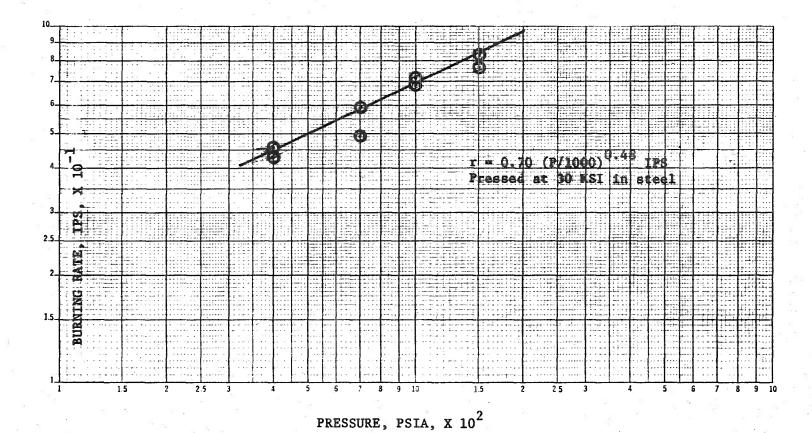
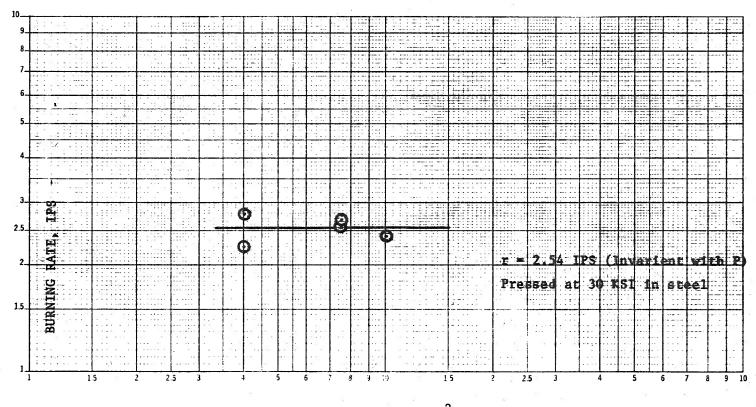


Figure 3.4. Burning Rate of HDB.



PRESSURE, PSIA, X 10²

Figure 3.5. Burning Rate of Comp 22 - 19.285% H_2WO_4 + 54.612% Zr + 26.103% NH_4CIO_4 .

For given values of propellant mass burning rate (\hat{m}) and combustion pressure (P_c), the nozzle discharge coefficient (C_D) is used to size the nozzle throat area (A_t) by means of the continuity equation,

$$\dot{m} = P_c C_D A_t \tag{1}$$

By continuity, $\mathring{\mathbf{m}}$ is also related to propellant burning surface area(S), burning rate (r), and density (ρ_p) and to the gas velocity (\mathbf{v}_t) and density (ρ_g) at the throat:

$$\dot{\mathbf{m}} = \operatorname{Sr} \rho_{\mathbf{p}} \tag{2}$$

$$= A_{t} v_{t} \rho_{g}$$
 (3)

Combining Equation 1 and 3 yields

$$C_{D} = v_{t} \rho_{g}/P_{c} \tag{4}$$

For adiabatic combustion at an assigned value of P_c coupled with isentropic flow to sonic velocity, it is seen that C_D is thermodynamically derived. For the motors of interest, it is determined from only the gaseous combustion products which are ejected from the motor, since the condensed species are assumed to be retained within the motor as a clinker or a filter cake. For this case of uncoupled flow (gas only), the derived values of C_D are the first-reported numbers in Table 3.11. In the event that the condensed species are ejected from the motor because no clinker is formed or because of failure of the filter system, than a higher value of C_D is computed. These values are reported in parenthesis in Table 3.11 and represent an upper limit on the theoretical value for this parameter. In this case (coupled flow, i.e., gas and solids), considerably greater values of P_c are computed for a given value of A_t/S . Combining equation 1 and 2, and expressing $r = kP_c^{\ n}$, then

$$\left(\frac{P_{c}, \text{ coupled}}{P_{c}, \text{ uncoupled}}\right) = \left(\frac{C_{D}, \text{ uncoupled}}{F \cdot C_{D}, \text{ coupled}}\right)^{\frac{1}{1-n}}$$

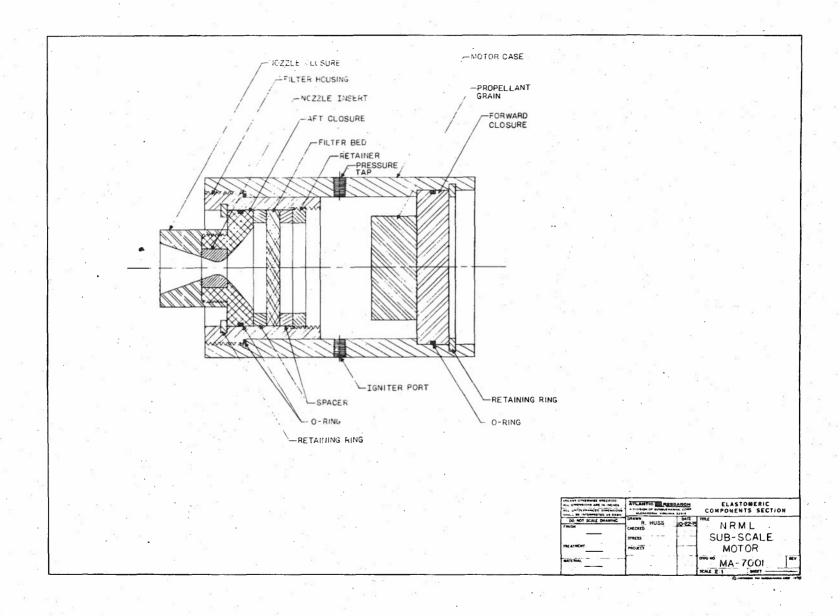
where F = mass fraction of gaseous combustion products.

For compositions 3.3.4, 3.8.1, and 3.8.8, the P_c ratios are 6.26, 9.55 and 8.03, respectively. Thus, for a given value of A_t/S computed on the basis of uncoupled flow (which is the desired phenomenon), very high pressure rises can occur if the propellant does not clink or if the filter system fails. This is why a step-wise approach was taken in sizing the throat for the motor firings.

3.5.2 Sub-scale Motor Configurations

The motor hardware that was constructed was comprised of a heavy-wall case and closures, with a filter housing located between the combustion chamber and nozzle throat. The intent was to allow for maximum flexibility in filter bed configuration and propellant mass burning rate. Of the three candidate propellants, 3.8.1 clinks best when burned. For this composition, a simple filter bed comprised of tungsten screens was evaluated. This configuration is shown in Figure 3.6.

At the other extreme, composition 3.8.8 forms no clinker, and the solids content of the combustion products is 82.7% by weight. For this composition, cyclones or centrifugal separators are best. Attributes of cyclones include simplicity of design and operation, while centrifugals offer a more compact design envelop. Re-entrainment of particles is a phenomenon peculiar to both techniques, since high gas velocities are involved. For the cyclone, a tangential velocity on the order of 200 ft/sec is required for separation of 10-u particles in the sub-scale envelop (3" dia). For this reason, either an impingement separator or filter bed located upstream of the nozzle outlet may be required to remove residual solids. Cyclones and centrifugals can both be operated at high pressure. The principal drawback of the cyclone is size, while that of centrifugals is the need for an external power source to drive the turbine and the increased mechanical complexity. It should be noted that centrifugation is an extremely effective technique for solids removal. In previous work performed with spinning motors, complete removal of condensed combustion products (A1 $_2$ 0 $_3$ (l), C(c), and B(l) was obtained for a 6-inch diameter motor spinning at 900 rpm (centrifugal force ~70 g's). The centrifugal effect for a full-scale motor probably can be attained by either spinning the motor, as above, or (more conventionally) by spinning the combustion products by means of a rotoclone, or cinder fan. For compositions 3.8.8 and 3.3.4, a cyclone separation chamber was employed, primarily because of design simplicity and ease of use at high pressure. This motor configuration is shown in Figure 3.7. The cyclone effect is imposed on the combustion products by the deflector plate located above the propellant burning surface to direct the flow tangentially along the motor wall. A re-entrant nozzle tube was employed to minimize the effects of re-entrainment.



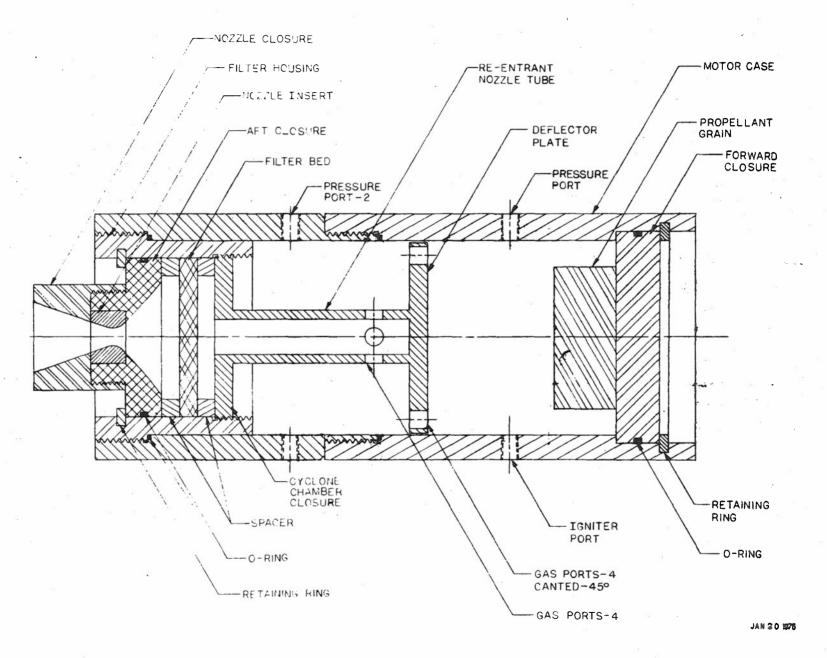


Figure 3.7. Motor in Cyclone Chamber.

3.5.3 Subscale Motor Firing Results

3.5.3.1 Composition 3.8.1 (H₂WO₄/ZrH₂/Zr(OH)₄/NH₄ClO₄)

The simplest motor configuration, that is shown in Figure 3.6, was used for these firings. The propellant forms excellent clinkers when burned, so that the only filter arrangement was comprised of 3 tungsten screens in series located in the filter housing. The screens are 1.60" diameter x 4 mil x 24 mesh, spaced 1/4 inch apart in series. Firing results are given in Table 3.15. Initial firings were oversized nozzles. It is seen that as $P_{\rm C}$ increases, the residue weight increases. Very little ash was found on the screens indicating a clean gas situation. Burn rates were difficult to measure from the motor firings at the lower values of $D_{\rm T}$, because the ratio of motor volume to $D_{\rm T}$ is very large (motor volume $\stackrel{\sim}{}$ 350 cc), and the pressure decay times are significant fractions of the burn time. For no heat losses the rate at which pressure in the motor chamber decays to ambient is derived from the Perfect Gas Law and the continuity equation:

$$\frac{dP}{dt} = -\frac{RT}{V} \frac{dn}{dt}$$

$$= -\frac{RT}{V} \frac{P C_D^A t}{(MW)}$$

where

P = Pressure in the chamber t = Time

T = Temperature of the gas R = Gas constant

n = Number of moles of gas V = Volume of chamber

MW = Molecular weight of gas

For firings 10-5, -6 and -7, the total burn times were 6.26, 6.12 and 6.52 sec, respectively; and the computed decay times are 0.22, 0.59 and 1.44 sec, respectively; resulting in computed burn times of 6.26 - 0.22 = 6.04, 5.53, and 5.08 sec, respectively. The tabulated burning rates, which are plotted in Figure 3.3, are based on the computed times and are seen to agree very well with strand data.

TABLE 3.15 Subscale Motor Firing Results: Composition 3.8.1 $({\rm H_2WO_4/ZrH_2/Zr(OH)_4/NH_4C1O_4})$

Firing Number	Grain Dimensions Dia(in) x L(in) x Wt(gm)	Nozzle Throat Diameter Dt, in	Filter Configuration	P _c , psia	Average* Burn Rate ips	Residue, % of Theoretical	Remarks
10-1	2.512 x 0.586 x 122.6	0.402	3 W screens	Ambient	- ·	81.5	Ambient burn
10-2	Igniter shot				1 2		
10-3	1.688 x 0.845 x 98.4	0.128	3 W screens	55	, , , , , , , , , , , , , , , , , , ,	93.0	3-sec fignition delay
10-4	1.688 x 0.860 x 99.4	0.089	3 W screens	197	-	100.0	See text
10-5	1.688 x 0.842 x 98.5	0.069	3 W screens	200	0.14(0.14)	99.7	See text
10-6	1.688 x 0.858 x 101.5	0.048	3 W screens	294	0.155(0.15)	97.0	See text
10-7	1.688 x 0.845 x 98.5	0.033	3 W screens	631	0.17(0.185)	98.5	See text

^{*}Computed from the quotient web : computed burn time (see text). Numbers in parenthesis are from strand data (see Figure 3.3).

3.5.3.2 Composition 3.3.4 (NH, NO /ZrH /Zr(OH) /NH, (CO,)

Results of these firings are given in Table 3.16. In the first firing, 9-1, the motor configuration of Figure 3.6 was used. The filter configuration was comprised of 3 Tungsten screens, 1/6" diameter x 4 mil x 24 mesh arranged in series 1/4 inch apart. The firing trace was very erratic and characterized by two pressure peaks. Post-firing inspection indicated the screens had been ruptured. It is known that the composition does not clink well at low pressures, and apparently the screens filled and ruptured, causing the second pressure peak. The average pressure is a poor value, and the resultant burning rate is therefore also poor.

In the second motor, also based on the 3.6 configuration, the filter housing was outfitted with three baffles made of 5 mil tungsten sheet and followed by a tungsten screen. The spacing between each baffle was 1/4 inch. The baffles were 1.6-inch dia discs with a 1/4 inch segment removed, each segment located on opposite diameters so that the flow through the baffles was serpentine. Post-firing examination revealed that the baffles were bent concave in the direction of flow, that the void space between baffles was filled with residue, and that the screen had become filled and ruptured above the open segment of the last baffle. The motor pressure had built up gradually to a value of 1500 psia 0.4 sec after ignition and dropped instantaneously to ~400 psia and remained at approximately this value for the remainder of the firing (0.4 sec). Apparently, the rise to the pressure spike is due to filling the baffles and screen, and the drop to operating pressure is due to failure of the screen after it had become completely filled. As above, the average pressure is not a very good one, due to the spike, and therefore the motor and strand burning rates, which are shown in Figure 3.2, do not agree well.

These firings indicate that the propellant does not clink well enough to permit the use of a simple filter configuration. For this reason, the remainder of the firings were carried out with the motor configuration of Figure 3.7, which includes a cyclone chamber and re-entrant nozzle. Firing 9-3 was carried out with this motor and the same filter configuration of 9-2. Even though the \overline{P}_c was lower, the retention of residue was considerably better, and the agreement between motor burning rate and the strand value improved. For this firing, most of the residue was caught in the cyclone chamber. None was found between the baffles nor on the screen, both of which were intact and not deformed or ruptured as in the previous firings. Evidently,

Subscale Motor Firing Results: Composition 3.3.4 $(NH_4NO_3/ZrH_2/Zr(OH)_4/NH_4C1O_4)$

TABLE 3.16

1	Firing	Grain Dimensions	Nozzle Throat Diameter	Filter	Average	Average* Burn Rate	Residue,	
	Number	Dia(in) x L(in) x Wt(gm)	D_{t} , in	Configuration	P _c , psia		Theoretical	Remarks
	9-1	2.50 x 0.947 x 194.2	0.334	Figure 3.6; 3 W screens	300	0.86(0.33)	55.1	Screens failed
	9-2	2.51 x 0.481 x 97.7	0.250	Figure 3.6; 3 W baffles 1 W screen	400	0.60(0.36)	67.3	Screens failed
	9-3	1.688 x 0.564 x 53	0.099	Figure 3.7; 3 W baffles 1 W screen	264	0.40(0.32)	93.5	See text
	9-4	1.688 x 1.076 x 97.22	0.067	Figure 3.7; ZrO ₂ matt and W screen	335	0.43(0.35)	100	See text
	9-5	1.688 x 1.061 x 96.61	0.055	Figure 3.7; ZrO ₂ matt and W screen	627	0.50(0.42)	98.7	See text

^{*}Computed from the quotient web : burn time, Numbers in parenthesis are from strand data (see Figure 3.2).

the cyclone chamber is very effective in removing the solids from the combustion products.

The next two firings were also made with the re-entrant motor and with ZrO_2 -fiber filter matt supported by tungsten screens in the filter housing replacing the baffle screen arrangement of the previous firings. For these firings, retention of the slag is essentially quantitative. In all three motor firings (9-3, -4, and -5), the motor burning rates are significantly greater than the strand rates ($\sim 20\%$). The reason for this is unexplained. It is possible that the grain breaks up near burnout, resulting in apparently shorter burn times. This could be due to inhibitor failure which consisted of a thinly painted coat of titanium dioxide filled epoxy.

3.5.3.3 <u>Composition 3.8.8 (H₂WO₄/Zr/HDB/NH₄C1O₄)</u>

As reported previously, this composition exhibited combustion instability when burned at 1500 psia. The possibility exists that pressure spikes occurring at ignition can trigger explosions instead of uniform combustion, so that no attempts were made to burn pellets of this composition in motors. Therefore, the bi-propellant two-compartment approach was taken, and motors comprised of HDB pellets and Composition 22 pellets were made up. The web lengths of the pellets were sized to permit equal burning times, based on strand data, and the relative number of pellets was governed by the required mass burning rates (39.97% HDB + 60.021% Composition 22). The motors were comprised of 10 pellets of HDB of dimensions of 0.5" diameter x 0.2" L x 0.52 gm weight, and one pellet of Composition 22, 0.5" diameter x 0.87" L x 8.9 gm weight. The pellets were inhibited on one end and the circumferential surface and functioned as end burners in the motor configuration of Figure 3.7 (cyclone chamber and reentrant nozzle).

Three motors were fired. Because of the small amount of propellant, residue weights could not be measured accurately. Firing results (P_c vs D_T) are given in Table 3.17. The first motor was "fired for effect." The next two were fired with D_T = 0.10 inches, the 8A-2 with a ZrO_2 fiber matt filter supported by a tungsten screen and the second (8A-3) with two ZrO_2 matt filters in series, each supported by a tungsten screen and separated by 1/4 inch. In the first firing (8A-2) post firing observation indicated a light fluffy ash (BN) throughout the combustion chamber, cyclone chamber, and re-entrant tube.

TABLE 3.17

Subscale Motor Firing Results: HDB + Composition 22 $(H_2WO_4/Zr/NH_4ClO_4)$

Firing Number	Nozzle Throat Diameter , in		Filter Configuration		Average Pressure P _C , PSIA
					1000
8A-1	0.237	Motor	3.6, screen only	19	80
8A-2	0.100	Motor	3.7, ZrO ₂ Matt + W	Screen	575
			, -		
8A-3	0.100	Motor	3.7, 2 ZrO, Matts +	W screens	560

The combustion chamber also contained an acorn-sized metallic clinker weighing 3.7 gm, which was the residue from Composition 22. The filter and screen were intact.

The residue from 8A-3 was very similar to the above, but the filters had been bent. Apparently the use of two filters gives rise to a value of ΔP across the filters that is too high. Since there was no difference in the amount of residue on the filters between the second and third test, one filter would prove sufficient.

The motor firing traces for firing 8A-2 and 8A-3 are shown in Figure 3.8a and 3.8b. As can be seen from the traces, the chamber pressure is essentially constant. Thus, even though the HDB is tempermental when combined with other ingredients, in a two-compartment motor the ballistic properties of this system are excellent.

Theoretical calculations of the propellant combustion products have been made and have been compared with experimental values. The experimental procedure used in the analysis was to combust small pressed propellant samples in a stainless steel bomb. The interior of the bomb and all exposed metal surfaces were sandblasted, cleaned with trichloroethylene, then teflon coated. The teflon coating was primarily used to prevent absorption of HCl on cold metal surfaces — a phenomenon observed under previous programs. Before and after teflon coating, the bomb was baked out in an oven overnight at 800°F to remove any absorbed gases. All fittings, valves and transfer tubing was constructed of teflon or a similar HCl unreactive plastic.

The bomb was purged a minimum of 5 times with argon and vacuum pumped between purges to ensure no contamination with residual air. Bomb volume was exactly measured via displacement and the propellant weight recorded. Ignition was carried out by a hot wire technique, to avoid any introduction of contaminating gases, under one atmosphere of argon. Upon cooling, a plastic syringe sampled gases through a septum fitted directly to the bomb and transferred the gases directly to a calibrated gas chromatograph for analysis. Transit time and residence time in the bomb were minimized. The results of a gas chromatograph of the sample gave results that essentially agreed with predicted thermochemistry calculations.

Burn Time = 0.37 sec versus 0.34 from strands P = 575 psia

Grain Configuration:

10 HDB pellets 0.2 in. L X 0.495 dia.

1 Comp. 22 pellet 0.87 in. L X 0.495 dia. All pellets end burners burning on one end

Total propagation weight = 14.1 gm

Motor Configuration:

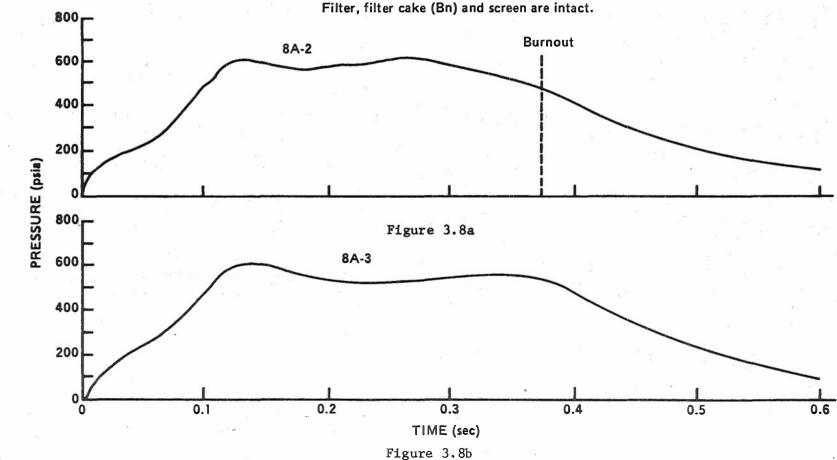
Re-entrant nozzle with cyclone chamber

One ZrO2 fiber mat filter with tungsten screen

 $D_T = 0.1$ in.

Post-firing Configuration:

Comp. 22 clinker in combustion chamber. Bn is uniformaly distributed in combustion cyclone, and re-entrant chambers.



3.6 SUMMARY

The objective of the solid propellant technology effort was to select and characterize three solid propellant candidates. Selection was based on the results of calculations of theoretical flame temperature and combustion product composition. For these calculations, propellant ingredient selection was constrained by state-of-the-art solid propellant chemistry. Successful sub-scale motor firings (1/4 lb nominal propellant weight) were carried out for all three candidates. Design factors for filtration and particle separation techniques for each candidate were delineated.

Of the three candidates, the first represents the best clinker propellant. This propellant is comprised of H2WO4, ZrH2, Zr(OH)4, and NH4ClO4, has a flame temperature of 2799°K and an (H2 + HC1) yield of 2.59% by weight at 98.0% gas purity and $H_2/HC1 = 32.6$ (molar). A simple filter comprised of a tungsten screen was found to be satisfactory for retaining entrained solids in the combustion products. Typical clinker weight >99% of theoretical values were measured for this propellant. The second propellant represents the best attainable from state-of-the-art ingredients. The propellant is comprised of $NH_{\Lambda}NO_3$, ZrH_2 , $Zr(OH)_{\Lambda}$, and $NH_{\Lambda}CIO_{\Lambda}$, has a flame temperature of 2801°K, and an $(H_2 + HC1)$ yield of 3.84% by weight at 89.5% gas purity and $H_2/HC1 = 38.2$ (molar). Although this propellant does not form clinkers as well as the $\mathrm{H_2WO}_{\Delta}$ system, clinker weights >99% of theoretical were attained through the use of a cyclone chamber and a filter bed comprised of ZrO, fiber matt. A similar particle separator/filtration technique was used satisfactorily for the third propellant which is the best attainable from use of advanced propellant ingredients. This propellant is comprised of H2WO, Zr, HDB and $NH_{L}C10_{L}$, has a flame temperature of 2808°K, and an (H₂ + HCl) yield = 10.8% by weight at 95.9% gas purity and $\rm H_2/HC1$ = 33.3 (molar). In this propellant, HDB (hydrazine diborane) is the coolant, HDB is a compound that is easily synthesized and is available on special order. The motor firings for this propellant are especially unique because a bipropellant grain configuration was used in the motors due to the reactivity during combustion of HDB with $NH_{\lambda}C10_{\lambda}$. One grain was comprised of $H_{2}W0_{\lambda}$, Zr, and $NH_{\lambda}C10_{\lambda}$ while the other was comprised of only HDB. The firings yield uniform combustion and pressure time traces, and the filtration system withstood the firing conditions without damage.

The high values of yield and purity of (H₂ + HCl) offered by use of HDB in propellants leads to the conclusion that additional work with this and possibly other HDB propellants is warranted. Synthesis of additional HDB is needed to permit propellant evaluation with a single lot of HDB, eliminating batch effects. Scale-up firings should be made, with the emphasis on BN(c) separation. The combustion products of H₂WO₄/Zr/NH₄ClO₄ form a good clinker when admixed with the auto-combustion products (H₂, BN) of HDB, and for this reason the use of a cyclone or centrifugal separator for only HDB combustion products minimizes the envelope and mechanical impact on the system as well as the heat sink effect of the filter hardware.

4.0 PROOF EXPERIMENT AND LASER SYSTEM CHARACTERISTICS

The end objective of the work described in this report is to develop an H₂-HCl laser device. Prior to design of such a device, it is anticipated that a laser proof experiment will be performed. The objective of the proof experiment is to experimentally determine laser performance at conditions that are of practical interest. The current phase is concerned with determining a low risk approach for the proof experiment and to develop the techniques that will be used in design of the proof experiment hardware.

An early objective of this task was directed at obtaining performance maps as a function of gas composition, chamber conditions, and nozzle contour. The one-dimensional harmonic oscillator program was used to generate these maps using the vibrational kinetics for the H₂/HCl system listed in Table 2.2. These calculations provided (i) a limited performance optimization, (ii) a base line performance for conditions expected to be achievable by propellant technology, and (iii) some effects on H-atom performance degradation.

Some preliminary performance sensitivity studies on stagnation and gas expansion conditions has been described in Phase I. It was shown that best performance was obtained at highest stagnation temperature (neglecting atom effects). However, it was also determined that there was an optimum value of p_0 h where p_0 in stagnation pressure and h is the nozzle throat height. This effect is similar to that observed in the $\mathrm{CO}_2/\mathrm{N}_2$ GDL and is easily explained in terms of the H_2 vibrational energy freezing criterion. In addition, it was shown that maximum performance occurred at lowest HCl concentrations. This latter conclusion may be expected to be modified in favor of larger HCl concentrations in the A.H.O. model and when cavity power extraction is included in the modeling.

Some additional calculations illustrating the points mentioned above are given in Figures 4.1 and 4.2. It is seen that maximum available energy occurs at [HC1] 1 2% for the smaller throat height and $P_{0} = 50$ atm. At the higher stagnation pressure the throat height has a large effect on the performance whereas at the lower chamber pressure the performance is less sensitive. This effect is similar to that which occurs in the 1 2 GDL where performance drops off at large values of the parameter 1 3 has a large parameter 1 4 has a large values of the parameter 1 5 has a large values of the parameter 1 6 has a large values of the parameter 1 6 has a large values of the parameter 1 6 has a large values of the parameter 1 6 has a large values of the parameter 1 6 has a large values of the parameter 1 6 has a large values of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 6 has a large value of the parameter 1 8 has a large value of the parameter 1 8 has a large value of the parameter 1 8 has a large value of the parameter 1 8 has a large value of the parameter 1 8 has a large value of the parameter 1 8 has a large value of the parameter 1 8 has a large value of the parameter 1 8 has a large value of the parameter 1 9 has a large value of the para

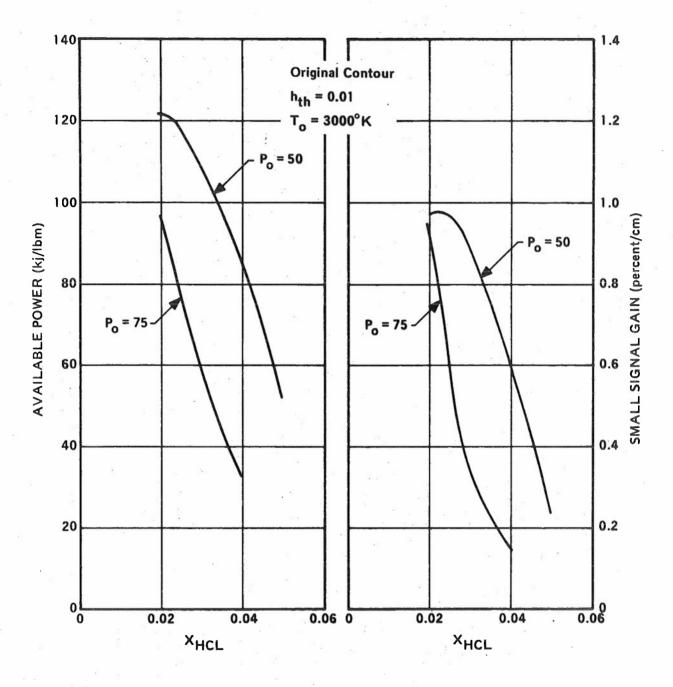


Figure 4.1

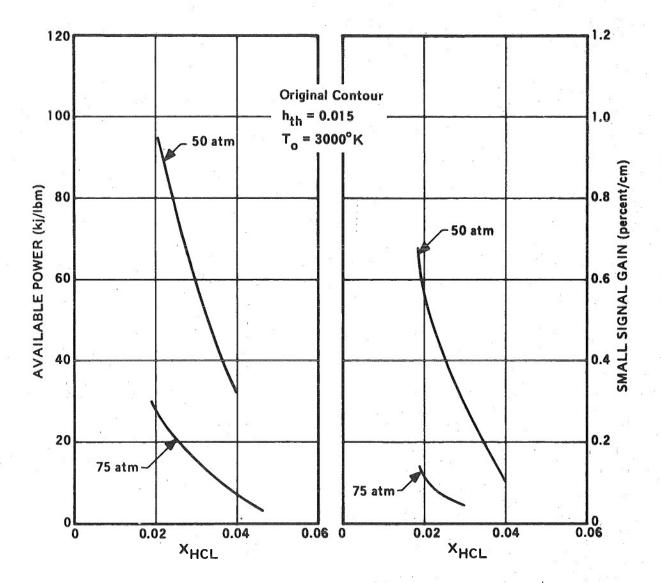


Figure 4.2

Some calculations were also performed using various assumed nozzle contours (or gas density distributions). The assumed density distributions are shown in Figure 4.3. Results of this limited study is Table 4.1. It is clear that some improvement in performance is possible by suitably tailoring the nozzle contour to maximize $\rm H_2 \rightarrow HC1~V \rightarrow V$ pumping. This study was not pursued further pending resolution of the basic kinetics of the $\rm H_2/HC1$ system.

Another technique was used to estimate the optimum area ratio for GDL performance. A calculation was performed comparing the characteristic reaction times for the various vibrational energy transfer processes. An isentropic flow was used to calculate the thermodynamic state associated with each expansion ratio starting from some assumed stagnation conditions. characteristic time of the reaction is equivalent to the vibrational relaxation time for that process at the particular temperature and density that result from the expansion. An example of such a comparison is shown in Figure 4.4. Also in Figure 4.4. is a dashed line that indicates the time for an expanded gas sample to traverse a one meter long cavity. For an efficient device we wish the V+V transfer reaction time to be short compared to the cavity flow time and in turn shorter than all V+T deactivation processes. It can be seen that in Figure 4.4 this criterion occurs generally for A/A* < 30. Calculations like this were performed for a variety of concentration</p> ratios and stagnation conditions to develop insight into optimum expansion ratio for H₂/HCl vibrational pumping.

The results of this analysis is shown in Figure 4.5. The available energy in the HCl at a distance of 100 cm from the throat is given as a function of area ratio for several HCl concentrations. Energies greater than 100 kJ/lb are obtained for HCl mole fractions of 0.01 - 0.02 for A/A* ~ 50. This result is comparable to the previous results in the sense that it gives a rough estimate for the nozzle expansion ratio. It is also apparent from this limited optimization study that the stagnation temperature would be lowered to 2800°K where the propellant technology is more practical with no performance degradation. It was recognized in the Phase I feasibility study that H-atoms which occur at the high stagnation temperature and freeze during the rapid expansion can seriously degrade calculated performance by rapidly deactivating HCl. During this program a series of calculations was performed at successively lower stagnation temperature. Of course, it was expected

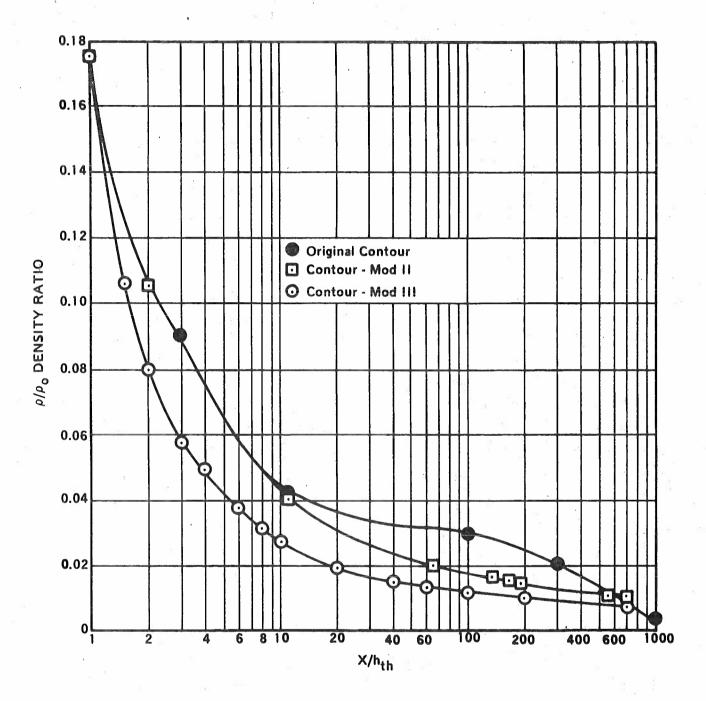


Figure 4.3

TABLE 4.1

Initial conditions: T_o = 3000°K, P_o atm, [HC1] = 0.02

Contour	Available Energy kJ/lb
old	93
I ,	109
II	116

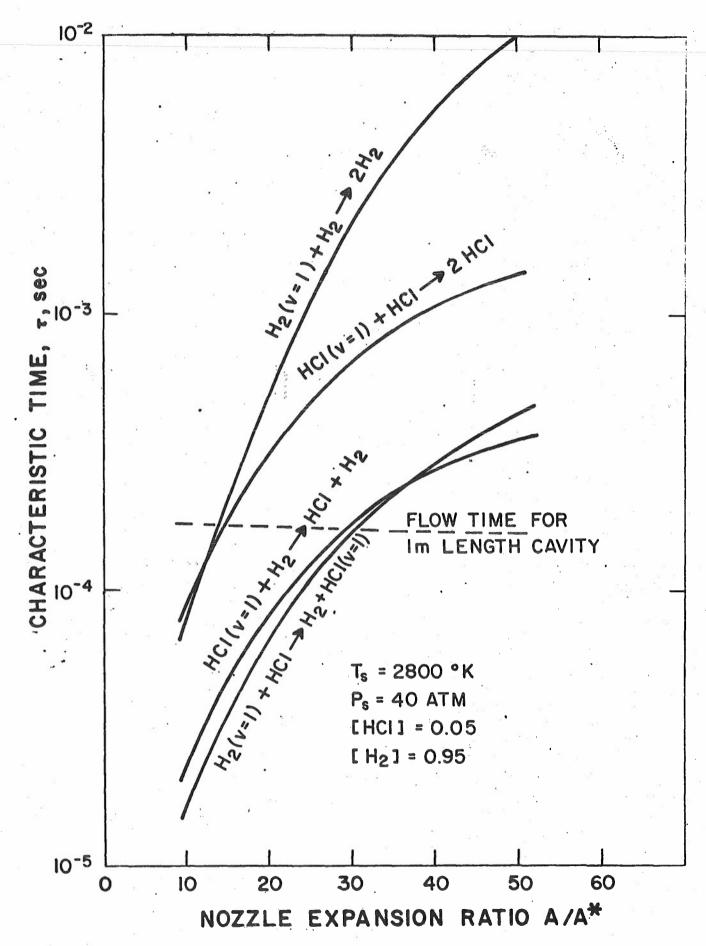


Figure 4.4 Characteristic Times for Various Energy
Transfer Processes as a Function of Expansion Ratio.

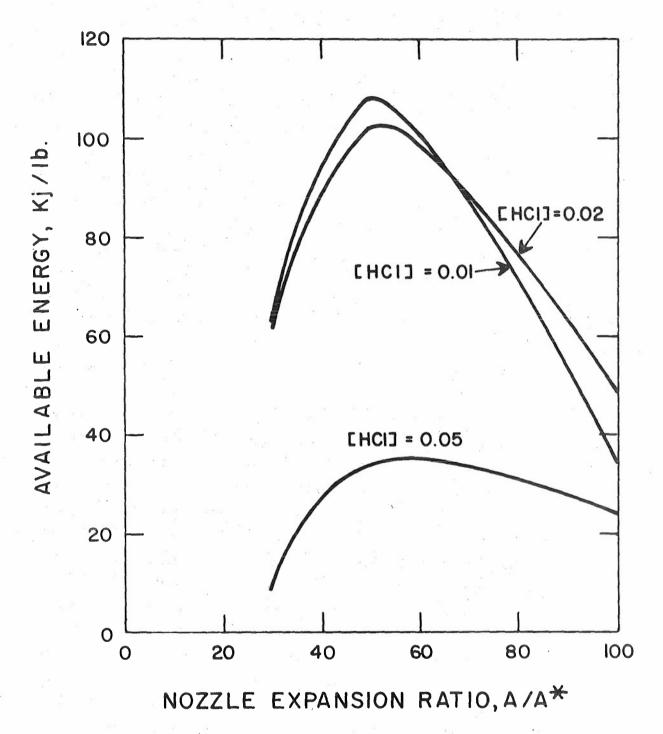


Figure 4.5 Optimized performance map for baseline condition of $T_s = 2800^{\circ} \text{K}$ and $P_s = 40 \text{ atm.}$

that specific energy would decrease, but, at the lower $\mathbf{T}_{\mathbf{O}}$, H-atoms would also decrease and have less effect of performance.

The method employed was to systematically modify the nozzle contour and observe the change in the vibrational temperature of HCl versus axial distance. At temperatures of 2400°K and 2300°K, the specific power computed is 45 kJ/lb and 25 kJ/lb, whereas the gain is .55%/cm and .35%/cm, respectively. Although these values are not as attractive as the values obtained at the higher temperatures, the effect of H-atoms at 2300°K is estimated to be negligible. Thus, from the point of view of a low risk proof experiment a stagnation temperature of 2300°K appears to be a good choice.

These estimates use the equilibrium throat value of H-atoms. The area ratio of the resulting nozzle giving the best performance was 15. For a throat height of .01 cm, the overall length of the nozzle is 2 cm. Since the nozzle has such a low area ratio, a slight increase in gain can be obtained without significant degradation of pressure recovery by placing an additional expansion downstream of the nozzle. The vibrational temperature of HCl for this case is shown in Figure 4.6 along with the results of other calculations at different expansion ratios and chamber pressures. Another implication of Figure 4.6 is that the performance is not as critical on the H₂ vibrational temperature as originally believed. A stagnation temperature of 2300°K for the H₂-HCl mixture gives a frozen vibrational temperature of 2100°K for the H₂. This relatively low H₂ vibrational temperature pumps the HCl to a vibrational temperature of 6500°K.

The harmonic oscillator program was then modified to include the H-atom recombination kinetics to see if the equilibrium H-atom throat value could be obtained using a quick freeze nozzle. The original contour had a subsonic approach section of .1 cm long. This resulted in essentially no H-atom recombination in the nozzle. The subsonic approach section was increased to 1.6 cm and 3.2 cm as shown in Figure 4.7. The resulting H-atom concentration versus x is shown in Figure 4.8. The theoretical results indicate that the equilibrium throat value of H-atom concentration is not achieved. A similar calculation was performed for chamber conditions of 80 atmospheres and 2800°K. The H-atom concentration is shown in Figure 4.9. Again, the H-atoms do not reach the throat equilibrium value although at the large subsonic lengths the differences are not large.

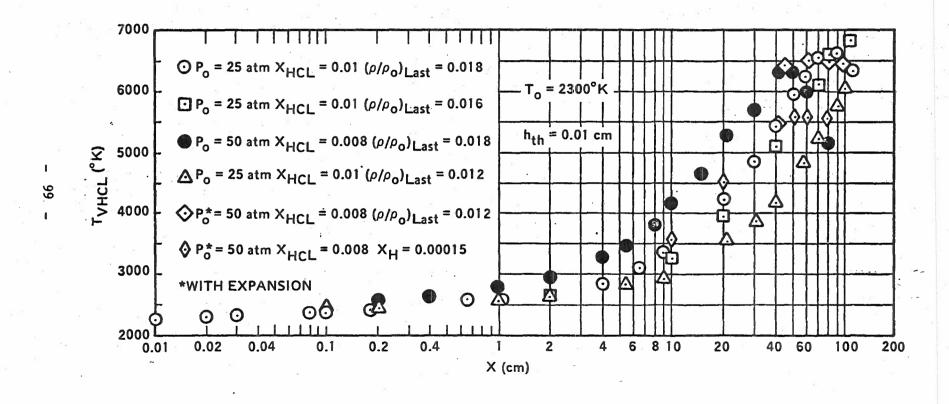


Figure 4.6 HCL Vibrational Temperature.

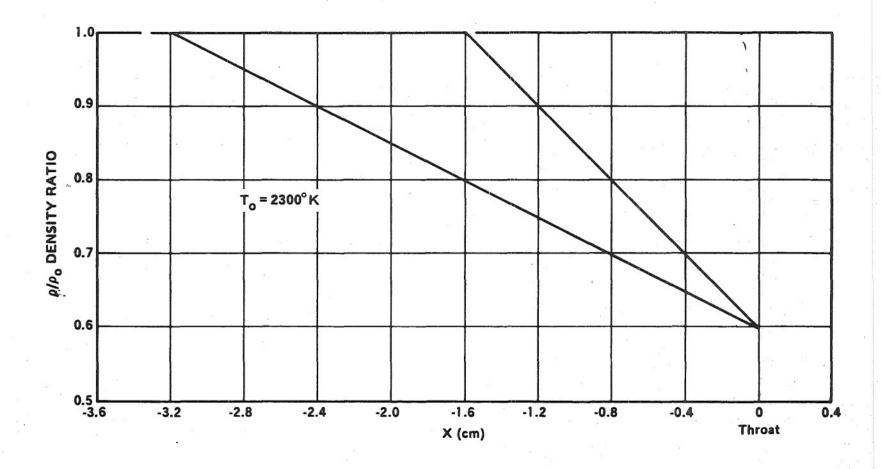


FIGURE 4.7. NOZZLE SUBSONIC SECTION DENSITY DISTRIBUTION

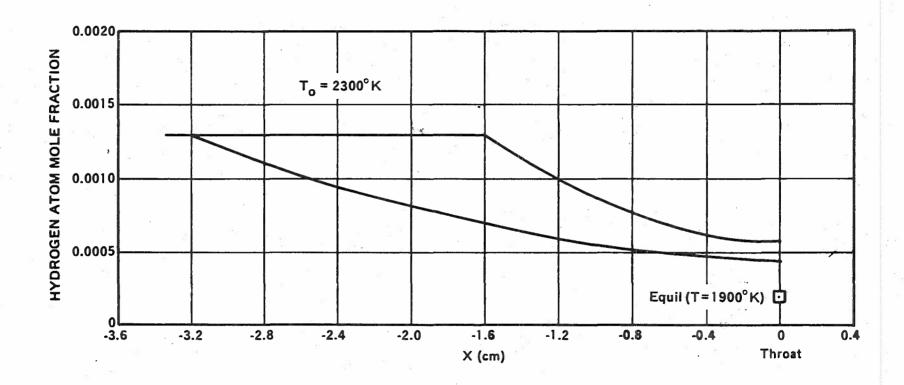


FIGURE 4.8. H-ATOM CONCENTRATION IN SUBSONIC NOZZLE REGION

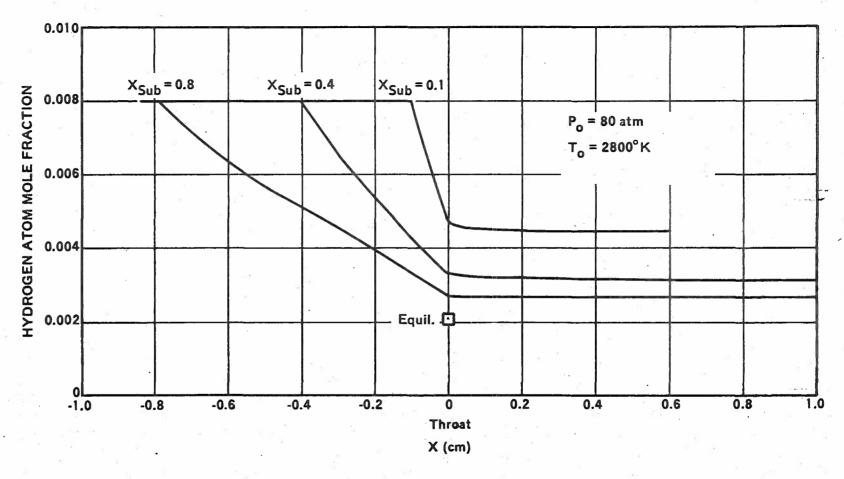


FIGURE 4.9. H-ATOM CONCENTRATION

4.1 Method of Characteristics Analysis

The harmonic oscillator or anharmonic oscillator analysis is not capable of generating a smooth nozzle contour that will produce a nozzle exit flow that is uniform and shock free. For this reason, a method of characteristics computer program that takes into account the nonequilibrium nature of the flow was developed. This program is a combination of a method of characteristics program and the harmonic oscillator analysis program. The modifications required to the method of characteristics program were to use the frozen sound speed defined as

$$a_f^{\ 2} = \frac{p}{\rho} \, \left(\frac{C_p}{C_p - R} \right)$$
 where p = pressure $$R$ = gas constant ρ = density$

where C are comprised of only the rotational and translational contributions to the specific heat. Also, the velocity u and temperature T that appear in the compatibility equations are determined from the solution of the energy and momentum equations which are:

$$\frac{u}{dx} \frac{du}{dx} = -\frac{dh}{dx}$$

$$\frac{u}{R} \frac{dh}{dx} = \left(-\frac{\dot{c}}{E_{V}} + \frac{C_{p}}{R} \right)^{T_{U}} \frac{d\ln\rho}{dx} \left(\frac{C_{p}}{R} - 1 \right)$$

$$u^{C_{p}} \frac{dT}{dx} = \frac{u}{dx} \frac{dh}{dx} - \frac{R}{E_{V}}$$

The terms involving $E_{\rm V}$ represent the effect of vibrational nonequilibrium and are the sum of the compounds from each specie. The quantity h is the static enthalpy as

$$h = \Sigma X_{1} E_{V1} + C_{D}T$$

where X_{i} is the mole fraction of specie i.

The compatability equations for the characteristics solution also contains a heat release term which accounts for nonisentropic vibrational energy changes. The compatibility equations are

$$A dp \pm d\theta + B \begin{Bmatrix} d\eta \\ d\xi \end{Bmatrix} + \frac{a_f}{u^2 C_p T} \sum_{\chi_i} \hat{\omega}_i \begin{Bmatrix} d\eta \\ d\xi \end{Bmatrix}$$

where
$$A = \left(\frac{u^2}{a^2} - 1\right)^{1/2} / \rho u^2$$

$$B = \frac{a_f \frac{\sin \theta}{uy}}{a_f \frac{\sin \theta}{uy}}$$

The increments $d\eta$ and $d\xi$ are shown in Figure 4.10 and are the distances along the uprunning and downrunning characteristics, respectively.

The computer program has been used to determine a preliminary contour for the Mod VIII nozzle contour. The area distribution from the harmonic oscillator program was used to determine an initial estimate for the contour. This was then used to determine the nozzle slope as a function of x. This initial flow angle distribution is shown in Figure 4.11. This initial contour produced the wave diagram in Figure 4.12 and caused the program to terminate at about x = .125 cm on the centerline. This was due to the formation of a shock. The contour was then subsequently modified to that designated as A in Figure 4.11. The wave diagram for this case is shown in Figure 4.13. As can be seen from the figure, a weaker shock is now forming due to too rapid turning of the nozzle wall back to zero. The shock is weaker than the initial contour but is still present. Further reduction in the nozzle wall angle would eventually result in completely eliminating the shock.

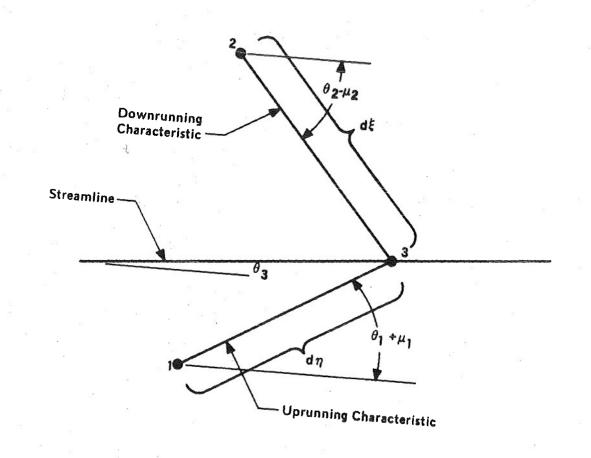


Figure 4.10. Characteristic Mesh Details.

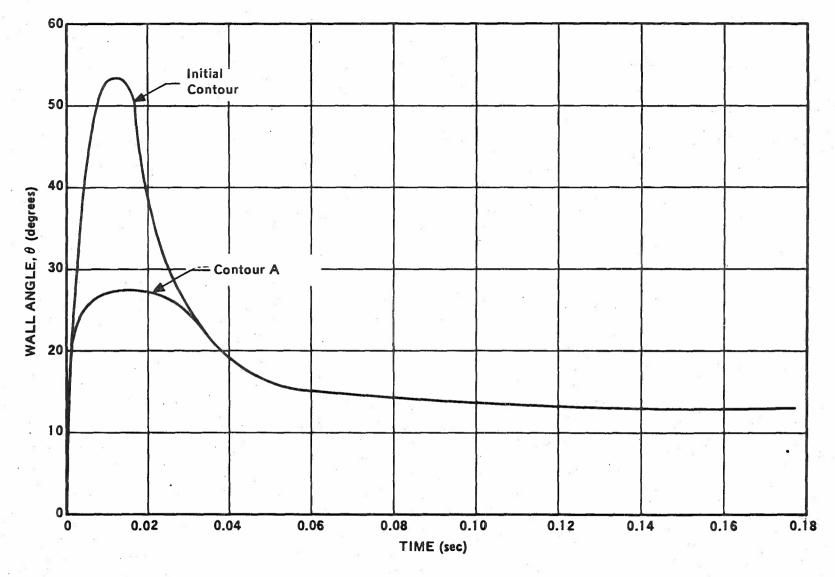


Figure 4.11. Mod VIII Nozzle Contour, θ Versus

Figure 4.12. Nozzle Contour Test Case.

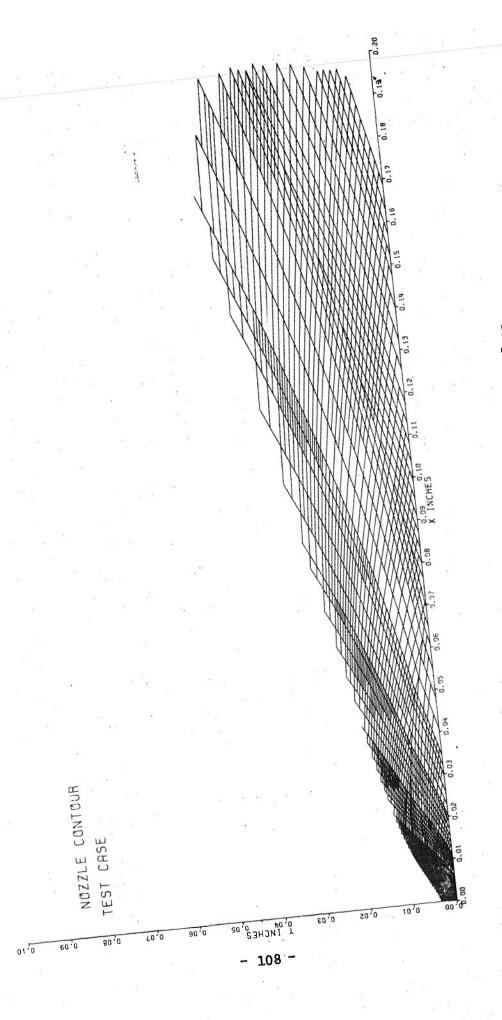


Figure 4.13. Nozzle Contour Test Case.

4.2 Boundary Layer Analysis

The inviscid nozzle contour as determined from the method of characteristics analysis has to be corrected for viscous effects due to boundary layer growth. To accomplish this, a boundary layer integral technique was used. Due to the extremely large flow gradients in the throat region of a quick freeze gas dynamic laser nozzle, an integral approach requires considerably less computer time than a finite difference technique. The integral technique employs the momentum integral equation which is

$$\frac{d\theta}{dx} + \frac{\theta}{\rho_e U^2} \frac{d}{dx} (\rho_e U^2_e) + \frac{\delta^*}{U_e} \frac{dU_e}{dx}$$

$$= \frac{\tau_{\tilde{u}}}{\rho_e U^2_e}$$

where U = boundary layer edge velocity

 $\delta *$ = displacement thickness

 θ = momentum thickness

 δ = boundary layer thickness

 ρ_e = boundary layer edge density

 τ_{m} = wall shear stress

There are three unknowns in this equation, θ , δ * and τ_{ω} . The quantities δ * and θ are defined as

$$\delta * = \delta \int_{0}^{1} \left(1 - \frac{\rho}{\rho_{e}} \frac{u}{u_{e}} \right) d\eta$$

$$\theta = \delta \int_{0}^{1} \frac{\rho u}{\rho_{e} u_{e}} \left(1 - \frac{u}{u_{e}} \right) d\eta$$

The evaluation of δ * and θ require specifying the velocity profile and the density profile. The velocity profile selected is a power law profile of the form

$$\frac{u}{u_e} = (\eta)^{1/N}$$

The quantity N is determined using the empirical data of Reference 40. A curve fit of N is given in Figure 4.14. As can be seen in the figure, the value of θ and subsequently $R_{e\theta}$ determines the value of N which in turn determines the velocity profile. This approach is simpler than resorting to the use of an additional equation to determine N and since the real test of the method must rely on experimental data, for the present this technique is probably sufficient.

The density profile and shear stress were determined using the results of Reference 41. The work of Reference 41 reviewed the literature on compressible turbulent boundary layers. One of the results of Reference 41 was to develop an expression for the skin friction coefficient \mathbf{C}_f from which the wall shear stress \mathbf{T}_ω can be determined. The skin friction relation developed has the form

$$C_f = \frac{F_{rd}^{R}e_{\theta}}{F_{c}}$$

where

$$F_{r_d} = \left(\frac{T_w}{T_\infty}\right)^{-.702} \left(\frac{T_{oe}}{T_w}\right)^{.772}$$

$$F_{c} = \left[\int_{0}^{1} \left(\frac{\rho}{\rho_{\infty}} \right)^{1/2} d(u/u_{e}) \right]^{-2}$$

The density distribution used in the above relation is

$$\frac{\rho}{\rho_{\infty}} = \left[\frac{T_{W}}{T_{\infty}} + 1 + \frac{1}{2} r (\gamma - 1) M_{\infty}^{2} - \frac{T_{W}}{T_{\infty}} \frac{u}{u_{c}} - \frac{r}{2} (\gamma - 1) M_{\infty}^{2} \left(\frac{u}{u_{c}} \right)^{2} \right]^{-1}$$

where r is the recovery factor and has a value of about .89. A plot of F_c^c versus $F_{r\theta}$ $R_{e\theta}$ taken from reference 41 shown in Figure 4.15. These values are also tabulated in the reference and are given in Table 4.2. The two sets of data, however, do not agree. In order to eliminate this problem a test case for

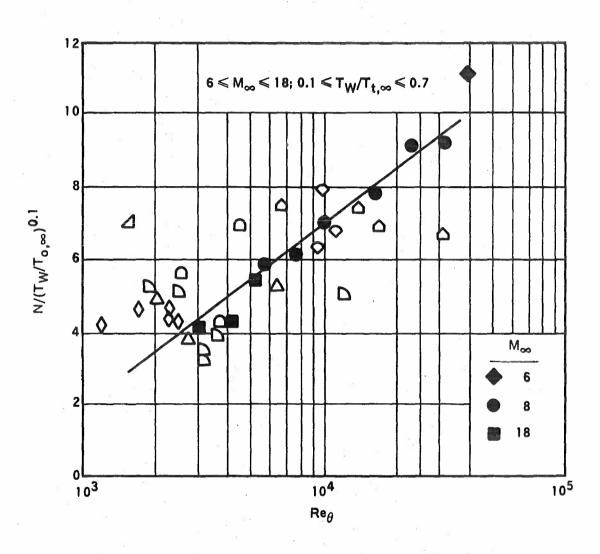


Figure 4.14 Variation of N with Re_{θ} for Axisymmetric Air and Nitrogen Nozzie Data.

Figure 4.15 Comparison Between Theoretical and Experimental F_cC_f Versus $F_{Rd}R_{e\theta}$. \Box , Experiments, Adiabatic; O, Experiments with Heat Transfer; –, Theory.

BOUNDARY LAYER CHARACTERISTICS *

TABLE 4.2

$^{\mathbf{F}}\mathbf{c^{c}}\mathbf{f}$	$^{F}R_{d}^{R}_{e\theta}$	$^{\mathrm{F}}\mathrm{c}^{\mathrm{c}}\mathrm{f}$	$^{\mathrm{F}}\mathrm{R_{d}}^{\mathrm{R}}\mathrm{e}\theta$
0.0010	2.878×10^7	0.0060	233.0
0.0015	3.955 x 10 ⁵	0.0065	177•6
0.0020	5.425 x 10 ⁴	0.0070	140 • 4
0.0025	1.386×10^4	0.0075	114.4
0.0030	5030	0.0080	95 • 62
0.0035	2283	0.0085	92•49
0.0040	1208	0.0090	70.91
0.0045	716 • 0	0.0095	62.55
0.0050	462 • 3	0.0100	55.87
0.0055	319 • 4	0.0105	50•46

^{*}Reference 41 - "The Drag of a Compressible Turbulent Boundary Layer on a Smooth Flat Plat with and without Heat Transfer," D.B. Spalding and S.W. Chi, Journal of Fluid Mechanics, Vol. 18, Part I, January 1964.

air was run that employed the analysis of reference 42. The results of this case are shown in Figure 4.16. The curve fit for $C_{\rm f}$ employed in the present technique that gave the best agreement with the method of reference 42 was

$$C_f = .0462 \left(F_{rd} R_{e\theta} \right) - .283826 / F_c.$$

The technique was then used to determine the boundary layer properties for the Mod VIII nozzle contour. The results of this calculation are shown in Figure 4.17. The results indicate that the boundary layer displacement thickness is relatively small and represents only a small percentage of the nozzle exit flow area. It must be remembered that these results are valid only within the limits of the analytical model employed but represent a good first estimate of the viscous effects.

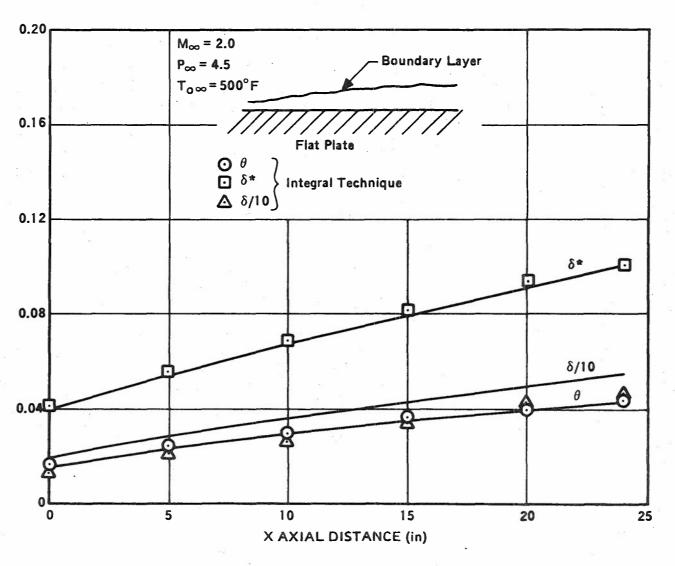


Figure 4.16 Test Problem Flow Configuration.

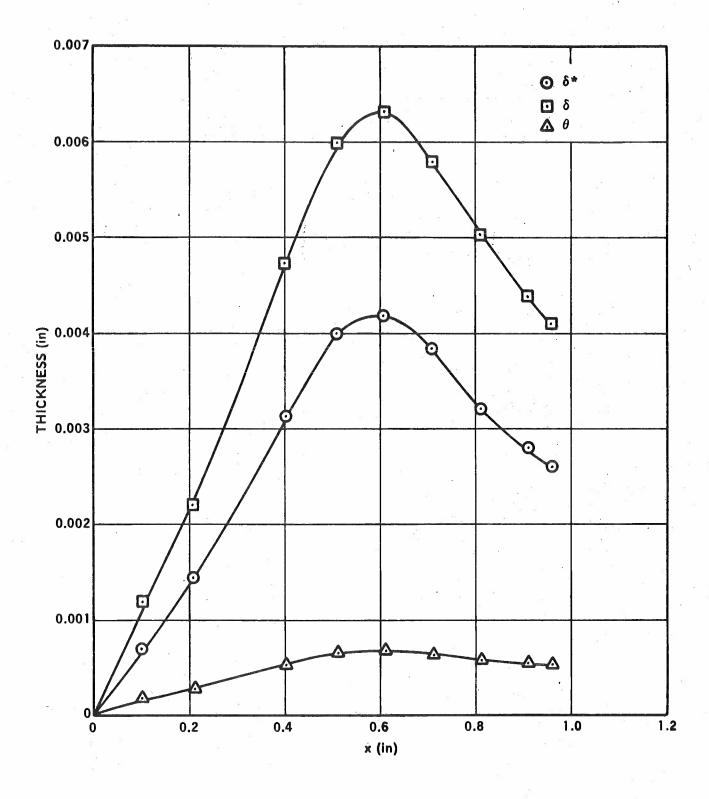


Figure 4.17 Nozzle Boundary Layer Properties.

4.3 Propellant Impurity Effects

During the solid propellant gas generator development, it was anticipated that various species other than H₂/HCl and their atoms might also be produced in small concentrations. A list of potential contaminant species is given in Table 4.3. The effect of these impurity species on predicted laser performance is important, since it could well determine the technical direction of any further propellant development work undertaken.

Most of the impurity species listed in Table 4.3 are metallic halides, oxides, or hydroxides with relatively high boiling points. Because they are not easy to produce in the gas phase, there is no experimental data on their efficiency of vibrational relaxation on HCl or H_2 . The fundamental vibrational frequencies of these refractory species are quite small compared to those for H_2 and HCl. Therefore, it is anticipated that these impurities will not act as $V \rightarrow V$ deactivators. On the other hand, the efficiency of a $V \rightarrow T$ deactivation might reasonably be expected to vary between a probability per collision, P, of 10^{-1} to 10^{-5} . For illustration, we will assume a value of 10^{-3} which is also similar to the probability for the HCl/HCl V \rightarrow T process at $T = 300^{\circ}$ K. Therefore, for impurity specie concentrations <<[HCl], the effect of the contaminant should be insignificant. To allow for the possibility of a larger probability per collision, a more conservative estimate of [impurity] $\stackrel{<}{\sim} 0.1$ mole percent is suggested. This appears to be generally within the capability of the current propellant technology.

Of the species listed in Table 4.3 only two possess fundamental frequencies close to those of HCl, e.g., $\rm H_2O$, and $\rm N_2$. Some data do exist on the vibrational coupling of $\rm N_2$ to HCl, so it was decided to do a more quantitative estimate of the effect of these species on laser performance.

 $\rm H_2^{0}$ is a special case. It is known that $\rm H_2^{0}$ deactivates HCl very rapidly 35 , although no temperature dependence is available. If the large cross-section for deactivation persists to elevated temperature, the effect of $\rm H_2^{0}$ on laser performance would be very similar to that of H-atoms. For this reason, the concentration of $\rm H_2^{0}$ 0 was kept to a low value in the propellant formulations selected.

TABLE 4.3

POTENTIAL GAS CONTAMINANTS

SPECIES		v	IBRATIONAL FREQUENCIES 33,3
			(cm ⁻¹)
BF			1410
BF ₂			1120, 470, 1450
BF ₃		98	881, 670, 1463, 480
ZrC1			437
zrc1 ₂			346, 92, 461
ZrCl ₃			490, 185, 333, 131
ZrC1 ₄			376, 99, 418, 110
A1C1			480
AlCl ₂			430, 210, 540
AlCl ₃			371, 146, 185, 610
A10C1			450, 350, 900
A1 ₂ 0			715, 238, 994
А10Н			1000, 1300, <u>3600</u>
H ₂ O	; at a		<u>3652</u> , 1595, <u>3756</u>
N_2			<u>2331</u>
	BF BF ₂ BF ₃ ZrC1 ZrC1 ₂ ZrC1 ₃ ZrC1 ₄ A1C1 A1C1 ₂ A1C1 ₃ A1OC1 A1 ₂ O A1OH H ₂ O	BF BF ₂ BF ₃ ZrC1 ZrC1 ₂ ZrC1 ₃ ZrC1 ₄ A1C1 A1C1 ₂ A1C1 ₃ A1OC1 A1 ₂ O A1OH H ₂ O	BF BF ₂ BF ₃ ZrC1 ZrC1 ₂ ZrC1 ₃ ZrC1 ₄ AlC1 AlC1 ₂ AlC1 ₃ AlOC1 Al ₂ 0 AlOH H ₂ 0

Vibrational rate constants have been assembled for $N_2/H_2/HC1$ molecular systems from the available data 7,36,38 and are listed in Table 4.4. Figure 4.18 shows the result of a calculation of available energy as a function of N_2 mole fraction for the assumed base line condition. The results of this calculation indicate that 2-3% of N_2 is acceptable without appreciably affecting performance. A major part of the performance change is due to the increased effective mass of gas as N_2 is added. Some of the loss can be recovered through improved gas yields made possible by permitting N_2 formation. Thus, the gas generators selected are sufficient for good performance; however, a more detailed assessment of N_2 must await better kinetic data or shock tunnel gain experiments.

4.4 Power Extraction

The initial feasibility calculations were all based on the concept of "available energy," i.e., the amount of potentially abstractable energy existing in HCl at some selected position in the cavity. To provide a more realistic analytic capability, power extraction models based on the "gain equals loss" approximation were developed and became operational for both the H.O. and A.H.O. codes. These analyses for power extraction are based on the Rigrod approach ³⁵ employing the concept that lasing can occur when the gain exceeds the single pass cavity loss. The photon amplification condition in the laser cavity can be expressed as:

$$e^{(2g_v^{v-1} \ell)} \ge [R(1-L_c)]^{-1}$$
 (32)

where g_V^{v-1} is the gain for vibrational transition $v \rightarrow v-1$, ℓ is the cavity length between mirrors of "effective" reflectivity R (which includes all useless losses), and L_c is the cavity output coupling. Equation (32) can be rewritten as the lasing condition

$$g_{v}^{v-1} \ge -\frac{\ln [R (1 - L_{c})]}{2 l}$$
 (33)

The gain is determined from the level populations which are in turn determined by the kinetic equations.

The effect of stimulated emission is incorporated into the differential equations for the level populations by adding to the kinetic terms the effects of lasing to and from the particular level; namely

TABLE 4.4

VIBRATIONAL KINETICS FOR N2/H2/HC1 SYSTEMS

(Rate constants in units of cm^3 - particles⁻¹ sec^{-1})

$$V \rightarrow T$$
 Process (26)

$$\text{HC1*} + \text{N}_2 \rightarrow \text{HC1} + \text{N}_2$$

 $k_f = 5.53 \times 10^{-14} \text{ T exp (-109.62 T}^{-1/3}) \text{ cm}^3 - \text{molecule}^{-1} - \text{sec}^{-1}$

$$H_2^* + N_2^- + H_2^- + N_2^-$$
 (27)
 $k_f = 5.44 \times 10^{-10} \exp (-114.91 \text{ T}^{-1/3})$

$$N_2^* + HC1 \rightarrow N_2 + HC1$$
 (28)
(not known, but unimportant when H_2 HC1)

$$N_2^* + H_2 \rightarrow N_2 + H_2$$
 (29)
 $k_f = 1.78 \times 10^{-7} \exp (-172.69 \text{ T}^{-1/3})$

$$N_{2}^{*} + N_{2} \rightarrow N_{2} + N_{2}$$

$$k_{f} = 8.53 \times 10^{-7} \exp \left(-273.10 \text{ T}^{-1/3}\right)$$
(30)

V → V Process

$$\text{HC1*} + \text{N}_2 \rightarrow \text{HC1} + \text{N}_2^*$$

$$k_f = 8.10 \times 10^{-17} \text{ T}$$
(31)

(The H_2/N_2 V \rightarrow V process is considered as unimportant and thus neglected)

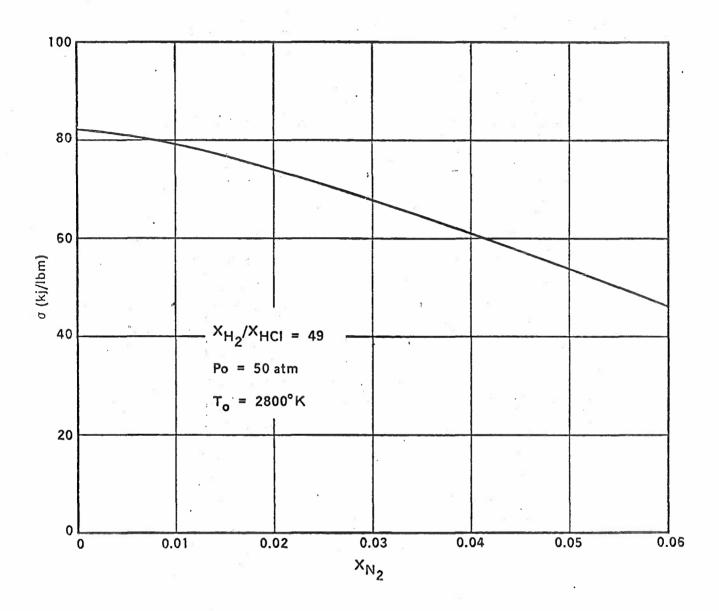


Figure 4.18. Effect of N₂ on Performance.

$$\frac{d n_{v}}{dt} = \frac{d n_{v}}{dt} + g_{v+1}^{v} F_{v+1}^{v} - g_{v}^{v-1} F_{v}^{v-1}$$
(34)

 $F_{\rm v}^{\rm v}$ is the cavity flux determined from the imposition of Equation (33). This model assumes rotational equilibrium. For a given vibrational level, only the rotational line of largest gain is allowed to laser, but simultaneous lasing on different vibrational levels is allowed. In the GDL case, the total power extracted is calculated from the expression:

$$E = \sum_{v} \frac{1}{2^{*}} \int_{0}^{+} h v_{v}^{v} - 1 F_{v}^{v} - 1 A/A^{*} \frac{L_{c}}{2^{*}} dx$$
 (35)

where ρ^* and u^* are the gas density and velocity at the throat, and $h\nu_V^{V-1}$ is the energy of the transition. The results of a calculation using this model are illustrated in Figures 4.19 - 4.20.

Figure 4.19 shows the calculated small signal gain, $g_{_{\rm O}}$, determined for the given conditions from the A.H.O. kinetics model. The maximum gain occurring on vibrational levels v=1-7 is shown as a function of distance downstream of the nozzle throat. Assuming an optical cavity condition of R=0.98 and $L_{_{\rm C}}=0.2$ a critical gain can be calculated from Equation (33) and is shown on Figure 4.19 as the dashed line. Lasing should occur for $x\geq 10$ cm downstream of the throat.

Figure 4.20 shows the same calculation as Figure 4.19 including power extraction at distances greater than 15 cm downstream. The effect of lasing is seen to surpress the gain to values around the critical value. Only a few gain histories are plotted on Figure 4.20 for clarity. The observed oscillatory behavior is due to rotational line shifting on a particular vibrational level.

On Figure 4.21 the cavity fluxes F_V^{V-1} for a number of vibration levels and the total flux, are given for the same conditions. From this result the total output power is calculated to be 54 kj/lb with most of this power radiated from the first three vibrational levels. As is seen in Figure 4.21 the cavity was allowed to extend arbitrarily to 300 cm downstream.

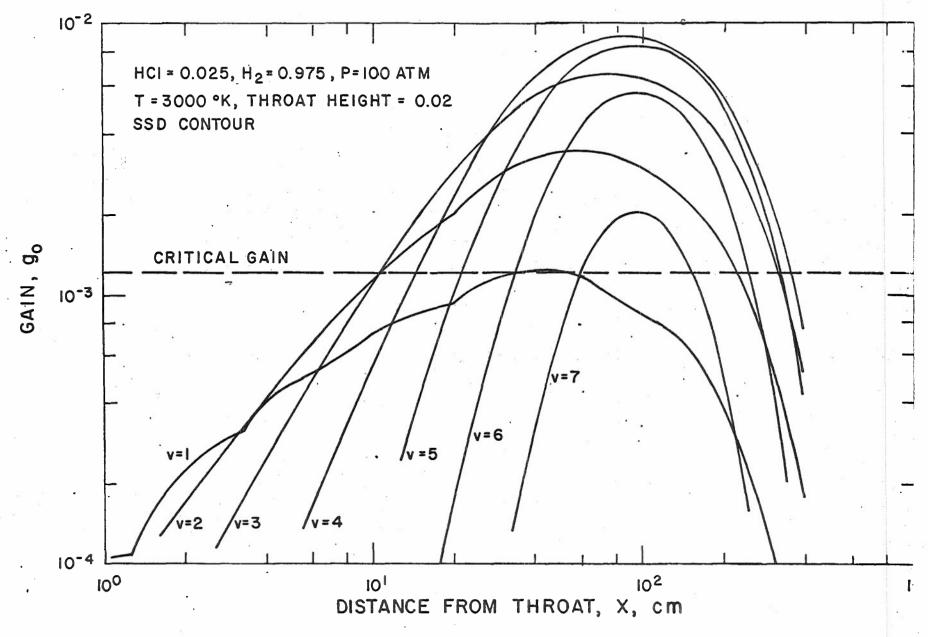


Figure 4.19. Maximum small signal gain for the v=1 to 7 vibrational levels of HCl. Critical gain is calculated for ℓ =100 cm, R=0.98, and L=0.2.

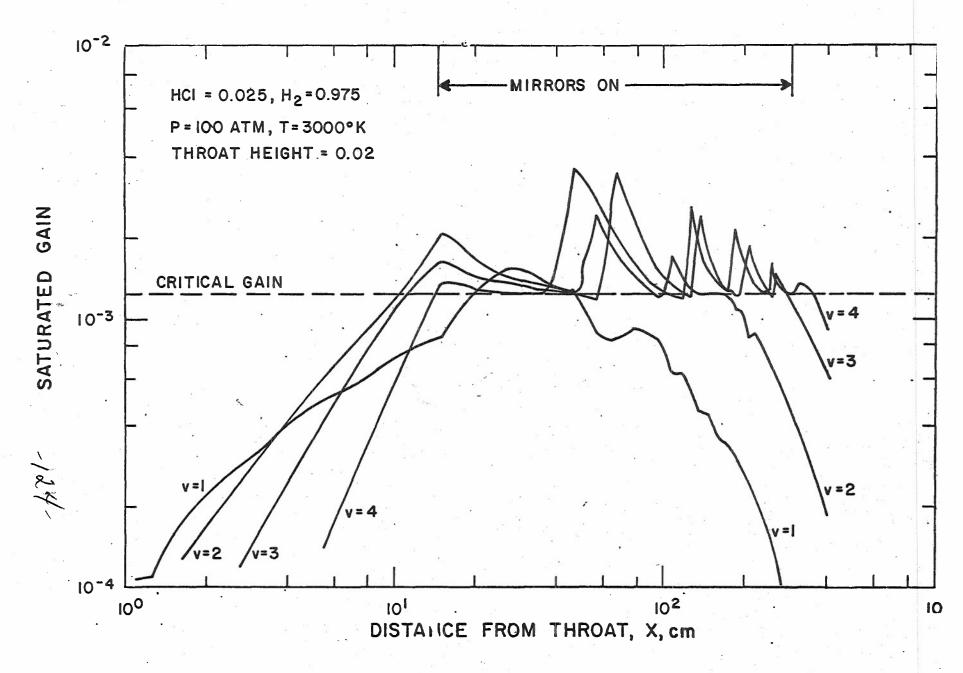


Figure 4.20. Saturated gain for v=1 to 4 vibrational levels of HC1. Cavity conditions: l=100 cm, R=0.98, and L=0.2.

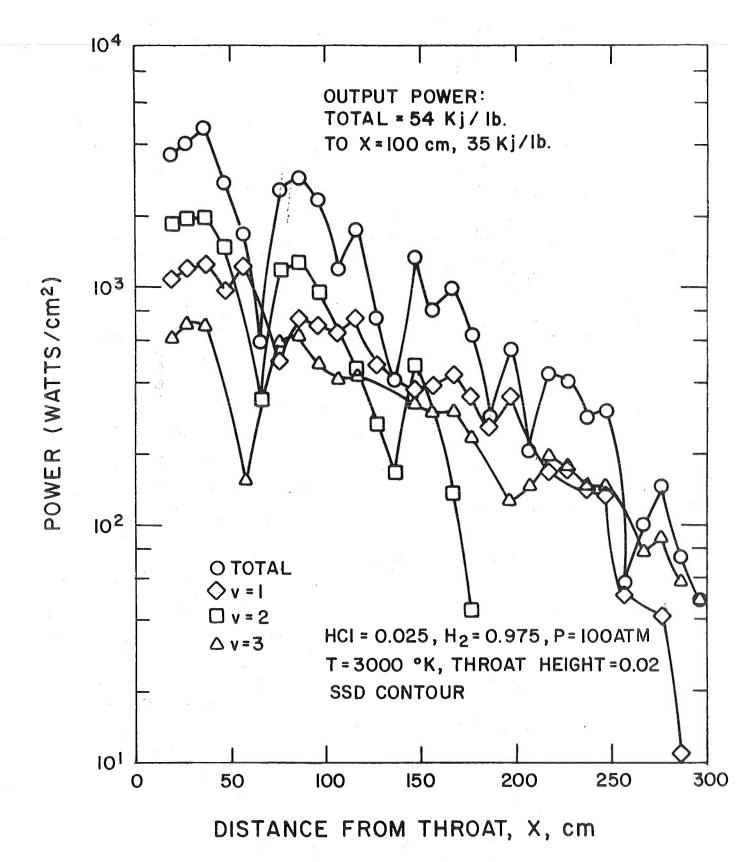


Figure 4.21. Calculated intracavity flux for v = 1 to 3 vibrational transition of HCl.

Also shown is the total flux, i.e., the sum of all levels, including those not shown.

However, the bulk of the energy is extracted early in the flow. For example, at X = 100 cm the output power is already 35 kj/lb. It is emphasized that the calculation presented in the figures is only illustrative, and no attempt has yet been made to optimize these results.

This power extraction routine can also be used with the H.O. model. In this situation, there is only a single vibrational level and the calculation is correspondingly simpler. The calculated output power in the H.O. approximation for the same assumed conditions as were shown previously for the A.H.O. case is given in Figure 4.22. A comparison of Figure 4.22 with the A.H.O. result indicates that the H.O. calculation yields about twice the extracted power as the A.H.O. result. This illustrates one severe limitation of this simple H.O. model of a vibration transfer laser.

4.5 Advanced H₂-HCl Gas Dynamic Laser System Characteristics

The end objective of the work performed is to develop a high energy laser device. The work performed during this phase is now sufficient to determine the weight and volume characteristics of a high energy laser system. A typical HEL device size was selected. The run time for the device consists of 20-three second bursts for a total run time of 60 seconds. The weight and volume of the laser system depends on the propellant weight yield and the specific power. The specific power used to estimate the system size is 100 kj/lb. This value represents an upper limit which depends on the effect of H-atoms on specific power. The H-atom problem will be addressed in much greater detail in the next year.

The laser device has been divided into three basic components which are: (1) the combustor; (2) the nozzle bank; and (3) the cavity and diffuser. A sketch of the device is shown in Figure 4.23. For the purpose of weight and volume estimates, the combustor is assumed to be comprised of 1" thick steel, the nozzle bank is assumed to be solid steel and the cavity and diffuser are assumed to consist of a 3" thermal skin. The thermal skin is comprised of a .25 inch thick steel plate outer shell and a .25 inch thick steel plate inner shell as shown in the inset of Figure 4.23.

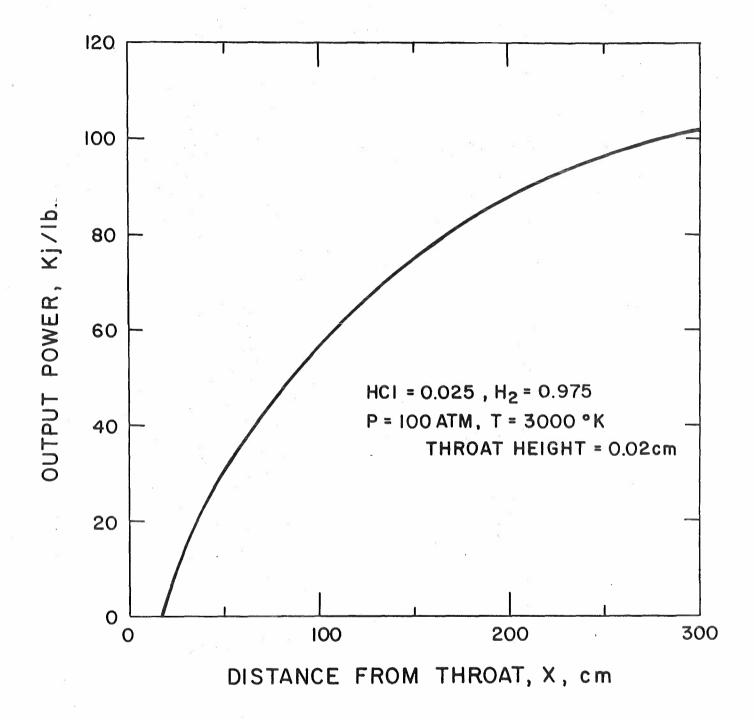
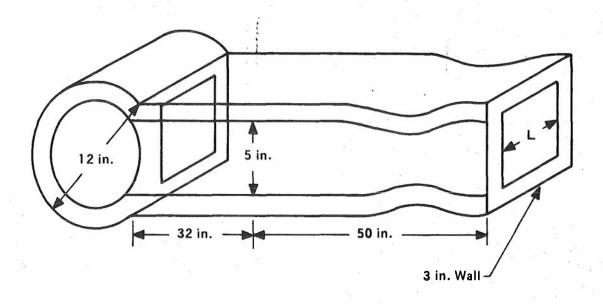


Figure 4.22. Calculated output power for H.O. calculation. Cavity conditions $\ell = 100$ cm, R = 0.98 and L = 0.2



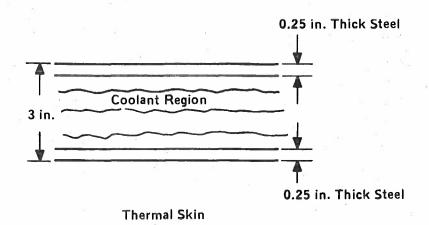


Figure 4.23. Laser Device

The weight and volume of the components is shown in Table 4.5 for each of the three candidate propellants. The overall envelope of the device is obtained using a packing factor of 2. The weight and volume of the individual gas generators are shown in Table 4.6. The gas generator weight and volume estimates include a filter system for propellant candidates 2 and 3. A more detailed breakdown of the gas generators is shown in Table 4.7.

The overall system weights and volumes are shown in Table 4.8. Also shown in Table 4.8 are the specific weight power of the laser devices, the specific volumetric power and the specific nozzle exit area power. The best device performance is obtained from the HDB propellant system, whereas the other two systems have about the same system weight and volume. Depending on whether the application of the laser is weight or volume limited either the device specific weight or volume parameter should be selected as the basis for comparison to other laser systems.

TABLE 4.5

WEIGHT AND VOLUME OF COMPONENTS

	PROP. 1 and	3 (L=2 ft)	PROP. 2 (L	=2.5 ft)
ITEM	WEIGHT (1b)	VOLUME(ft ³)	WEIGHT(1b)	VOLUME(ft ³)
COMBUSTOR	330	1.6	420	2
NOZZLE	378	0.73	500	1
CAVITY AND DIFFUSER	836	12.0	1,040	16
TOTAL	1,544	14.33	1,960	19
PF = 2		28.66		38

TABLE 4.6

WEIGHT AND VOLUME OF GAS GENERATOR

P	ROPELLANT	WEIGHT (1b)	VOLUME (ft ³)
	1	4,900 (245)	45 (2.24)
	2	10,100 (506)	88 (4.4)
	3	11,580 (579)	104 (5.2)
		() DENOTES PER SH	TOT

		(H ₂ +HC1)			Derived Propellant System Characteristics						
Comp	Composition, Wt %	H ₂ /HC1 Molar	<u>Yi</u>	eld Vol %	T _f ,	Propellant Wt (1b)	Weight of Case	Total Wt (1b)	Motor Vol(ft ³)	Motor Dia (in)	Motor Length (in)
1	11.821% H ₂ WO ₄ 33.475% Zr 38.704% HDB 16% AP	33.3	10.8	95.9	2808	139	121	260	2.8	15	26
2	10.482% NH ₄ NO ₃ 47.132% ZrH ₂ 38.186% Zr(OH) ₄ 4% AP	38.2	3.84	89.5	2799	391	234	625	5.5	15 ,	51
3	40.879% 42.968% ZrH ₂ 13.153% Zr(OH) ₄	32.6	2.59	98.0	2799	579	275	854	6.5	15	61

TABLE 4.8

WEIGHT AND VOLUME OF SYSTEM

PROPELLANT	WEIGHT(ft)	VOLUME (ft 3)	kj/lb	kj/ft ³	kj/in ²
1	6,444	73	0.062	5.47	4.7
2	12,060	126	0.033	3.17	3.2
3	13,124	132	0.030	3.03	4.7

5.0 CONCLUSIONS

The results of the Phase I study indicated that an advanced H2-HCl gas dynamic laser was indeed feasible. The objective of the current phase, Phase II, was to experimentally verify the V-V and V-T kinetics for the important H2-HCl processes and to develop solid propellant gas generators that could be used to generate the required gases at attractive weight yields. The results of the current phase have verified that the kinetics are capable with operation of a H_2 -HCl laser. In addition, the gas generator development program has succeeded in producing the required gases at weight yields that make the system attractive for Army field application. The potential of even higher gas yields that would make the system even more attractive is also possible but requires additional propellant development. The fundamental remaining problem area is the effect of H-atoms on performance. This problem area is currently under investigation and represents the only remaining barrier that prevents the development of an advanced H2-HC1 gas dynamic laser. It is believed that this problem may have to be resolved by the use of an adequate H-atom scavenging technique.

REFERENCES

- "Final Report for Phase I of the NRML Laser," ARC Report No. 47-5629, 1. Contract No. DAAHO1-74-C-0611, Atlantic Research Corporation, Alexandria, Virginia 22311, February 1975.
- 2. W. D. Breshears and P. F. Bird, J. Chem, Phys. 52, 999 (1970).
- 3. C. T. Bowman and D. J. Seery, J. Chem. Phys. 50, 1904 (1969).
- B. M. Hopkins, H. L. Chen and R. Sharma, J. Chem. Phys. 59, 5758 (1973).
- 5. P. F. Zittel and C. B. Moore, J. Chem. Phys. 59, 6636 (1973).
- 6. B. M. Hopkins and H. L. Chen, J. Chem. Phys. 57, 3161 (1972).
- 7. J. F. Bott and N. Cohen, J. Chem. Phys. 63, 1518 (1975).
- 8. M. M. Audibert, C. Joffries and J. Ducuing, Chem. Phys. Lett. 25, 158 (1974).
- 9.
- J. H. Kiefer and R. W. Lutz, J. Chem, Phys. 44, 668 (1966).
 J. E. Dove and H. Tietelbaum, J. Chem., Chemical Physics (1974) to be supplied. 9a.
- 10. The authors are indebted to Captain Charles Holmes of the Air Force Weapons Laboratory, Kirtland Air Force Base, Albuquerque, N.M., for his help in the theoretical calculation.
- 11. J. C. Decius, J. Chem. Phys. 32, 1262 (1960).
- 12. R. V. Steele, Jr. and C. B. Moore, J. Chem. Phys. 60, 279 (1974).
- 13. D. J. Seery, J. Chem. Phys. 58, 1796 (1973).
- 14. R. N. Schwartz, Z. I. Slawsky, and K. R. Herzfeld, J. Chem, Phys. 20, 1591 (1952).
- R. L. Taylor, M. Camac and R. M. Feinberg, "Measurements of Vibration-15. Vibration Coupling in Gas Mixtures," Proceedings of the Eleventh Symposium (International) on Combustion, 14 - 20 August 1966, The Combustion Institute, Pittsburg, Pennsylvania, 1967, pp. 49-65.
- V. Sato, S. Tsuchiya and K. Kuratani, J. Chem. Phys. 50, 1911, (1969) plus 16. Erratum, J. Chem. Phys. 53, 1304 (1970).
- 17. K. Hancock, private communication.
- 18. D. Arnoldi and J. Wolfrum, Chem. Phys. Lett. 24, 234 (1974).
- 19. R. F. Heidner III and J. F. Bott, J. Chem. Phys, 63, 1810 (1975).
- 20. R. D. H. Brown, G. P. Glass and I. W. M. Smith, results to be published.
- 21. R. L. Wilkens, J. Chem. Phys. 63, 534 (1975).
- M. A. A. Clyne and D. H. Stedman, Trans. Far. Soc. <u>62</u>, 2164 (1966). 22.

APPENDIX A

Theoretical values of flame temperature and combustion product compositions are tabulated here for selected propellants as a function of propellant compositions. Values of flame temperature are for adiabatic combustion with chemical equilibrium at 750 psia (generally). Values for combustion product compositions are for throat conditions, assuming expansion of only the gaseous species that exist at chamber conditions. The expansion is assumed to be isentropic, with chemical equilibrium throughout. The calculations are computerized results, using JANAF thermochemical data.

It should be noted that ${\rm Al}_2{}^0{}_3({\rm c})$ and BN(c) appear as combustion products at throat conditions of all propellants containing Al (or AlH $_3$) and HDB, respectively. This phenomenon has been discussed in the text and is essentially the result of condensation reactions of gaseous species during expansion. The effect is not large, amounting to $\sim\!2\%$ by weight of combustion products appearing as condensed at the throat.

In the tables, the gas composition is expressed in terms of mol/100 gm of propellant. Some of the tables show the column entitled "Eng" which refers to the total moles of gaseous combustion products at the throat. The molar ratio $\rm H_2/HC1$ is seen to be quotient of the entry in Column 7 ($\rm H_2$) divided by that in Column 8 (HC1). The yield ($\rm H_2$ + HC1) is seen to be the sum of the products

$$2.016$$
 (Col. 7) + 36.461 (Col. 8)

for the expression for weight percent. The yield (${\rm H_2}$ + HCl) for volume percent is the molar content of (${\rm H_2}$ + HCl) in the combustion products, and is

$$(H_2 + HC1)/\Sigma ng.$$

For most of the tables, this quantity is derived.

APPENDIX A INDEX OF TABLES (1)

Propellant System

Fue1 ZrH2	Zr(OH) ₄	Chlorides NH ₄ C1	Number
21112	LF(OH) ₄	1 140 1.1	۸1
+	i i	i '	A1
i	(NH ₄ C10 ₄	A2
	(CH ₂) _n	1 '	A3
10 m ²⁴	20	PL PAGE	. A 4
	(84)	1 ' ' 1	A5
Zr	1.5	1 ' '	. A6
•	HDB	1	A7
(V)	8	NH ₄ C10 ₄	8A
**	HDB/CH ₂		A9
A1,	A1(OH) ₃	NH ₄ C1	A10
r		NH ₄ C1O ₄	A11
	(CH ₂) _n	1	A12
		1 ' ' '	A13
	нтрв		A14
A1H ₂	A1 (OH) 2	NH,C1	A15
3		4	A15
13	3 2 n		
ZrH	Zr(OH),	NH, C1	A16
2	4	1 '	A17
	FE 75		A18
, ,	(CH ₂)	' '	A19
* *	2'n	1 -	A20
i.	2)	1	A21
8	HDB		A22
7r			A23
			A24
	UDB /2~ (OU)		
t _d	105/2F(OR)4	21014	A25
1			
	10 10	• V	72
	. ~ =		
121	ľ	8	
	Zr A1 A1H ₃ ZrH ₂	HDB HDB/CH ₂ A1 (OH) ₃ (CH ₂) _n HTPB A1 (OH) ₃ A1 (OH) ₃ /(CH ₂) _n ZrH ₂ Zr(OH) ₄ (CH ₂) _n HDB	HTPB

APPENDIX A

INDEX OF TABLES (1)

(Continued - Page 2)

Propellant System

Oxidizer	Fuel	Coolant	Chlorides	Table Number
N2H5NO3	A1	Al(OH) ₃	NH ₄ C1	A26
2 3 3	×		N ₂ H ₅ C1	A27
			NH ₄ C1O ₄	À28
e * • * *		(CH ₂) _n	NH ₂ C1	A29
0.		2 11	N ₂ H ₅ C1	A30
¥ 2		: ** ** .	NH ₄ C1O ₄	A31
		HDB	NH ₄ C1	A32
	A1H ₃	A1 (OH) 3	NH ₄ C1	A33
		(CH ₂) _n	7	A34
		A1 (OH) 3	NH ₄ C1O ₄	A35
		Zr(OH) ₄	NH _Z C1	A36
		* *	NH ₄ C1O ₄	A36
14 (A)	*	, af		
	a _{ex}		0. 10	e y se
NH ₄ C10 ₄	A1	HDB		A37
	Zr		* *	A38
	A1		ZrCl ₄	A38
* a	Zr	$Zr(OH)_4/(CH_2)_n$		A39
¥3		ZrH ₂ /Zr(OH) ₄	·	A40
		7		. "
				14.1
H ₂ WO ₄	ZrH ₂	Zr(OH) ₄	NH ₄ C10 ₄	A41
4	1 -	(CH ₂) _n		A42
*	*		ZrC1 ₄	A43
8 8 Y		79	ZrC1 ₄ /NH ₄ C10 ₄	A44
	Zr	HDB	NH ₄ C1O ₄	A45
		HDB/(CH ₂) _n		A46
	25	HDB	NH ₄ C1	A47
,	Zr/ZrH ₂	$Zr(OH)_4/(CH_2)_n$	NH ₄ C1	A48
19		2 y		,
8		2 2		
		. *		
140			1 1 1	l

APPENDIX A

INDEX OF TABLES (1)

(Continued - Page 3)

Propellant System

•		9 6		Table
Oxidizer	Fue1	Coolant	Chlorides	Number
		·		10
WO ₃	ZrH ₂	Zr(OH) ₄	NH ₄ C1O ₄	A49
J ,	2	(CH ₂) _n	7 7	A50
8		Z n	ZrCl ₄	A51
	Zr	ZrH ₂	N ₂ H ₅ C1	A52
		2	NH ₄ C10 ₄	A53
120 W		ZrH ₂ /Zr(OH) ₄	4 2 4	A54
•	e sk *·	HDB	NU C1	A55
RC (2)		ļ	NH ₄ C1	
	.00	HDB	NH ₄ C10 ₄	A56
€	Zr/ZrH ₂	H H		A57
46				
*	- 2	81 S		
ZrO(NO ₃) ₂ ·2H ₂ O	ZrH ₂	Zr(OH) ₄	·NH ₄ C1	A57
: :	a a	٠	NH ₄ C1O ₄	A58
<i>2</i>		(CH ₂) _n	NH ₄ C1	A59
	2		NH ₄ C10 ₄	A60
		·	4 4	
Sa.			13	. *
			s . H	
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·		<i>y</i>	20	
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8	,	·	19	.12
a a	,			141
e.	·			
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